

The Running Coupling in a Coupling-Coherent Approach

to

Light-Front Field Theory

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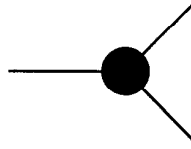
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## GOAL

- Coupling coherence with a smooth transformation
- Example:  $\langle \phi_2 \phi_3 | \mathcal{M}_n^2 | \phi_1 \rangle$  in  $\phi_{5+1}^3$



## OUTLINE

- 1 - Overview of the Approach
- 2 - A General Coupling-Coherent Solution
- 3 - A Second-Order Matrix Element
- 4 - The Third-Order Matrix Element
  - a - The Running Coupling
  - b - Noncanonical Operators
- 5 - Summary

## OVERVIEW OF THE APPROACH

- $\mathcal{M}_f^2$  is the free mass-squared operator

$$\mathcal{M}_f^2 |K\rangle = M_K^2 |K\rangle$$

$$\square_{KJ} \equiv M_K^2 - M_J^2$$

- $\mathcal{M}_n^2$  is the mass-squared operator at scale  $\Lambda_n$ .

$$\langle F | \mathcal{M}_n^2 | I \rangle = M_I^2 \langle F | I \rangle + e^{-\Lambda_n^{-4} \square_{IF}^2} \langle F | \mathcal{M}_{I_n}^2 | I \rangle$$

- Infinite number of scales

$$\Lambda_0 > \Lambda_1 > \Lambda_2 > \cdots > \Lambda_{n-2} > \Lambda_{n-1} > \Lambda_n > \Lambda_{n+1} > \cdots$$

- Unitary Transformation that lowers the cutoff

$$\mathcal{M}_{n+1}^2 = U(\Lambda_{n+1}, \Lambda_n) \mathcal{M}_n^2 U^\dagger(\Lambda_{n+1}, \Lambda_n)$$

$$\frac{dU(\Lambda_{n+1}, \Lambda_n)}{d\Lambda_{n+1}^{-4}} = [\mathcal{M}_f^2, \mathcal{M}_{n+1}^2] U(\Lambda_{n+1}, \Lambda_n)$$

$$U(\Lambda_n, \Lambda_n) \equiv 1$$

F. Wegner, Ann. Physik **3**, 77 (1994)

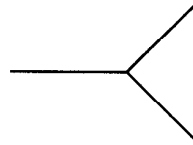
- Perturbative Expansion

$$\begin{aligned}
\langle F | \mathcal{M}_{I(n+1)}^2 | I \rangle &= \langle F | \mathcal{M}_{I_n}^2 | I \rangle \\
&+ \frac{1}{2} \sum_K \langle F | \mathcal{M}_{I_n}^2 | K \rangle \langle K | \mathcal{M}_{I_n}^2 | I \rangle D_1(F, K, I) C_1(n, F, K, I) \\
&+ \frac{1}{4} \sum_{K,L} \langle F | \mathcal{M}_{I_n}^2 | K \rangle \langle K | \mathcal{M}_{I_n}^2 | L \rangle \langle L | \mathcal{M}_{I_n}^2 | I \rangle T(n, F, K, L, I) \\
&+ \dots
\end{aligned}$$

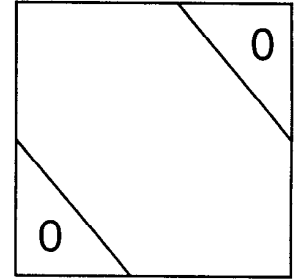
- Postulate an interaction

$$\begin{aligned}
\mathcal{M}_{I_n}^2 &= \hat{\mathcal{P}}^+ (2\pi)^5 g_n \int D_1 D_2 D_3 \left[ a_3^\dagger a_1 a_2 \delta^{(5)}(p_3 - p_1 - p_2) \right. \\
&\quad \left. + a_2^\dagger a_3^\dagger a_1 \delta^{(5)}(p_2 + p_3 - p_1) \right] + \mathcal{O}(g_n^2)
\end{aligned}$$

$$D_i \equiv \frac{d^4 p_{i\perp} dp_i^+}{64\pi^5 p_i^+}$$



$$g_n \equiv \frac{\langle \phi_2 \phi_3 | \mathcal{M}_n^2 | \phi_1 \rangle \Big|_{\vec{p}_{2\perp} = \vec{p}_{3\perp} = 0}}{64\pi^5 p_1^+ \delta^{(5)}(p_1 - p_2 - p_3)} \Big|_{\text{constant part}}$$



- Noncanonical operators have strengths governed by coupling coherence

R. J. Perry, Ann. Phys. **232**, 116 (1994)

# A GENERAL COUPLING-COHERENT SOLUTION

- Recursion Relation

$$\langle F | \mathcal{M}_{I(n+1)}^2 | I \rangle \Big|_{\mathcal{O}(g_n^r)}^{NC} = \langle F | \mathcal{M}_{I_n}^2 | I \rangle \Big|_{\mathcal{O}(g_n^r)}^{NC} + \langle F | \Delta_n | I \rangle \Big|_{\mathcal{O}(g_n^r)}^{NC}$$

- All interactions are approximately transverse-local

$$\begin{aligned} \langle F | \mathcal{M}_{I_n}^2 | I \rangle \Big|_{\mathcal{O}(g_n^r)}^{NC} &= 64\pi^5 \mathcal{P}_I^+ \delta^{(5)}(\mathcal{P}_I - \mathcal{P}_F) \left[ \prod_{i=1}^{N_s} 64\pi^5 s_i^+ \delta^{(5)}(s_i - s'_i) \right] \\ &\times \Lambda_n^{6-2N_{int}} \sum_{\{m_{pc}\}} C_{\{m_{pc}\}}^{(n)} h_{\{m_{pc}\}} \left( \left\{ \frac{k_p^+}{\mathcal{P}_I^+} \right\} \right) \prod_{p=1}^{N_{int}} \prod_{c=1}^4 \left( \frac{k_{p\perp}^c}{\Lambda_n} \right)^{m_{pc}} \end{aligned}$$

- $\langle F | \Delta_n | I \rangle \Big|_{\mathcal{O}(g_n^r)}^{NC}$  is restricted to have the form

$$\begin{aligned} \langle F | \Delta_n | I \rangle \Big|_{\mathcal{O}(g_n^r)}^{NC} &= 64\pi^5 \mathcal{P}_I^+ \delta^{(5)}(\mathcal{P}_I - \mathcal{P}_F) \left[ \prod_{i=1}^{N_s} 64\pi^5 s_i^+ \delta^{(5)}(s_i - s'_i) \right] \\ &\times g_n^r \sum_{\{m_{pc}\}} B_{\{m_{pc}\}} h_{\{m_{pc}\}} \left( \left\{ \frac{k_p^+}{\mathcal{P}_I^+} \right\} \right) \\ &\times \left[ \Lambda_{n+1}^{6-2N_{int}} \prod_{p,c} \left( \frac{k_{p\perp}^c}{\Lambda_{n+1}} \right)^{m_{pc}} - \Lambda_n^{6-2N_{int}} \prod_{p,c} \left( \frac{k_{p\perp}^c}{\Lambda_n} \right)^{m_{pc}} \right] \end{aligned}$$

- Match powers of transverse momentum
- Assume Coupling Coherence

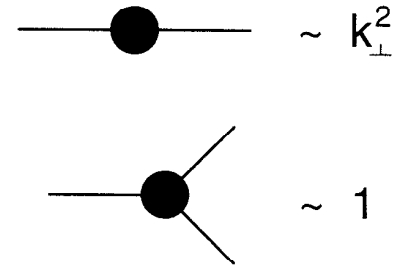
$$C_{\{m_{pc}\}}^{(n)} = \alpha_{\{m_{pc}\}} g_n^r + \mathcal{O}(g_n^{r+1})$$

- This leads to

$$\left( \alpha_{\{m_{pc}\}} - B_{\{m_{pc}\}} \right) \left[ \Lambda_{n+1}^{6-2N_{int}} \prod_{p,c} \left( \frac{k_{p\perp}^c}{\Lambda_{n+1}} \right)^{m_{pc}} - \Lambda_n^{6-2N_{int}} \prod_{p,c} \left( \frac{k_{p\perp}^c}{\Lambda_n} \right)^{m_{pc}} \right] = 0$$

- Marginal contributions to  $\langle F | \Delta_n | I \rangle \Big|_{\mathcal{O}(g_n^r)}^{NC}$  have

$$6 - 2N_{int} - \sum_{p,c} m_{pc} = 0$$



- Nonmarginal contributions to  $\langle F | \Delta_n | I \rangle \Big|_{\mathcal{O}(g_n^r)}^{NC}$  have

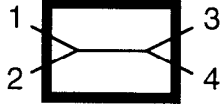
$$\alpha_{\{m_{pc}\}} = B_{\{m_{pc}\}}$$

- General Coupling-Coherent solution for nonmarginal contributions:

$$\langle F | \mathcal{M}_{I_n}^2 | I \rangle \Big|_{\mathcal{O}(g_n^r), \text{nonmarginal}} =$$

$$- \langle F | \Delta_n | I \rangle \Big|_{\mathcal{O}(g_n^r), \text{nonmarginal}, \Lambda_{n+1} \text{ terms} \rightarrow 0}$$

## A SECOND-ORDER MATRIX ELEMENT

- As an example, compute 

- Use jacobi variables ( $\mathcal{P} \equiv p_1 + p_2$ )

$$p_1 = (x\mathcal{P}^+, x\vec{\mathcal{P}}_\perp + \vec{r})$$

$$p_2 = ([1-x]\mathcal{P}^+, [1-x]\vec{\mathcal{P}}_\perp - \vec{r})$$

$$p_3 = (y\mathcal{P}^+, y\vec{\mathcal{P}}_\perp + \vec{q})$$

$$p_4 = ([1-y]\mathcal{P}^+, [1-y]\vec{\mathcal{P}}_\perp - \vec{q})$$

- Define initial and final mass squares

$$u \equiv \frac{r^2}{x(1-x)} \qquad v \equiv \frac{q^2}{y(1-y)}$$

$$\begin{aligned} & \langle \phi_3 \phi_4 | \Delta_n | \phi_1 \phi_2 \rangle \Big|_{\mathcal{O}(g_n^2), \text{ann}} = 64\pi^5 \delta^{(5)}(p_1 + p_2 - p_3 - p_4) \\ & \times \frac{1}{2} g_n^2 \mathcal{P}^+ \left( \frac{1}{u} + \frac{1}{v} \right) \left( e^{-2\Lambda_n^{-4} uv} - e^{-2\Lambda_{n+1}^{-4} uv} \right) \end{aligned}$$



$$\langle \phi_3 \phi_4 | \Delta_n | \phi_1 \phi_2 \rangle \Big|_{\mathcal{O}(g_n^2), ann} = 64\pi^5 \delta^{(5)}(p_1 + p_2 - p_3 - p_4) \\ \times \frac{1}{2} g_n^2 \mathcal{P}^+ \left( \frac{1}{u} + \frac{1}{v} \right) \left( e^{-2\Lambda_n^{-4} uv} - e^{-2\Lambda_{n+1}^{-4} uv} \right)$$

- Apply the general solution

$$\langle F | \mathcal{M}_{I_n}^2 | I \rangle \Big|_{\mathcal{O}(g_n^r), nonmarginal} = \\ - \langle F | \Delta_n | I \rangle \Big|_{\mathcal{O}(g_n^r), nonmarginal, \Lambda_{n+1} terms \rightarrow 0}$$

- Leads to the specific solution

$$\langle \phi_3 \phi_4 | \mathcal{M}_{I_n}^2 | \phi_1 \phi_2 \rangle \Big|_{\mathcal{O}(g_n^2), ann} = \begin{array}{c} 1 \\ \hline \square \\ \hline 2 \end{array} \begin{array}{c} 3 \\ \hline \square \\ \hline 4 \end{array} \\ = 64\pi^5 \delta^{(5)}(p_1 + p_2 - p_3 - p_4) \frac{1}{2} g_n^2 \mathcal{P}^+ \left( \frac{1}{u} + \frac{1}{v} \right) \left( 1 - e^{-2\Lambda_n^{-4} uv} \right)$$

## THE THIRD-ORDER MATRIX ELEMENT

- GOAL: compute  $\langle \phi_2 \phi_3 | \mathcal{M}_n^2 | \phi_1 \rangle$  through  $\mathcal{O}(g_n^3)$  ~~scribble~~
- The recursion relation:

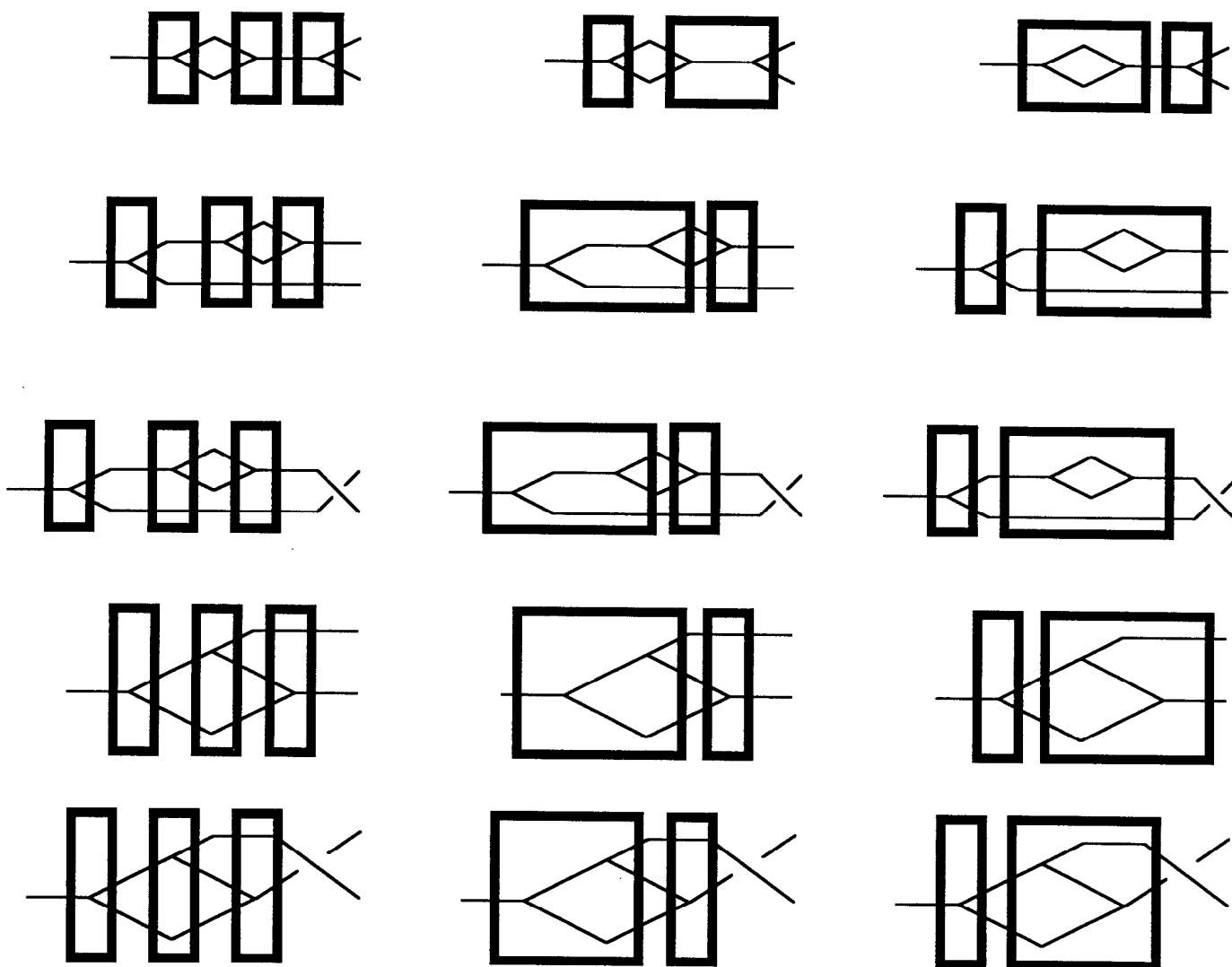
$$\langle \phi_2 \phi_3 | \mathcal{M}_{I(n+1)}^2 | \phi_1 \rangle = \langle \phi_2 \phi_3 | \mathcal{M}_{I_n}^2 | \phi_1 \rangle + \langle \phi_2 \phi_3 | \Delta_n | \phi_1 \rangle$$

$$\begin{aligned} \langle \phi_2 \phi_3 | \Delta_n | \phi_1 \rangle &= \frac{1}{2} \sum_K \langle \phi_2 \phi_3 | \mathcal{M}_{I_n}^2 | K \rangle \langle K | \mathcal{M}_{I_n}^2 | \phi_1 \rangle D_1() C_1() \\ &+ \frac{1}{4} \sum_{K,L} \langle \phi_2 \phi_3 | \mathcal{M}_{I_n}^2 | K \rangle \langle K | \mathcal{M}_{I_n}^2 | L \rangle \langle L | \mathcal{M}_{I_n}^2 | \phi_1 \rangle T() \end{aligned}$$

# Third-Order Contributions

to

$$\langle \varphi_2 \varphi_3 | \Delta_n | \varphi_1 \rangle$$



- The recursion relation:

$$\langle \phi_2 \phi_3 | \mathcal{M}_{I(n+1)}^2 | \phi_1 \rangle = \langle \phi_2 \phi_3 | \mathcal{M}_{I_n}^2 | \phi_1 \rangle + \langle \phi_2 \phi_3 | \Delta_n | \phi_1 \rangle$$

- Examine the case  $\vec{p}_{1\perp} = \vec{p}_{2\perp} = 0$

$$\begin{aligned} \langle \phi_2 \phi_3 | \mathcal{M}_{I(n+1)}^2 | \phi_1 \rangle \Big|_{\vec{p}_{1\perp} = \vec{p}_{2\perp} = 0} &= \langle \phi_2 \phi_3 | \mathcal{M}_{I_n}^2 | \phi_1 \rangle \Big|_{\vec{p}_{1\perp} = \vec{p}_{2\perp} = 0} \\ &+ \frac{3\pi^2}{8} g_n^3 p_1^+ \delta^{(5)}(p_1 - p_2 - p_3) \log \frac{\Lambda_n^2}{\Lambda_{n+1}^2} \end{aligned}$$

- Use the definition of  $g_n$

$$g_n \equiv \frac{\langle \phi_2 \phi_3 | \mathcal{M}_n^2 | \phi_1 \rangle \Big|_{\vec{p}_{2\perp} = \vec{p}_{3\perp} = 0}}{64\pi^5 p_1^+ \delta^{(5)}(p_1 - p_2 - p_3)} \Big|_{\text{constant part}}$$

- The Running Coupling

$$g_{n+1} = g_n + \frac{3}{512\pi^3} g_n^3 \log \frac{\Lambda_n^2}{\Lambda_{n+1}^2}$$

- Now determine the noncanonical operators in  $\langle \phi_2 \phi_3 | \mathcal{M}_n^2 | \phi_1 \rangle$

$$\begin{aligned} \langle \phi_2 \phi_3 | \Delta_n | \phi_1 \rangle \Big|_{1st\ 9\ diagrams} &= \frac{\pi^2}{48} g_n^3 p_1^+ \delta^{(5)}(p_1 - p_2 - p_3) \\ &\times \int_0^\infty du [f(u, v) + 2f(u, -v)] \end{aligned}$$

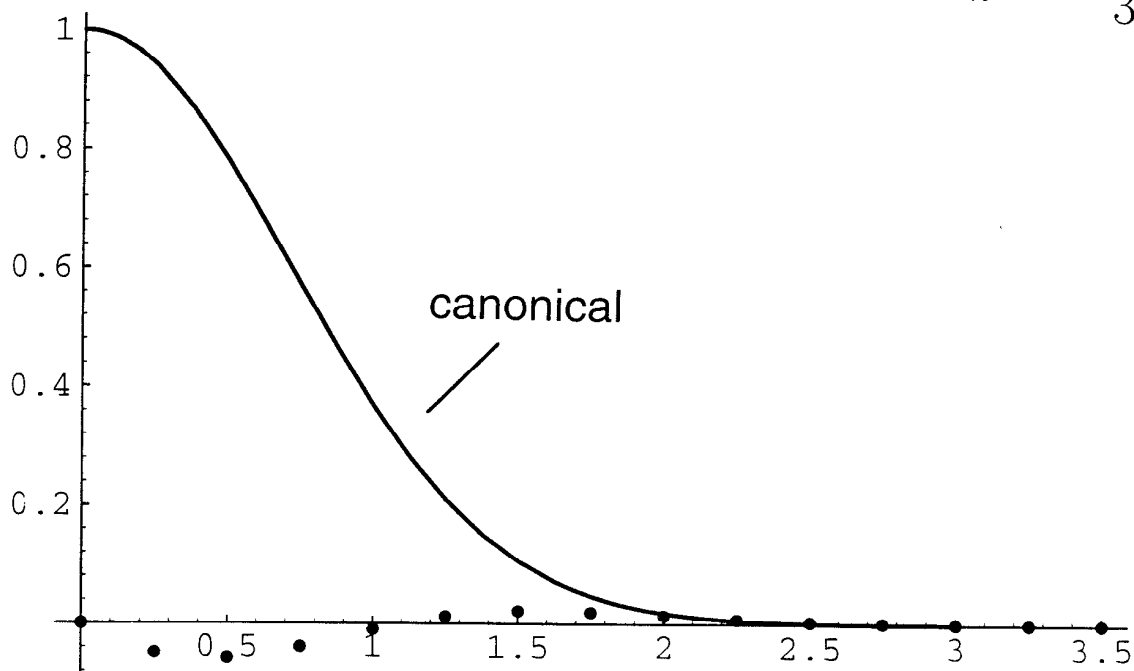
$$\begin{aligned} f(u, v) &= \left[ e^{-2\Lambda_n^{-4}u^2} - e^{-2\Lambda_{n+1}^{-4}u^2} \right] \left[ \frac{1}{u} + \frac{v}{u^2} + \frac{2}{v} \right] \\ + \left[ e^{-2\Lambda_n^{-4}(u^2-uv)} - e^{-2\Lambda_{n+1}^{-4}(u^2-uv)} \right] &\left[ \frac{1}{v-u} - \frac{1}{u} + \frac{u}{v(v-u)} - \frac{1}{v} \right] \end{aligned}$$

- Apply the general solution

$$\begin{aligned} \langle \phi_2 \phi_3 | \mathcal{M}_{In}^2 | \phi_1 \rangle \Big|_{NM, 1st\ 9} &= - \langle \phi_2 \phi_3 | \Delta_n | \phi_1 \rangle \Big|_{NM, 1st\ 9, \Lambda_{n+1}\ terms \rightarrow 0} \\ &= - \frac{\pi^2}{48} g_n^3 p_1^+ \delta^{(5)}(p_1 - p_2 - p_3) \\ &\times \int_0^\infty du [f(u, v) + 2f(u, -v) - 3f(u, 0)] \Big|_{\Lambda_{n+1}^{-2} \rightarrow 0} \end{aligned}$$

$$\langle \varphi_2 \varphi_3 | \mathcal{M}_n^2 | \varphi_1 \rangle$$

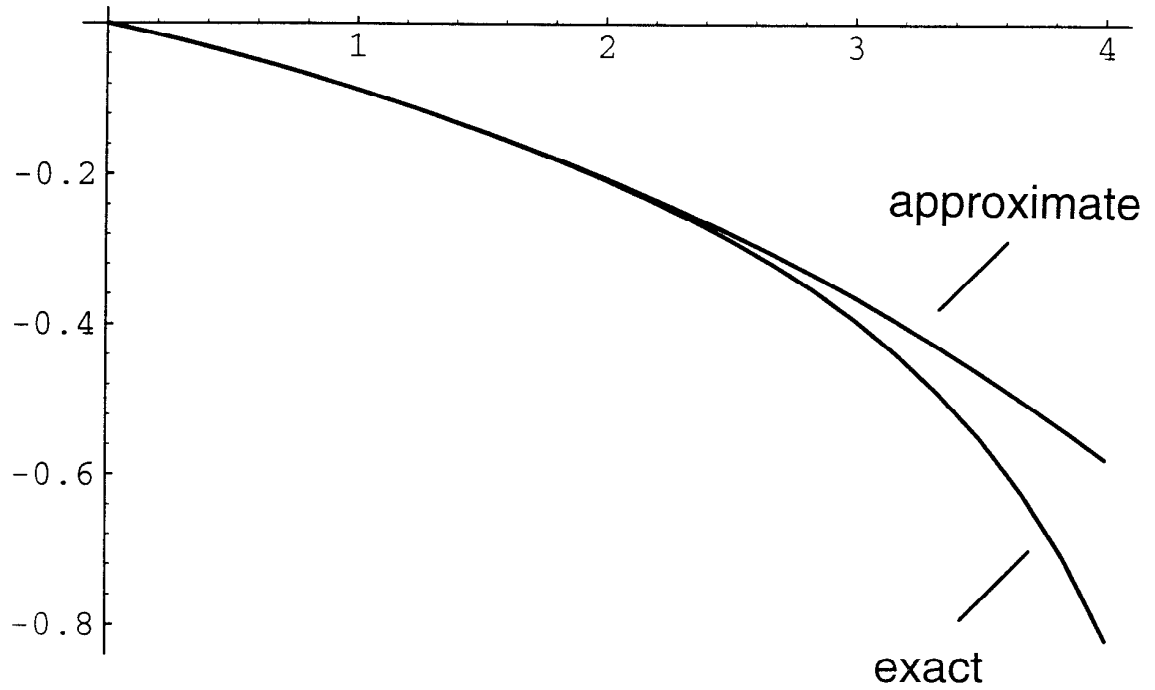
$$g_n^2 = \frac{512\pi^3}{3}$$



$$\frac{v}{\Lambda_n^2}$$

Noncanonical  $\langle \varphi_2 \varphi_3 | \mathcal{M}_{In}^2 | \varphi_1 \rangle$

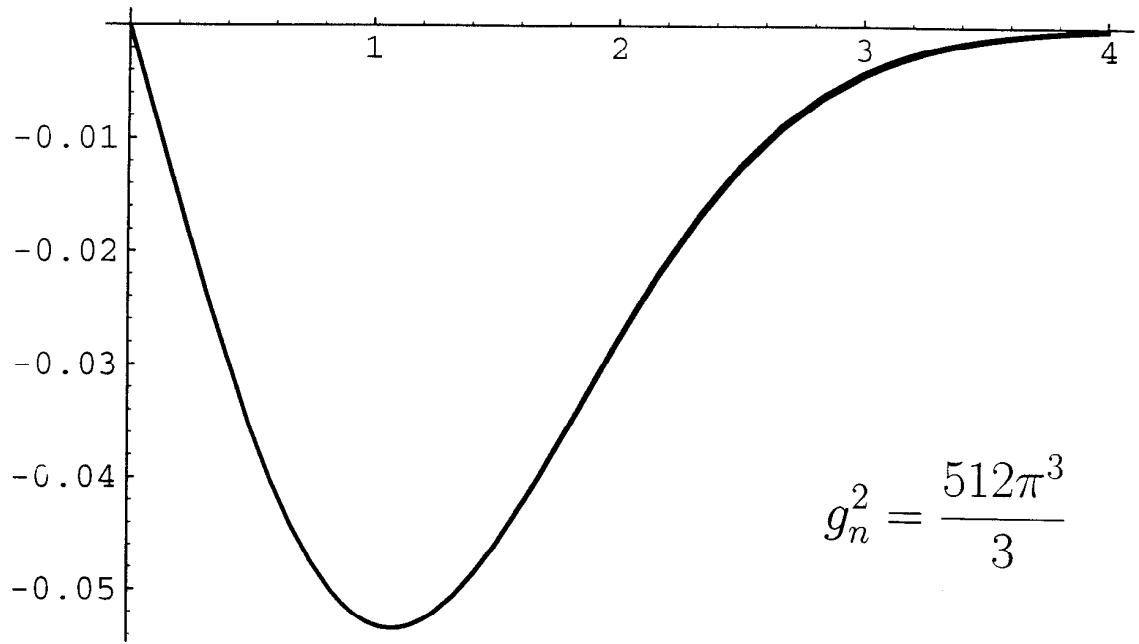
$$\frac{v}{\Lambda_n^2}$$



$$g_n^2 = \frac{512\pi^3}{3}$$

Noncanonical  $\langle \varphi_2 \varphi_3 | \mathcal{M}_n^2 | \varphi_1 \rangle$

$$\frac{v}{\Lambda_n^2}$$



$$g_n^2 = \frac{512\pi^3}{3}$$



## SUMMARY

- 1 - Developed a general method for coupling coherence with a smooth transformation
- 2 - Computed second- and third-order matrix elements of  $\mathcal{M}_n^2$  in  $\phi_{5+1}^3$  theory
- 3 - Showed that the nonmarginal noncanonical contributions to  $\langle \phi_2 \phi_3 | \mathcal{M}_n^2 | \phi_1 \rangle$  are small
- 4 - Verified the assumption of approximate transverse locality

# NONCANONICAL MARGINAL OPERATORS

- The general equation for the coupling is

$$h_{n+1} = h_n + ag_n^3 + bg_n^2 h_n + cg_n^5$$

- Use coupling coherence

$$h_n = \alpha g_n^3 + \beta g_n^5 + \mathcal{O}(g_n^7)$$

- Use the running coupling

$$g_{n+1} = g_n + dg_n^3 + \mathcal{O}(g_n^5)$$

- To third order

$$a = 0$$

- To fifth order

$$\alpha = \frac{c}{3d - b}$$