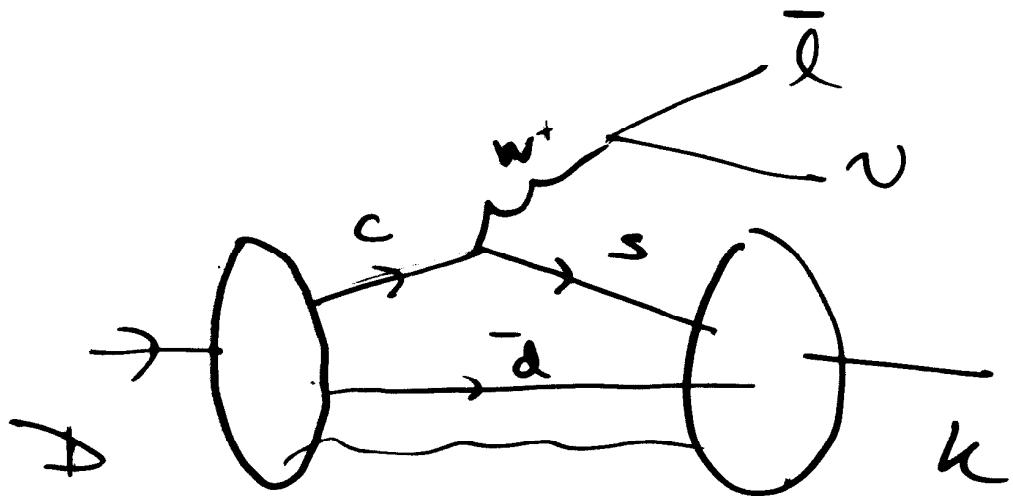


Light-Cone Wavefunction:

B, D

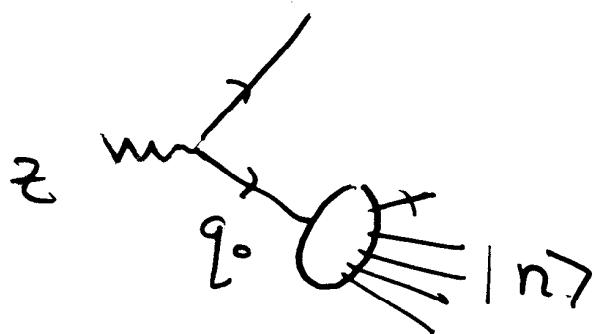
Exclusive decay



SFB  
Siegertsch.  
Hoyer

Jet hadronization in LCQ

(heuristic)



$$|q_0\rangle = \sum |n\rangle \langle n| q_0 \rangle$$

↑  
e. state  
 $H_{LC}^0$

↑  
e. state of  $H_{LC}$   
↑  
IR, UV regulated

Interesting to study in 1+1 theories  
collinear QCD

Hadronization at amplitude level!

# Proposal for DCCQ

OJR  
J.-K. Lin

## Prediagonalization

Diagonalize  $H_{\text{LC}}$  for  $q, \bar{q}$

$$I \rightarrow \left( \begin{array}{c} q_0 \\ \bar{q}_0 \end{array} \right) \langle \Psi_n \rangle \left( \begin{array}{c} q_0 \\ \bar{q}_0 \end{array} \right)^T$$

Regulate IR + UV :

$$\lambda^2 < \Delta m_n^2 < \Lambda^2$$

$$\Delta m_n^2 = \sum_{i=1}^n \frac{m_i^2 + \vec{k}_{\perp i}^2}{x_i} - M_q^2$$

Then use  $\{q, \bar{q}\}$  basis to diagonalize  
hadrons

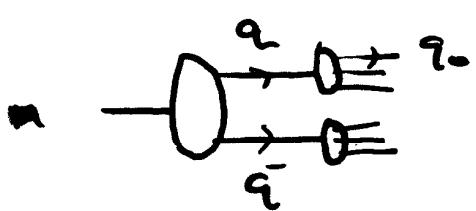
control

$$x \rightarrow 0$$

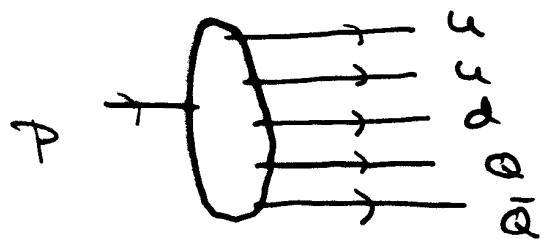
Ladder Reactions

Antonov, Dally  
 $\delta_{\delta 3}$

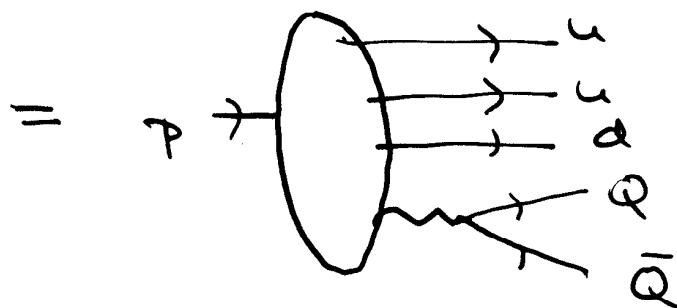
$q_0$  are "extrinsic" quarks



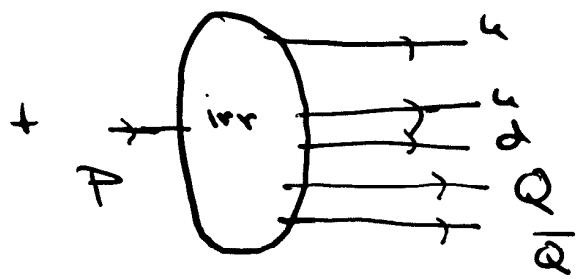
Two contributions to heavy quark sea!



$$Q = s, c, b, t$$

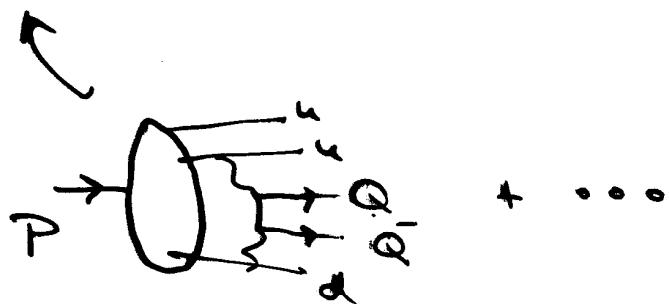


"extrinsic"  
subconstituents of  
gluon  
DGLAP evolution



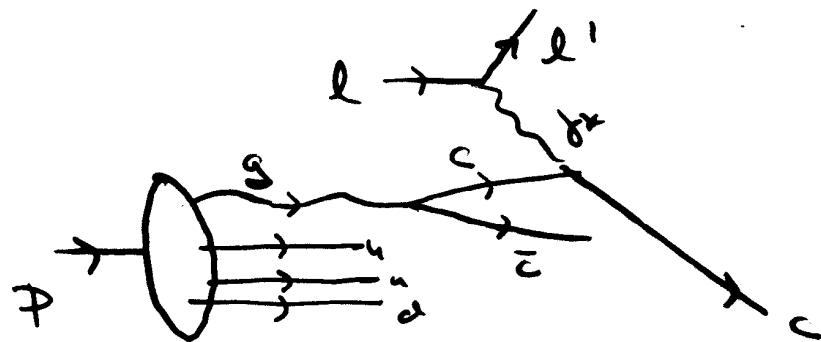
"intrinsic"  
1 gluon irreducible

(Physical source)  
 $A^+ = 0$

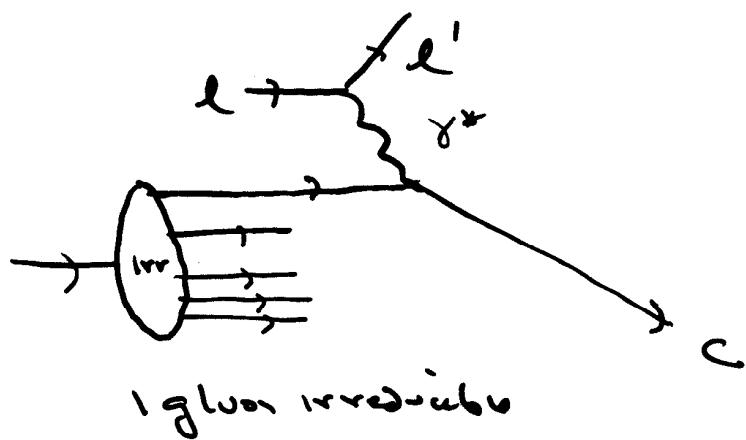


+ ...

## Two Contributions to Sea Quark Distributions



extrinsic =  
 photo-gluon  
 fusion  
 $\gamma g \rightarrow cc\bar{c}$



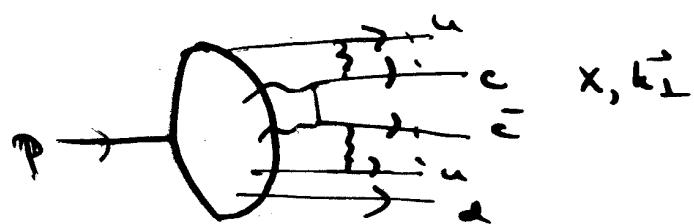
intrinsic  
 initial state  
 for DGLAP  
 evolution



$$Q^2 \gg 4m_c^2$$

$$c_I \propto \frac{1}{m_c^2 R^2}$$

Intrinsic



$$\sum x_i = 1$$

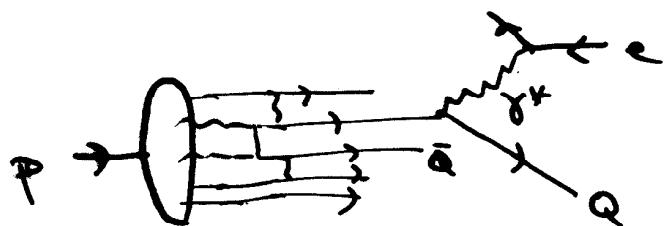
$$\sum k_{bi} = 0$$

$$\psi_{unacc}(x, k_\perp) \sim \frac{\Gamma(x, k_\perp)}{M_p^2 - \sum_{i=1}^n \frac{m_i^2 + k_{bi}^2}{x_i}}$$

wavefunction maximal at equal velocity (rapidity)

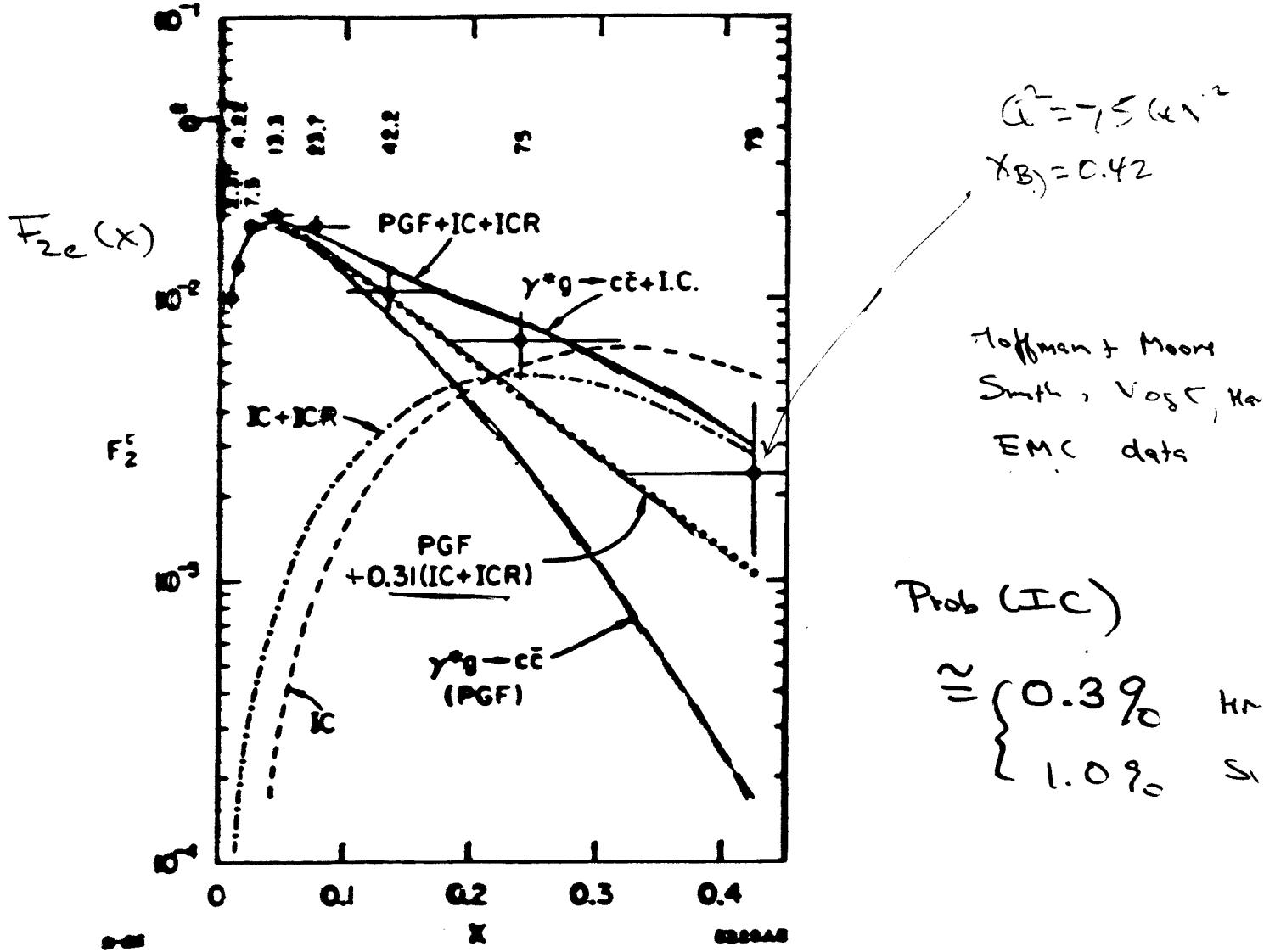
$$x_i = \frac{m_i}{\sum m_i}$$

$\therefore$  Heavy quarks tend to have largest momentum



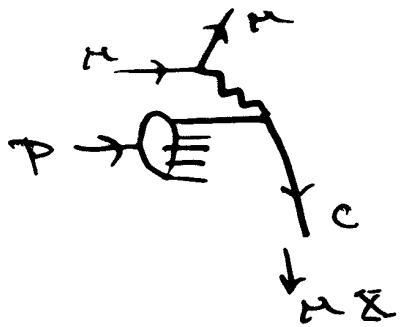
intrinsic heavy quark structure function

enhancement factor  $\frac{x}{1-e^{-x}}$ ,  $x = \frac{\pi C_F \alpha_s}{B}$



## Charm Puzzles

\* Charm Structure Function



$$c(x, Q^2) \approx 30 \text{ PDF !} \\ (\gamma^* g \rightarrow c\bar{c})$$

$$x = .42 \\ Q^2 = 75 \text{ GeV}^2$$

\*  $\pi N \rightarrow J/\psi J/\psi \cdot \chi$  NA3

$$\frac{d\sigma}{dx_F} \quad \underline{\text{only observed}} \text{ at } x_F > 0.4 !$$

\*  $\pi N \rightarrow J/\psi \chi$  NA3

$$\frac{d\sigma}{dx_F} = A^{.98} \frac{d\sigma^{\text{PDF}}}{dx_F} + A^{.77} \frac{d\sigma^{\text{diff}}}{dx_F}$$

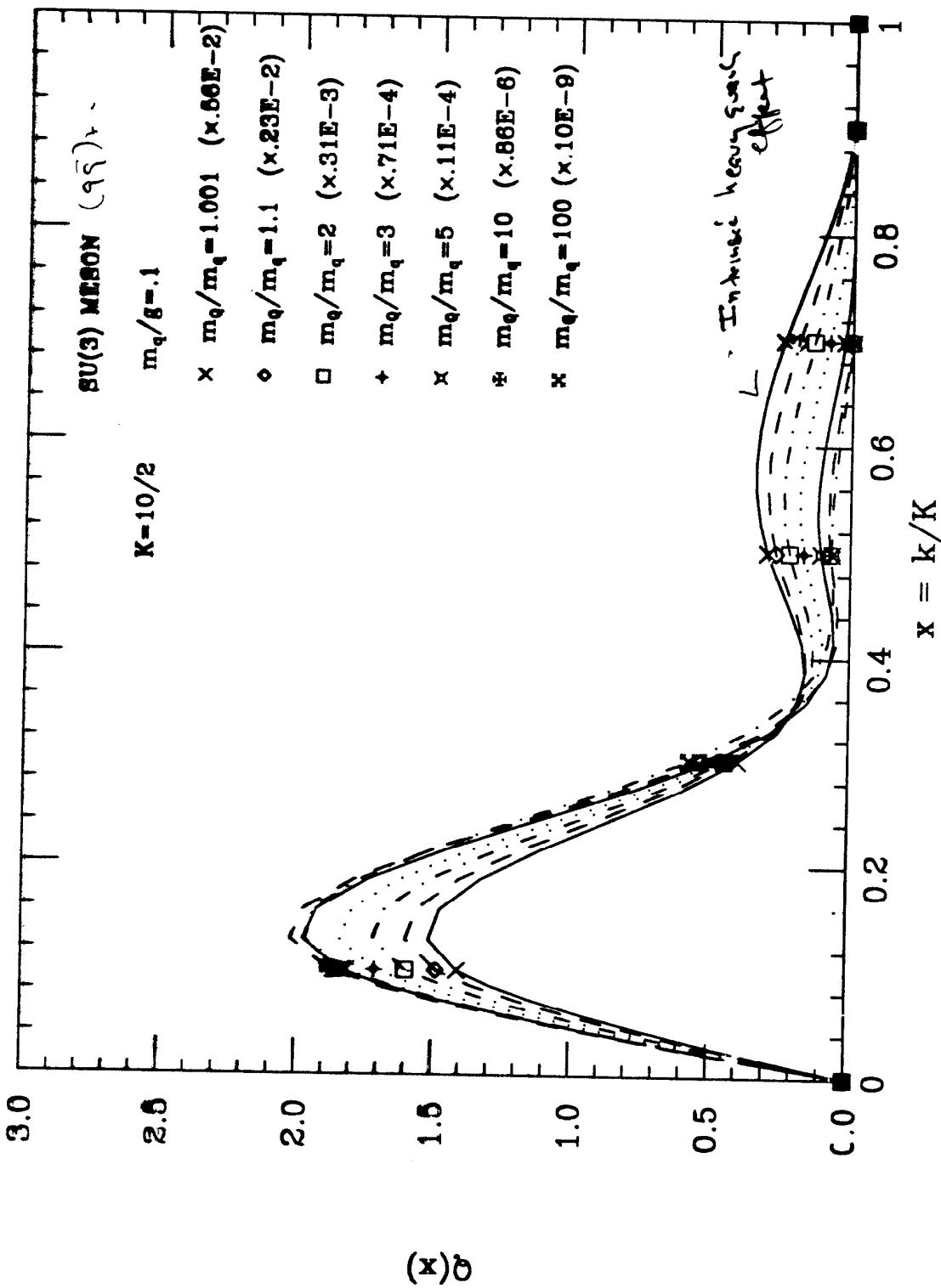
$A^{d(x_F)}$  not  $A^{d(x_2)}$   $\sim$  flat in  $x_F$ !

violates PDF factorization!

Hoyer, Sankar, Vanttre

# MOMENTUM DISTRIBUTION $q \bar{q} Q \bar{Q}$

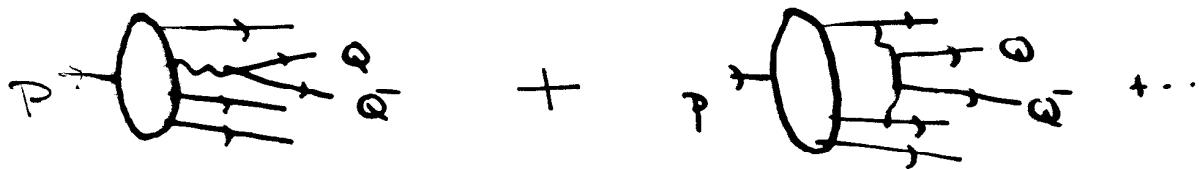
$\text{QCD}(1+1)$



Possible new source ~  
 $\delta, \alpha, b$   
 quarks at some  $x=1$

# Heavy Sea in Nucleon

Distinguishing characteristics:



extrinsic

intrinsic

$$P_{\text{ext}} \left( m_{Q\bar{Q}}^2 > 4M_Q^2 \right)$$

$$P_{\text{int}} \left( M_{\text{qq}}^2 > 4M_Q^2 \right)$$

Probability  $\sim \frac{\alpha_s}{\pi} \log \frac{m_{Q\bar{Q}}^2}{4M_Q^2} P_0$

high  
momentum  
probe)

$$\sim \frac{\alpha_s^2}{R^2 m_{Q\bar{Q}}^2} P_{\text{BS}}$$

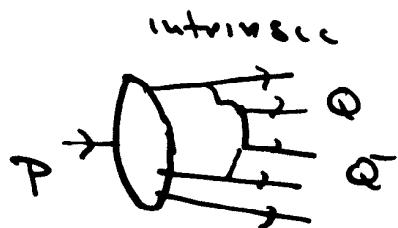
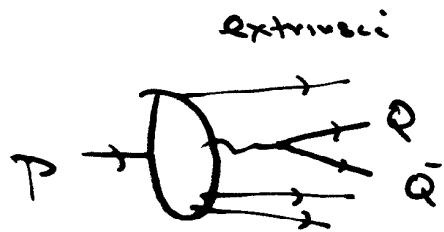
$\uparrow$   
interquark  
separation

$x_Q$  dist  $\left\{ \begin{array}{l} \langle x_Q \rangle < \langle x_g \rangle \\ (1-x_Q)^{5+7} \end{array} \right.$

$$\left\{ \begin{array}{l} \langle x_Q \rangle \sim \frac{m_{Q\perp}}{\sum m_\perp} \\ \text{high } x_Q! \end{array} \right.$$

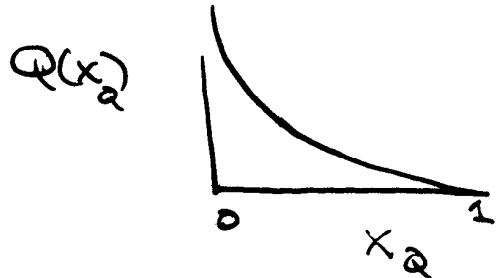
# Distinguishing Properties

$Q = S, C, b, t$



$$Q(x) \underset{\text{ext}}{\sim} (1-x) g(x)$$

$$\sim \frac{1}{x^2} (1-x)^{5-7}$$



$$Q_{int}(x) :$$

$$\text{peaks at } x_Q \sim \frac{m_1}{\sum m_i}$$



$$Q(x) \equiv \bar{Q}(x)$$

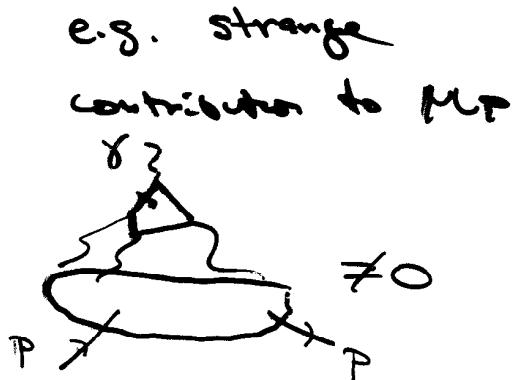
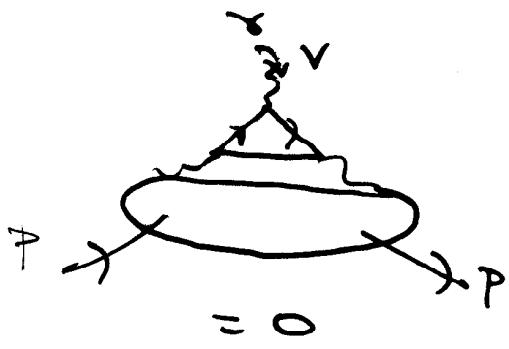
$Q(x)$  and  $\bar{Q}(x)$   
can be different!

---

zero contribution to  
vector current

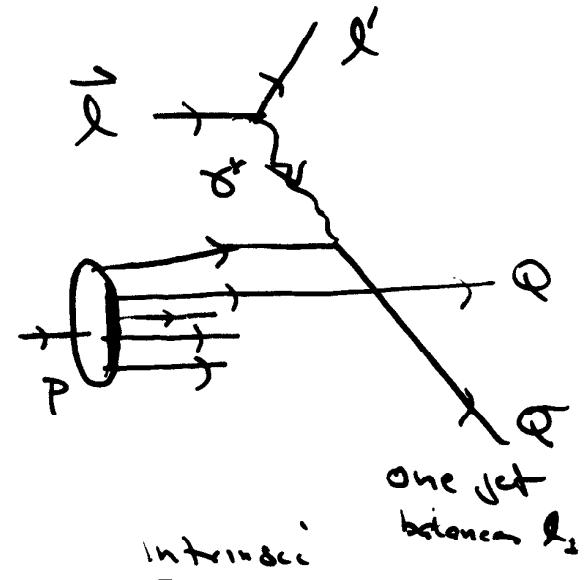
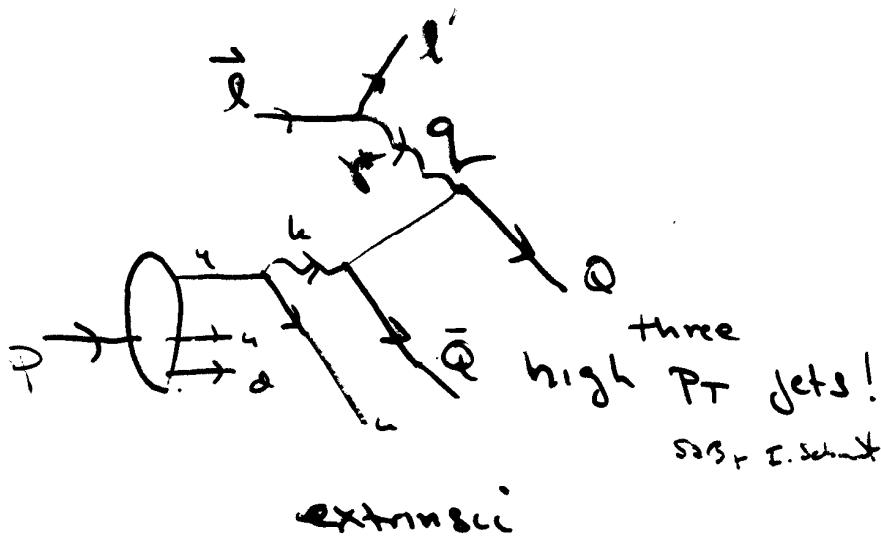
---

non-zero contribution  
vector current



# Distinguishing Characteristics

Axial Current  
and Polarized DIS



$$Q^2 \gg |k^2| \gg 4m\bar{Q}^2$$

$$Q^2 \gg 4m\bar{Q}^2$$

$$\Delta Q = \Delta \bar{Q}$$

$$\Delta Q \neq \Delta \bar{Q} !$$

$$\Delta Q_{TOT} = -\frac{\alpha_s}{2\pi} \Delta g$$

$\Delta S$  negative for  
 $P \rightarrow k\Lambda$

"anomalous contribution"

$$\Delta \bar{S} = 0$$

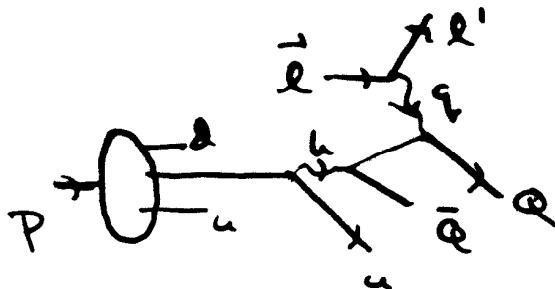
( $A^+ = 0$  gauge)

$$\Delta Q_{TOT} \propto \frac{1}{R^2 m_q}$$

$\Delta Q_{TOT}$  independent of  $m_q^2$ !  
 $Q^2 \gg |k^2| \gg 4m\bar{Q}^2$

so Two contributions to E-J Son Rule  
high  $Q^2$

Extrinsic



$$\Delta S_{\text{ext}} = -\frac{\alpha}{2\pi} \Delta q$$

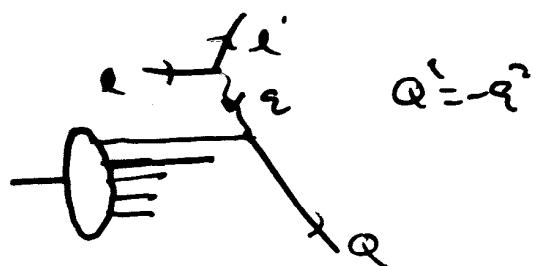
$$\Delta \bar{Q}_{\text{ext}} = \Delta Q_{\text{ext}}$$

$$\Delta C_{\text{ext}} = -\frac{\alpha}{2\pi} \Delta q$$

$$\Delta b_{\text{ext}} = -\frac{\alpha}{2\pi} \Delta q$$

if  $Q^2 \gg \Lambda^2 \gg 4m_e^2$

Intrinsic



$$\Delta Q_{\text{int}} \neq \Delta \bar{Q}_{\text{int}}$$

$$\Delta Q_{\text{int}} \sim \frac{1}{M_Q^2 R^2}$$

$$\frac{\Delta S_{\text{int}}}{\Delta C_{\text{int}}} \sim \frac{m_c^2}{m_S^2}$$

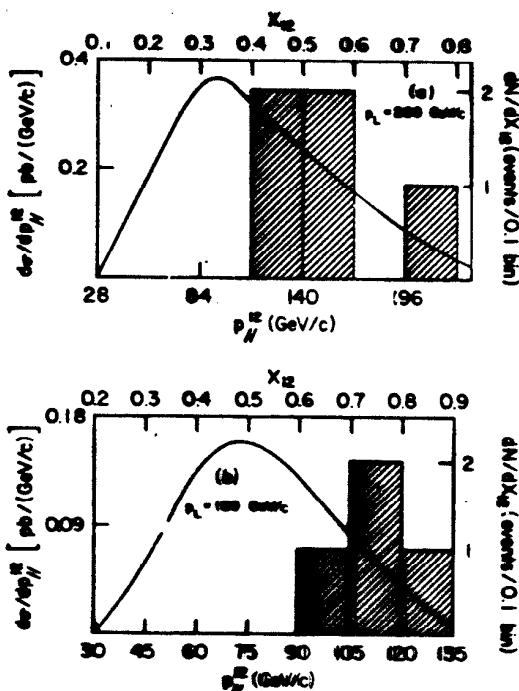
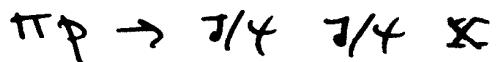
$Q^2 \gg 4m_\phi^2$

Phenomenology: perhaps roughly equal!

look at  $x_c$ -dist  $(x_c \text{ range})$   
# jet recoiling against  $\ell$

Note: resolution dependence!

Production of two  $J/\psi$



all events  
Total  $X_F \geq c. 4$  !

$$P_{\text{Lab}} = 280 \text{ GeV}/c$$

$$P_{\text{Lab}} = 150 \text{ GeV}/c$$

FIG. 5. The calculated longitudinal differential cross section  $d\sigma/dp_{TJ}^L$  for the  $J/\psi$  pair in  $\pi p \rightarrow J/\psi X$  at (a)  $p_L = 200 \text{ GeV}/c$  and (b)  $p_L = 150 \text{ GeV}/c$ , with experimental data. The scales for the experimental events are to the right of the figure.

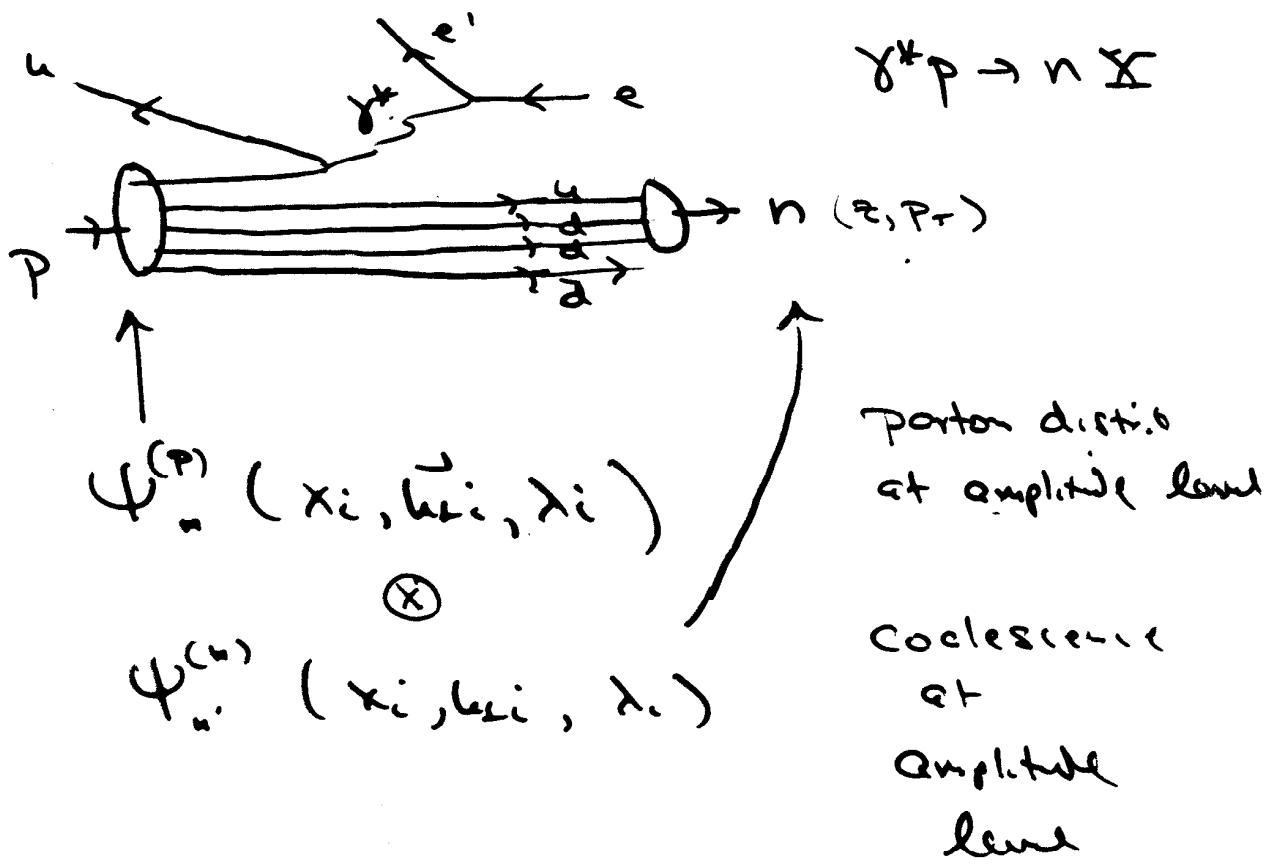
NA3 data

( $L_c + L_{c\bar{c}}$ )

$$\frac{\partial \sigma}{\partial x_{12}} (\pi p \rightarrow J/\psi_1 J/\psi_2 X)$$

# Light-Cone Wavefunctions and "Fracture Functions"

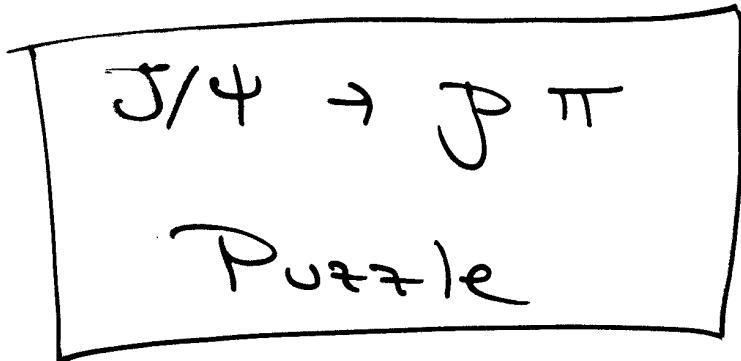
Veneziano  
Trentadue



$$\mathcal{M}^\mu \sim \langle \text{pl. } j^\mu(o) \ln u \bar{a} \rangle$$

extension  
of F.F.

SJB +  
M. Karkinen

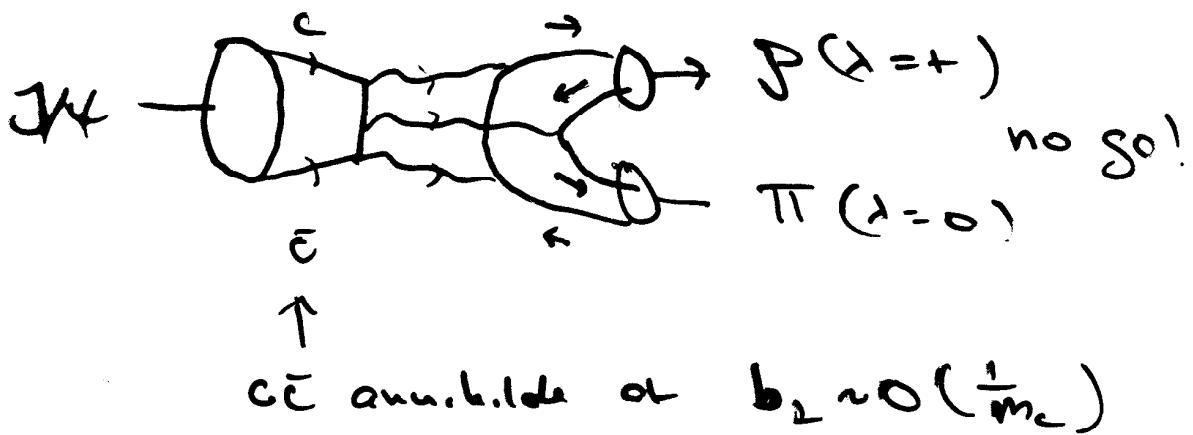


\*  $B[J/\psi \rightarrow J\pi] = 1.28 \pm 0.10\%$

largest two-body  
hadronic decay!

\* PQCD:  $J/\psi \rightarrow PV$  suppressed  
by helicity conservator

SJB, Tuc., Lepos



ref PDG  
6-92

## J/4 → p̄π Puzzle

$$\text{BR} (\text{J}/4 \rightarrow h^+ h^-) = 2.37 \pm 0.31 \times 10^{-4}$$

not suppressed! →

$$p\bar{\pi} = 1.28 \pm 0.10 \times 10^{-2}$$
$$h^+ h^{*\pm} = 5.0 \pm 0.4 \times 10^{-3}$$

$$\text{BR} (\psi' \rightarrow h^+ h^-) = 1.0 \pm 0.7 \times 10^{-4}$$

$$p\bar{\pi} < 8.3 \times 10^{-5} \quad (90\%)$$

$$h^+ h^{*\pm} < 1.8 \times 10^{-5} \quad (90\%)$$

Expect suppression ↴ vector + pseudoscalar

ψ' okay

J/4 disaster!

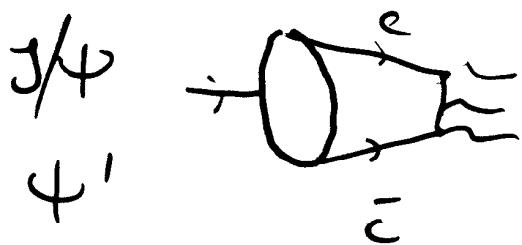
\* Strong violation of hadron helicity conservation

\* Extremely fast quenching of  $p\bar{\pi}$ ,  $h^+$  signal  
See: Chidian + Torgquist.

PQCD :  $e\bar{e}$  annihilation

Expect  $T_{ns} \propto |\Psi_s(\vec{0})|^2$

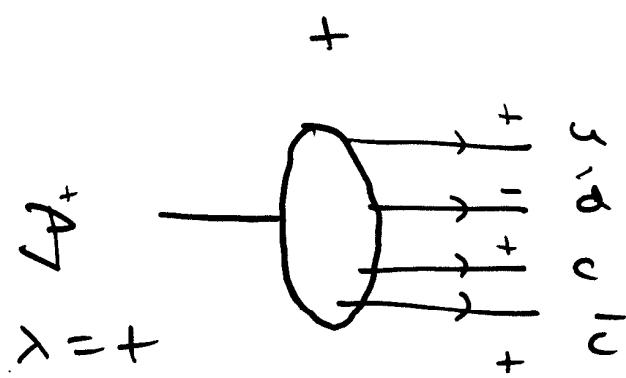
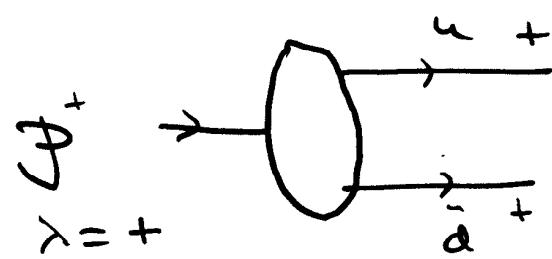
$$\frac{\Gamma_{2S}}{\Gamma_{1S}} \sim \frac{\Gamma_{+^3 \rightarrow e\bar{e}}}{\Gamma_{+^3 \rightarrow e\bar{e}}}$$



However  $\psi' \rightarrow \rho\pi$   
never observed!

$$B \psi' \rightarrow \rho\pi < \frac{1}{50} \quad \begin{matrix} \text{PQCD} \\ \text{expect} \end{matrix}$$

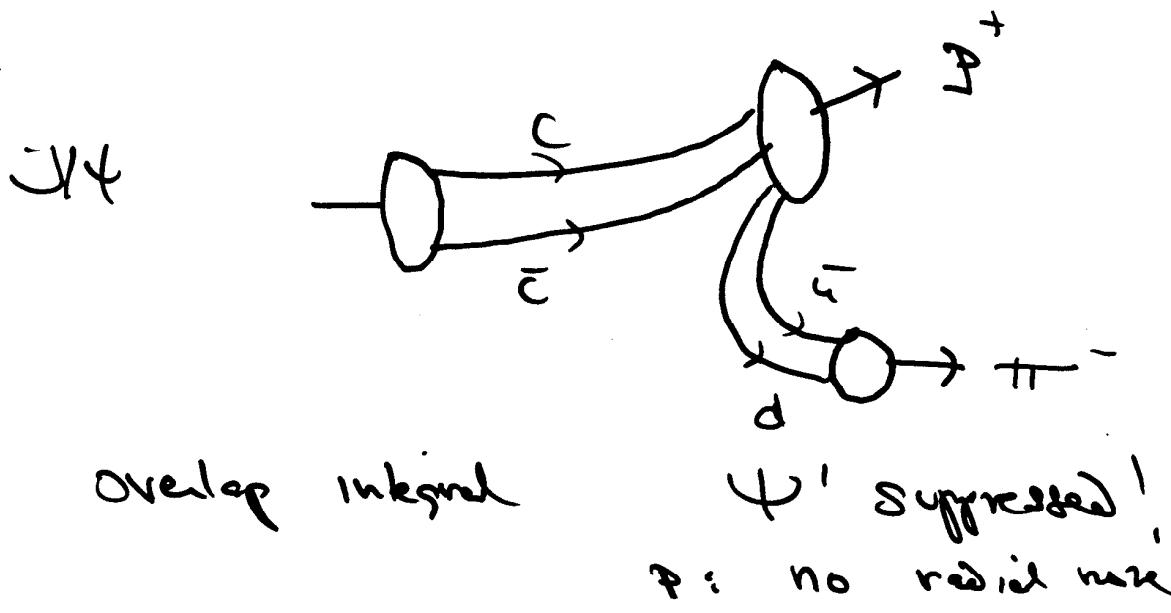
$\uparrow$   
 $< 3.6 \times 10^{-5}$



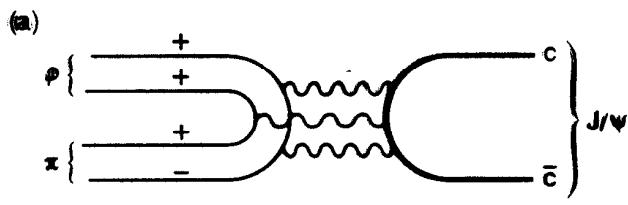
minimize

$$m^2$$

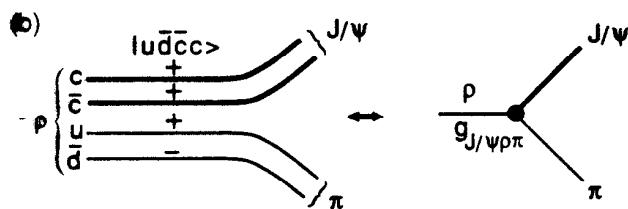
$$\beta \sim \pi/4 \text{ rad}$$



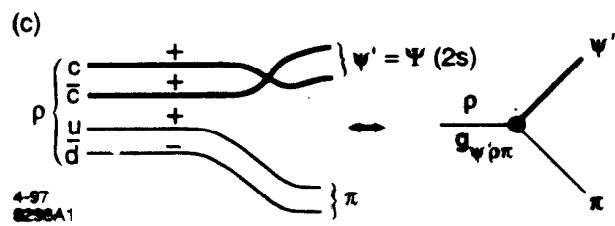
radial  
node



usual  
QCD  
ann. let.



$c\bar{c}$   
rearrangement  
 $\not\rightarrow J/\psi \pi$

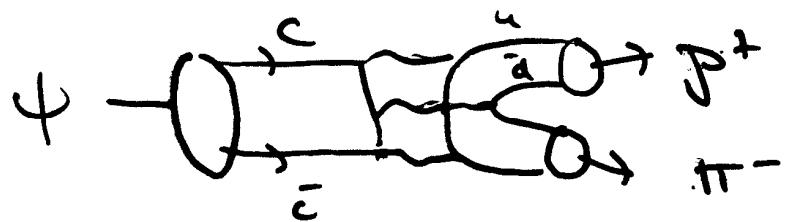


suppressed  
 $\not\rightarrow \psi' \pi$   
coupling

JP<sub>π</sub> Puzzle

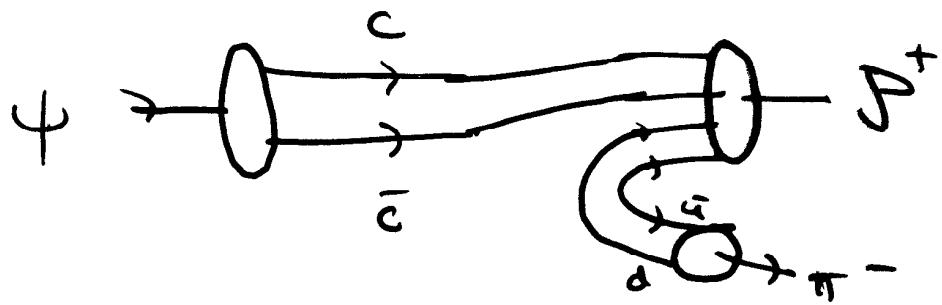
[1S]  $J/\psi \rightarrow J/\psi \pi$   $1.28 \pm 0.10 \%$

[2S]  $\psi' \rightarrow J/\psi \pi$   $B < \frac{1}{50}$  expectations!  
 $< 3.6 \times 10^{-5}$



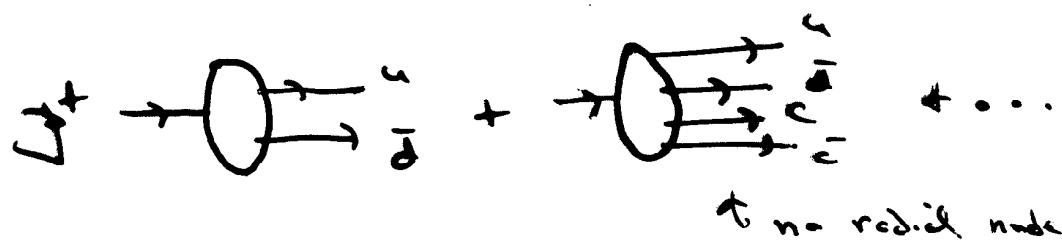
$$T \sim |\psi_{ns}(0)|^2$$

SJBB  
H. Kowalski



$\psi' \neq J/\psi$   
expressed  
(radial  
node)

rearrange  $c\bar{c}$  into 4-quark Fock state  $J/\psi$

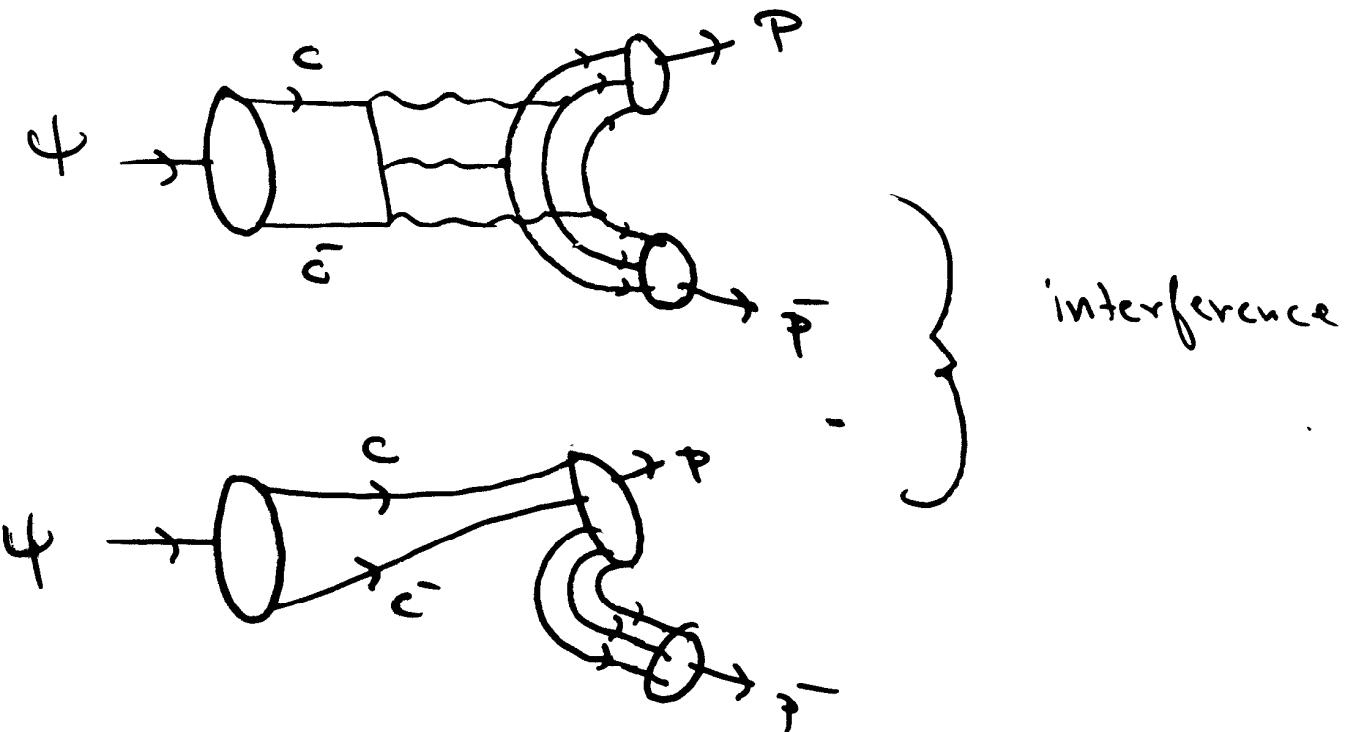


## Intrinsic Charm in light-hadrons

$$P_{Q\bar{Q}} \sim \frac{1}{m_{Q\bar{Q}}} \alpha_s^2 (M_{Q\bar{Q}}) \quad \text{scaling} \quad \bar{s}\bar{s} : \bar{c}\bar{c} : \bar{b}\bar{b}$$

$\Rightarrow \phi \rightarrow p\pi \text{ vs } \psi \rightarrow p\pi$

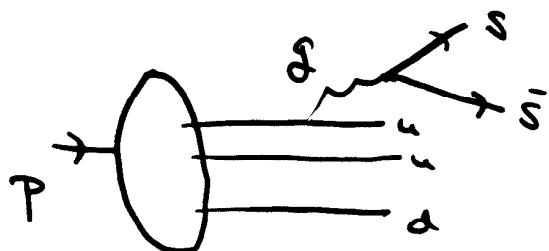
## Rearrangement plus annihilation



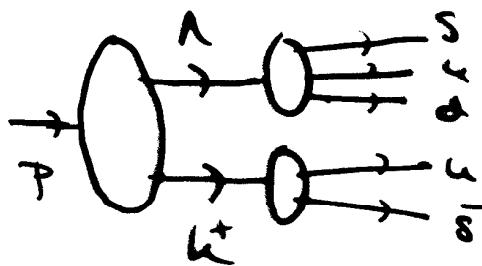
complicates  $B, \Upsilon$  decays.

Anti-symmetry of strange sea

$$S(x) \stackrel{?}{=} \bar{S}(x)$$



yes!



no!

Thomas et al  
SJS + Ma

Nielsen  
Neumann  
et al

In LN model

s helicity opposite to p helicity

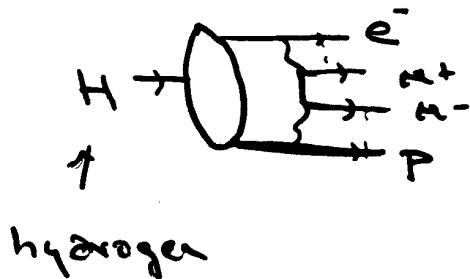
$$S(x) \neq \bar{S}(x)$$

probes new physics

a) \* non-perturbative QCD effect

\* k-T fluctuations

\* QED analog:



Coulomb distortion

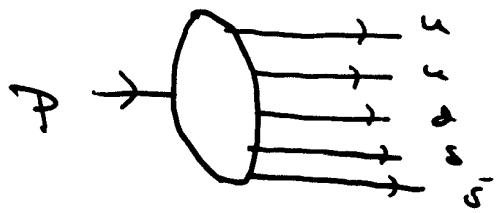
Oppositely charged  
have low  $V_{rel}$ 
 $\Rightarrow$  asymmetric  $\alpha^+, \alpha^-$   
 $\langle x_{\alpha^-} \rangle > \langle x_{\alpha^+} \rangle$ 
b) \* non-perturbative hadronization  
of jets

$$D_{p/s} \neq D_{\bar{p}/s}$$

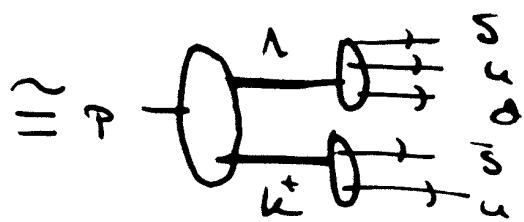
c) \* coalescence effects

$$\Delta S \neq \Delta \bar{S}$$

e) \* discriminate intrinsic vs extrinsic  
contributions



"intrinsic strangeness"



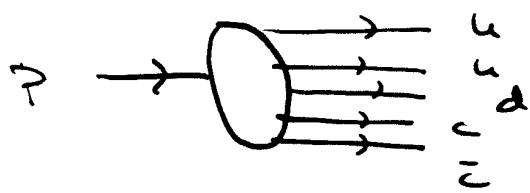
asymmetric distribution

$$\langle x_s \rangle > \langle x_{\bar{s}} \rangle$$

helicity of 8 anti-aligned!

$$\Delta S < 0, \Delta \bar{S} = 0$$

minimum  $M^2$  dominant



"intrinsic charm"

$$\langle x_c \rangle, \langle x_{\bar{c}} \rangle \text{ large}$$

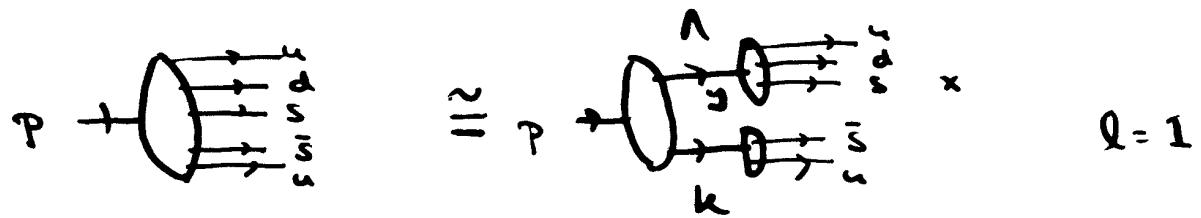
minimum  $M^2$  dominant  $\Rightarrow$

$$P_{Q\bar{Q}} \sim \frac{1}{m_{Q\bar{Q}}^2}$$

- $\hookrightarrow \left\{ \begin{array}{l} c(x, Q), b(x, Q) \text{ at large } x \\ J/\psi, \psi\psi, \Lambda_c, \Lambda_b \text{ at large } x_F \\ J/\psi \rightarrow p\pi \text{ puzzle} \\ \text{HERA anomaly} \end{array} \right.$

# Model for Intrinsic Strongness

SJR + Bo-Qiang Ma  
 SLAC-PUB-7126  
 PL



$$S(x) = \int_x^1 \frac{dy}{y} f_{\alpha/\mu+\lambda}(y) Q \sin\left(\frac{x}{y}\right)$$

$$\Psi(x_i, \omega_i) = \begin{cases} A e^{-m^2/2\alpha^2} \\ A (1+m^2/\alpha^2)^{-p} \end{cases}$$

$$m^2 = \sum \frac{\omega_i^2 + m^2}{x}$$

$p=3.5$

estimate:  $P_{K\Lambda} \sim 4 \text{ to } 15\% : \underline{\langle x_s \rangle > \langle x_{\bar{s}} \rangle}$   
 $K\Sigma, K^* \Lambda$  negligible

$$\Delta S_s = -\frac{1}{3} P_{K\Lambda} \quad \Delta S_{\bar{s}} = 0$$

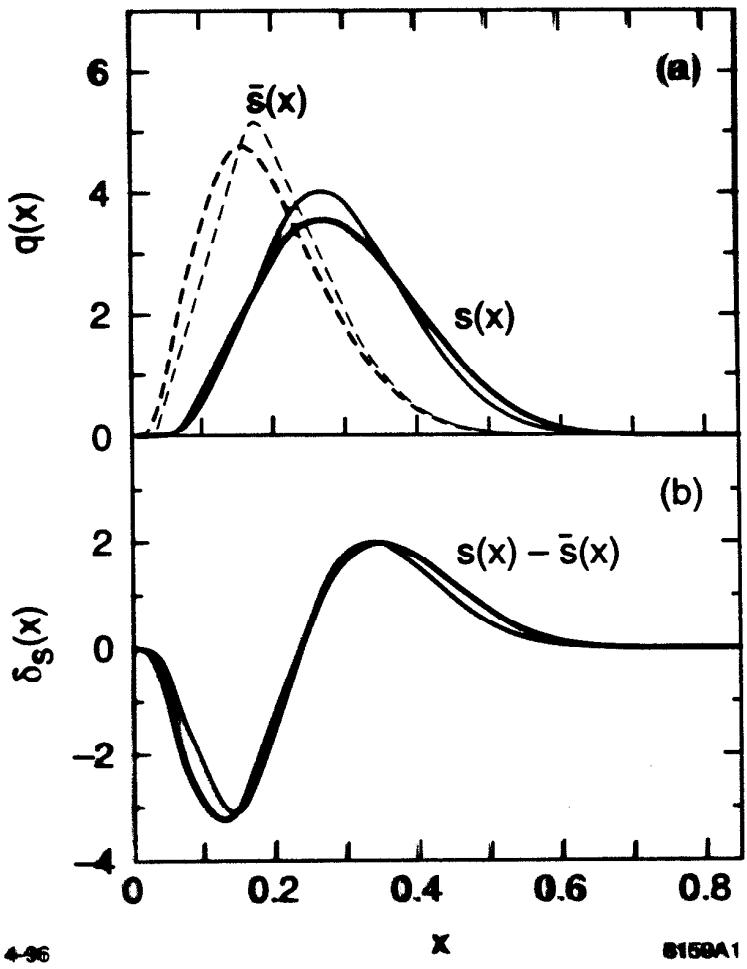
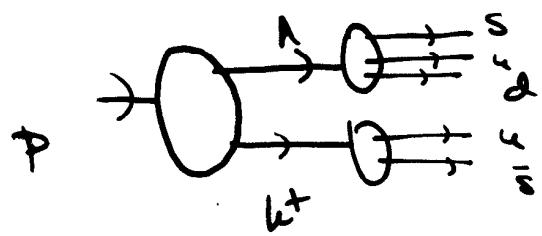


Figure 1: The momentum distributions for the strange quarks and antiquarks in the light-cone meson-baryon fluctuation model of intrinsic  $q\bar{q}$  pairs, with the fluctuation wavefunction of  $K^+\Lambda$  normalized to 1. The curves in (a) are the calculated results of  $s(x)$  (solid curves) and  $\bar{s}(x)$  (broken curves) with the Gaussian type (thick curves) and power-law type (thin curves) wavefunctions and the curves in (b) are the corresponding  $\delta_s(x) = s(x) - \bar{s}(x)$ . The parameters are  $m_q = 330$  MeV for the light-flavor quark mass,  $m_s = 480$  MeV for the strange quark mass,  $m_D = 600$  MeV for the spectator mass, the universal momentum scale  $\alpha = 330$  MeV, and the power constant  $p = 3.5$ , with realistic meson and baryon masses.

$$\text{If } G_{s/p}(x) \neq G_{\bar{s}/p}(x)$$

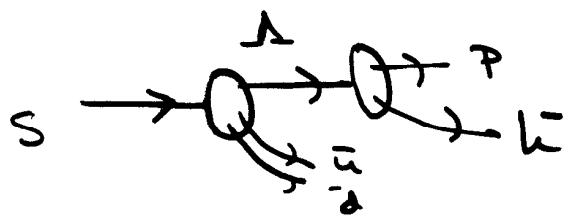
$$\text{Then } D_{p/s}(z) \neq D_{\bar{p}/s}(z)$$

$$D_{p/s}(z) \neq D_{\bar{p}/s}(z)$$



$$G_{s/p}(x) > G_{\bar{s}/p}(x) \quad \text{at large } x$$

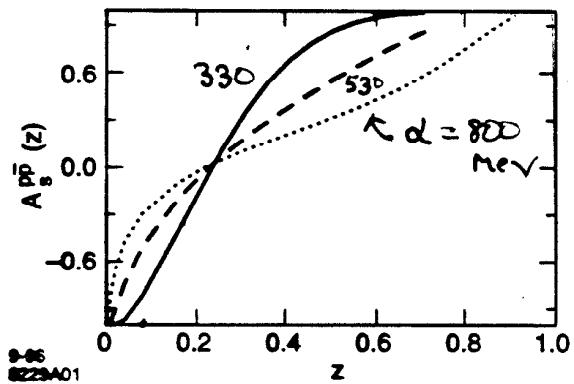
∴  $\swarrow$  lowest mass intermediate state



$$D_{p/s}(z) > D_{\bar{p}/s}(z) \quad \text{at large } z$$

$$A_s^{pp}(z) = \frac{D_{p/s}(z) - D_{\bar{p}/s}(z)}{D_{p/s}(z) + D_{\bar{p}/s}(z)}$$

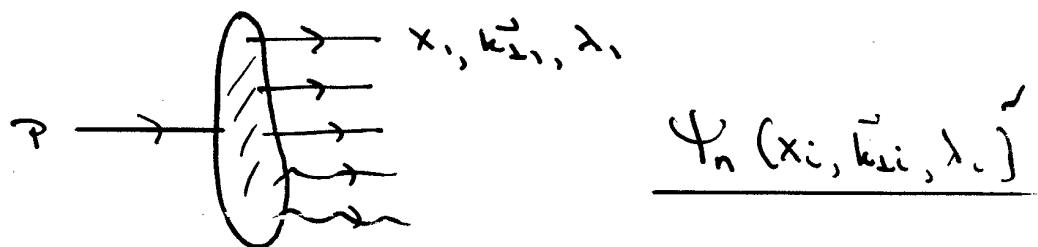
$$= \frac{G_{s/p}(x) - G_{\bar{s}/p}(x)}{G_{s/p}(x) + G_{\bar{s}/p}(x)}$$



Model for  $A_s^{p\bar{p}}(z)$   $\Psi = C e^{-m^2/2\alpha^2}$

$$A_s^{p\bar{p}}(z) = \frac{D_{p/s}(z) - D_{\bar{p}/s}(z)}{D_{p/s}(z) + D_{\bar{p}/s}(z)}$$

## Role of LC Wavefunctions in QCD



- \* Precise realization of Quark Parton Model
- \* Representations of { Leading Twist } Structure Functions  
+ { Non-Leading Twist }
- \* Representations of Exclusive Amplitudes  
including Form Factors, Weak Trans. Amplit.  
from low to high  $Q^2$
- \* Definition of Higher-Twist Coefficient Functions
- \* Densite Factorization Theorems  $(A^+ = 0 \text{ gauge})$   
Evolution Equations, R.G.E.
- \* Translation of Hadron to  $q + g$  degrees of freedom
- \* Novel Effects : Color Transparency  
Intrinsic gluons, heavy quark

## Summary / Outlook

$\text{QCD} + \text{LCQ} \Rightarrow$

relativistic many-body theory

boost-invariant

trivial vacuum

Symmetry breaking } zero nodes  
 $m_\nu \neq m_\mu$

$\text{DLCQ} \Rightarrow$

complete solns to quantum field theory!

spectrum, scattering, states,  $\Phi_n(x, k_\mu, p_i)$

exclusive, inclusive n.e., OPE

new directions  $\rightarrow B$ -decays, fragmentation,  $\Xi_L = \sum e^{-m_B/p_i^2}$

Novel QCD Phenomenology

## Summary:

Two contributions to heavy quark sea!

- \* LC. Wavefunctions, analyses
  - \* high  $x_c$  EMC data
  - \*  $\Delta q = \Delta q_{\text{int}} + \Delta q_{\text{extrinsic}}$   
    ↑  
     anomaly
  - \*  $\Delta q_{\text{int}} \neq \bar{\Delta q}_{\text{int}}$ ,  $q(x) \neq \bar{q}_{\text{int}}$
  - \* hadroproduction at large  $X_F$ , A-dep., 1
  - \* HERA: proton fragmentation
  - \* coalescence: leading particle effects
  - \* asymmetric hadronization
- $c\bar{c}$  { charm + strangeness together:  
                crucial probe of hadron structure