

**The Meson Mass Spectra and Mixing Angles of  $(\omega, \phi)$   
and  $(\eta, \eta')$  in the Light-Cone Quark Model**

Ho-Meoyng Choi and Chueng-Ryong Ji

Department of Physics

North Carolina State University

---

## Contents

**I.** Motivation

**II.** Quadratic mass formula and  $SU(3)$  breaking

**III.** Effective potential in the light-cone quark model

**IV.** Mass spectra of light pseudoscalar and vector mesons

**V.** Summary and Discussion.

## I. Motivation

1. How to resolve the sign of the puzzle of  $\omega - \phi$  mixing angle in LCQM,<sup>\*</sup>  
 $\delta_V = -3.3^\circ$ , which is OPPOSITE to the usual convention that is based upon  
the quadratic SU(3) mass formula(Gell-Mann-Okubo formula) ?

- Since the LCQM[1] fitting of the  $\omega - \phi$  mixing angle has included the effect of SU(3) breaking, *i.e.*,  $m_s/m_{u(d)} \sim 1.5$ , it would be the right way to analyze the meson mass spectra based on the SU(3) breaking not on the Gell-Mann-Okubo(GMO) mass formula, which does not account for SU(3) breaking.

2. How to implement a dynamical mechanism on LCQM<sup>\*</sup> ?

- Set up three different types of potentials to analyze the mass spectra of light pseudoscalar( $\pi, K, \eta, \eta'$ ) and vector( $\rho, K^*, \omega, \phi$ ) mesons:

(1) Harmonic Oscillator(HO) + spin-spin(S-S)

(2) Coulomb + HO + S-S

(3) Coulomb + linear confining + S-S

\* W. Jaus, Phys. Rev. D 44, 2851 (1991)

→ Analyze the meson mass spectra!!

## II. Quadratic mass formula and SU(3) breaking<sup>\*,\*\*</sup>

\* De Rújula, Georgi, and Glashow \*\* N. Isgur

The generic two particle ( $f_1, f_2$ ) mixing is given by

$$\begin{aligned} |f_1\rangle &= -\sin\delta|n\bar{n}\rangle - \cos\delta|s\bar{s}\rangle \\ |f_2\rangle &= \cos\delta|n\bar{n}\rangle - \sin\delta|s\bar{s}\rangle \end{aligned} \quad (2.1)$$

where  $n\bar{n} = (u\bar{u} + d\bar{d})/\sqrt{2}$  and  $\delta = \theta_{SU(3)} - 35.26^\circ$ . We shall identify ( $f_1, f_2$ ) with ( $\phi, \omega$ ) and ( $\eta, \eta'$ ). The quadratic mass eigenvalue equation is given by

$$\mathcal{M}^2|f_i\rangle = M_i^2|f_i\rangle \quad (i = 1, 2), \quad (2.2)$$

with

$$\mathcal{M}^2 = \begin{pmatrix} M_{n\bar{n}}^2 + 2\lambda & \sqrt{2}\lambda X \\ \sqrt{2}\lambda X & M_{s\bar{s}}^2 + \lambda X^2 \end{pmatrix}. \quad (2.3)$$

The parameter  $\lambda$  characterizes the strength of the quark-annihilation graph which couples the  $I=0$   $u\bar{u}$  state to  $I=0$   $u\bar{u}, d\bar{d}, s\bar{s}$  states with equal strength in the SU(3) limit. The parameter  $X$  pertains to SU(3) breaking, *i.e.*,  $X \rightarrow 1$  in the SU(3) limit, such that the quark-annihilation graph *factors* into its flavor parts, with

$$\lambda : u\bar{u} \rightarrow u\bar{u}(d\bar{d})$$

$$\lambda X : u\bar{u} \rightarrow s\bar{s} \text{ (or } s\bar{s} \rightarrow u\bar{u})$$

$$\lambda X^2 : s\bar{s} \rightarrow s\bar{s}$$

Solving Eqs.(2.1)-(2.3), we obtain

$$M_{f_1}^2 = (M_{n\bar{n}}^2 + 2\lambda) \sin^2 \delta + 2\sqrt{2}\lambda X \sin \delta \cos \delta + (M_{s\bar{s}}^2 + \lambda X^2) \cos^2 \delta, \quad (2.4)$$

$$M_{f_2}^2 = (M_{n\bar{n}}^2 + 2\lambda) \cos^2 \delta - 2\sqrt{2}\lambda X \sin \delta \cos \delta + (M_{s\bar{s}}^2 + \lambda X^2) \sin^2 \delta, \quad (2.5)$$

and

$$\tan^2 \delta = \frac{M_{f_2}^2 - \epsilon M_{f_1}^2}{\epsilon M_{f_2}^2 - M_{f_1}^2}, \quad (2.6)$$

where  $\epsilon = (M_{n\bar{n}}^2 + 2\lambda)/(M_{s\bar{s}}^2 + \lambda X^2)$ . The constraint requires the invariance of the trace of quadratic mass matrix:

$$M_{n\bar{n}}^2 + M_{s\bar{s}}^2 + 2\lambda + \lambda X^2 = M_{f_1}^2 + M_{f_2}^2. \quad (2.7)$$

**How to determine the sign of  $\delta$  in quark model?**

$$\tan 2\delta = \frac{2\sqrt{2}\lambda X}{M_{s\bar{s}}^2 - M_{n\bar{n}}^2 + \lambda X^2 - 2\lambda},$$

(A)  $\omega - \phi$  mixing with  $\delta_V = -3.3^\circ$  (or  $\theta_{SU(3)} = 31.96^\circ$ )

$$\lambda_V < 0, \quad X_V > 0 \quad (?) \quad \text{or}$$

$$\lambda_V > 0, \quad X_V < 0 \quad (?)$$

while for  $\delta_V = +3.3^\circ$

$$\lambda_V X_V > 0$$

(B)  $\eta - \eta'$  mixing with  $\delta_P = -54.26^\circ$  (or  $\theta_{SU(3)} = -19^\circ$ )

$$\lambda_P X_P > 0$$

### III. Effective potential in the light-cone quark model

The effective Hamiltonian for the description of the meson mass spectra is defined by

$$H_{q\bar{q}} = \sqrt{m_q^2 + k^2} + \sqrt{m_{\bar{q}}^2 + k^2} + V_{q\bar{q}}, \quad (3.1)$$

and the mass-squared operator is given by

$$M_{\text{Op}}^2 = 4(m^2 + k^2 + mV_{q\bar{q}}). \quad (3.2)$$

Three different types of potentials are as follows

$$(A) \quad V_{q\bar{q}} = \underbrace{V_0}_{V_{S.H.O}} + \underbrace{V_{\text{I}}}_{V_{S-S}} \\ = A_1 + B_1 r^2 + \frac{\vec{S}_q \cdot \vec{S}_{\bar{q}}}{m_q m_{\bar{q}}} \nabla^2 V_{Coul}, \quad (3.3)$$

$$(B) \quad V_{q\bar{q}} = V_{Coul} + V_{S.H.O} + V_{S-S} \\ = A_2 - \frac{\kappa}{r} + B_2 r^2 + \frac{\vec{S}_q \cdot \vec{S}_{\bar{q}}}{m_q m_{\bar{q}}} \nabla^2 V_{Coul}, \quad (3.4)$$

$$(C) \quad V_{q\bar{q}} = V_{Coul} + V_{lin} + V_{S-S} \\ = A_3 - \frac{\kappa}{r} + B_3 r + \frac{\vec{S}_q \cdot \vec{S}_{\bar{q}}}{m_q m_{\bar{q}}} \nabla^2 V_{Coul}, \quad (3.5)$$

where  $\nabla^2 V_{Coul} = 4\pi\kappa\delta^3(r)$ .

Using the ground-state  $S$ -wave variational wavefunction,

$$\phi(k^2) = \left(\frac{4}{\sqrt{\pi}\beta^3}\right)^{1/2} \exp(-k^2/2\beta^2), \quad (3.6)$$

we obtain

$$\langle V_{q\bar{q}} \rangle = A_1 + \frac{3}{2\beta_{q\bar{q}}^2} B_1 + 4\kappa \frac{\beta_{q\bar{q}}^3}{m_q^2 \sqrt{\pi}} \vec{S}_q \cdot \vec{S}_{\bar{q}}, \quad \text{for A} \quad (3.7)$$

$$= A_2 - \kappa \frac{2\beta_{q\bar{q}}}{\sqrt{\pi}} + \frac{3}{2\beta_{q\bar{q}}^2} B_2 + 4\kappa \frac{\beta_{q\bar{q}}^3}{m_q^2 \sqrt{\pi}} \vec{S}_q \cdot \vec{S}_{\bar{q}}, \quad \text{for B}, \quad (3.8)$$

$$= A_3 - \kappa \frac{2\beta_{q\bar{q}}}{\sqrt{\pi}} + \frac{2}{\beta_{q\bar{q}} \sqrt{\pi}} B_3 + 4\kappa \frac{\beta_{q\bar{q}}^3}{m_q^2 \sqrt{\pi}} \vec{S}_q \cdot \vec{S}_{\bar{q}}, \quad \text{for C}, \quad (3.9)$$

where

$$\langle \vec{S}_q \cdot \vec{S}_{\bar{q}} \rangle = \begin{cases} 1/4 & \text{for vector} \\ -3/4 & \text{for pseudoscalar.} \end{cases} \quad (3.10)$$

and

$$\langle M_{q\bar{q}}^2 \rangle = 4 \times \left\{ m_q^2 + \frac{3\beta_{q\bar{q}}^2}{2} + m_q \langle V_{q\bar{q}} \rangle \right\}, \quad (3.11)$$

where we use  $m_{u(d)} = 0.25$  GeV,  $m_s = 0.37$  GeV,  $\beta_{u\bar{u}} = 0.3194$  GeV and  $\beta_{s\bar{s}} = 0.3478$  GeV.

• ~~Methods of fitting three parameters,  $(A_i, B_i, \kappa)$  ( $i = 1, 2, 3$ ):~~

~~(M1) Fitting the  $\pi$ ,  $\rho$  and  $K^*$  meson masses;~~

~~$$M_{Vn\bar{n}}^2 = m_\rho^2, M_{Pn\bar{n}}^2 = m_\pi^2, \text{ and } M_{Vs\bar{s}}^2 = 2m_{K^*}^2 - m_\rho^2.$$~~

~~(M2) Fitting  $\pi$ ,  $\rho$  and  $K$  meson masses:~~

~~$$M_{Vn\bar{n}}^2 = m_\rho^2, M_{Pn\bar{n}}^2 = m_\pi^2, \text{ and } M_{Ps\bar{s}}^2 = 2m_K^2 - m_\pi^2.$$~~

\* Methods of fitting Parameters,  
( $A_i, B_i, K$ )

e.g.)  $V = \underbrace{V_{s-H.O}}_{V_0} + \underbrace{V_{s-s}}_{V_I}$

step 1) Variational Principle

$$\begin{aligned} \langle M^2 \rangle &= \langle M_0^2 \rangle + \langle M_I^2 \rangle \\ &= 4 \left\{ m^2 + \frac{3\beta^2}{2} + mV_0 \right\} + 4mV_I \end{aligned}$$

→ minimization

$$\therefore \frac{\partial \langle M_0^2 \rangle}{\partial \beta} = 0 \sim \beta - m\beta^3 B_1$$

$$\therefore B_1 = \frac{\beta_{u\bar{u}}^4}{m_u} = \frac{\beta_{s\bar{s}}^4}{m_s}$$

check with Jaus's Parameters

$$\frac{\beta_{u\bar{u}}^4}{m_u} = 0.042 \sim \frac{\beta_{s\bar{s}}^4}{m_s} = 0.040$$

→ model parameters seems O.K.



Step 2) Find  $A_1$  and  $K$   
using  $\rho$ - $\pi$  splitting

i.e.

$$\langle M_{n\bar{n}}^2 \rangle_V = m_\rho^2, \quad \langle M_{s\bar{s}}^2 \rangle_P = m_\pi^2$$

\* Numerical Results

$$A_1 = -1.024 \text{ GeV}, \quad B_1 = 0.042 \text{ GeV}^3$$

$$K = 0.488$$

and

$$\lambda_V = 0.6 m_\pi^2 \text{ GeV}^2, \quad X_V = -6.18$$

$$\lambda_P = 23.6 m_\pi^2 \text{ GeV}^2, \quad X_P = 0.68$$

Mass Spectra [MeV]

	Th.		Th.
$\pi$ (135)	135	$\rho$ (770)	770
$K$ (494)	469	$K^*$ (892)	848
$\eta$ (547)	547	$\eta'$ (958)	958
$\omega$ (782)	782	$\phi$ (1020)	1020