Gravitational formfactors, equivalence principle and shear viscosity



3D Structure of the Nucleon via Generalized Parton Distributions

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Oleg Teryaev JINR, Dubna

Outline

Gravitational formfactors and hadron spin structure gravitational field

Spacelike vs timelike

Viscosity

Space -> Time: T-odd-> exotics

Golographic bounds -> smallness of DA



Gravitational Formfactors (spin 1/2)

$$\langle p'|T_{q,g}^{\mu\nu}|p\rangle = \bar{u}(p')\Big[A_{q,g}(\Delta^2)\gamma^{(\mu}p^{\nu)} + B_{q,g}(\Delta^2)P^{(\mu}i\sigma^{\nu)\alpha}\Delta_{\alpha}/2M]u(p)$$

• Conservation laws - zero Anomalous Gravitomagnetic Moment : $\mu_G = J$ (g=2)

$$\begin{split} P_{q,g} &= A_{q,g}(0) & A_{q}(0) + A_{g}(0) = 1 \\ J_{q,g} &= \frac{1}{2} \left[A_{q,g}(0) + B_{q,g}(0) \right] & A_{q}(0) + B_{q}(0) + A_{g}(0) + B_{g}(0) = 1 \end{split}$$

- May be extracted from high-energy experiments/NPQCD calculations
- Describe the partition of angular momentum between quarks and gluons
- Describe interaction with scalar and tensor particles, with both classical and "TeV" gravity

Gravity and hadron structure: (OT'99)

Interaction – field vs metric deviation

$$M = \langle P'|J^{\mu}_{q}|P\rangle A_{\mu}(q)$$

 $M = \frac{1}{2} \sum_{q,G} \langle P' | T_{q,G}^{\mu\nu} | P \rangle h_{\mu\nu}(q)$

Static limit

$$\langle P|J_q^{\mu}|P\rangle = 2e_q P^{\mu}$$

$$\sum_{q,G} \langle P|T_i^{\mu\nu}|P\rangle = 2P^{\mu}P^{\nu}$$
$$h_{00} = 2\phi(x)$$

$$M_0 = \langle P|J_q^{\mu}|P\rangle A_{\mu} = 2e_q M\phi(q)$$

$$M_0 = \frac{1}{2} \sum_{q,G} \langle P | T_i^{\mu\nu} | P \rangle h_{\mu\nu} = 2M \cdot M\phi(q)$$

Mass as charge – equivalence principle

EP and hadron structure

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"Microscopic" EP (coupling of gravity to EMT)
Conservation law
 q(x) + G(x) = 1
"Macroscopic" EP (universal falling):
Tested VERY precisely
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Gravitomagnetism

• Gravitomagnetic field (weak, except in gravity waves) – action on spin from $M = \frac{1}{2} \sum_{G} \langle P' | T_{q,G}^{\mu\nu} | P \rangle h_{\mu\nu}(q)$

$$ec{H}_J = rac{1}{2} rot ec{g}; \; ec{g}_i \equiv g_{0i}$$
 spin dragging twice smaller than EM

• Lorentz force – similar to EM case: factor $\frac{1}{2}$ cancelled with 2 from $h_{00} = 2\phi(x)$ Larmor frequency same as EM

$$\omega_J = \frac{\mu_G}{J} H_J = \frac{H_L}{2} = \omega_L \ \vec{H}_L = rot \vec{g}$$

 Orbital and Spin momenta dragging – the same -Equivalence principle



Experimental test of PNEP

Reinterpretation of the data on G(EDM) search

VOLUME 68 13 JANUARY 1992 NUMBER 2

Search for a Coupling of the Earth's Gravitational Field to Nuclear Spins in Atomic Mercury

B. J. Venema, P. K. Majumder, S. K. Lamoreaux, B. R. Heckel, and E. N. Fortson Physics Department, FM-15, University of Washington, Seattle, Washington 98195 (Received 25 Sentember 1991)

If (CP-odd!) GEDM=0 -> constraint for AGM (Silenko, OT'07) from Earth rotation — was considered as obvious (but it is just EP!) background

$$\mathcal{H} = -g\mu_N \mathbf{B} \cdot \mathbf{S} - \zeta \hbar \boldsymbol{\omega} \cdot \mathbf{S}, \quad \zeta = 1 + \chi$$

$$|\chi(^{201}\text{Hg}) + 0.369\chi(^{199}\text{Hg})| < 0.042 \quad (95\%\text{C.L.})$$



If spin is just a (pseudo) vector: EP due to Earth rotation is trivial

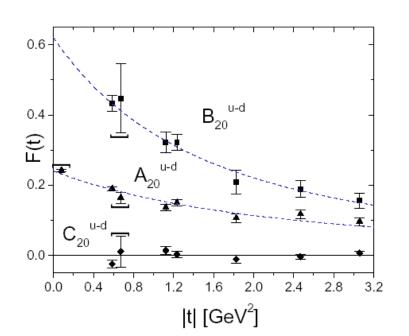
Crucial if measured by a device in rotating frame

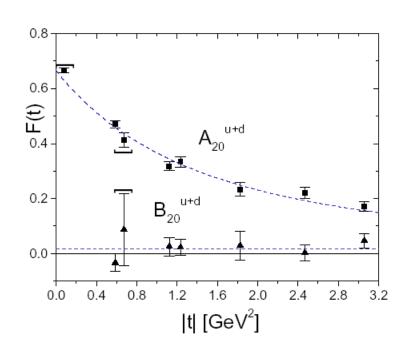
Quantum measurement problem becomes practical

Cf Unruh effect in HIC (Prokhorov, OT, Zakharov'19 and in progress)

Generalization of Equivalence principle (smallness of B_G used om gluons FF extraxtion- talk of Z.-E. Meziani)

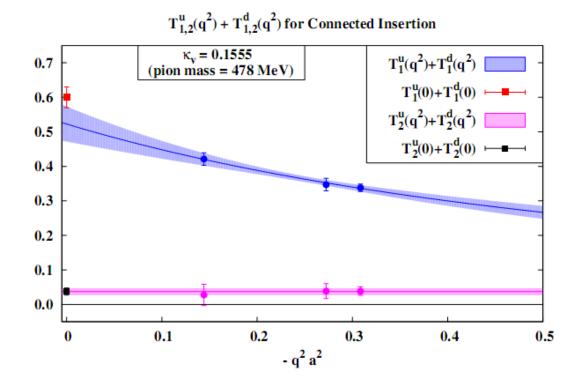
 Various arguments: AGM ≈0 separately for quarks and gluons – most clear from the lattice (LHPC/SESAM)





Further lattice study (M. Deka et al. Phys.Rev. D91 (2015) no.1, 014505)

 Sum of u and d for Dirac (T1) and Pauli (T2) FFs



Extended Equivalence Principle=Exact EquiPartition

- In QED, pQCD violated (Brodsky et al)
- Reason in the case of ExEP enforcing by subtracted DR - no smooth transition for zero fermion mass limit (Milton, 73)
- Conjecture (O.T., 2001 prior to lattice data)

 valid in NP QCD zero quark mass limit is safe due to chiral symmetry breaking (test in models? Cf talk of H.- Ch. Kim)
- Gravityproof confinement? Nucleons do not break even by black holes?
- Support by recent observations of smalness of (nucleon "cosmological constant") Cbar

One more gravitational formfactor (related to "D-term" of Maxim Polyakov and Christian Weiss)

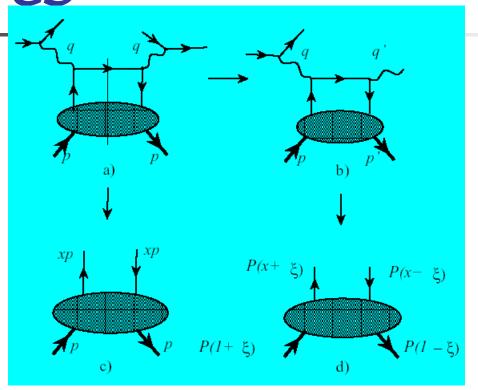
Quadrupole

$$\langle P + q/2 | T^{\mu\nu} | P - q/2 \rangle = C(q^2)(g^{\mu\nu}q^2 - q^{\mu}q^{\nu}) + \dots$$

• Cf vacuum matrix element – cosmological constant $\langle 0|T^{\mu\nu}|0\rangle = \Lambda g^{\mu\nu}$ $\Lambda = C(q^2)q^2$

- NO "vacuum-like" term EP, Smallness
 -ExEP
- How to measure experimentally DVCS (and DVMP?)

QCD Factorization for DIS and



Manifestly spectral

$$\mathcal{H}(x_B) = \int_{-1}^1 dx \frac{H(x)}{x - x_B + i\epsilon}$$

Extra dependence

on
$$\xi$$

$$\mathcal{H}(\xi) = \int_{-1}^{1} dx \frac{H(x,\xi)}{x - \xi + i\epsilon},$$

Unphysical regions

■ DIS : Analytical function – polynomial in $1/x_B$ if $1 \le |X_B|$

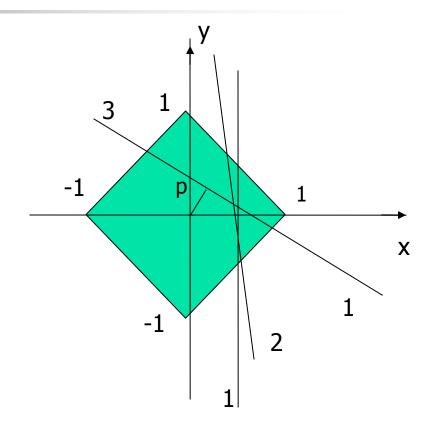
$$H(x_B) = -\int_{-1}^{1} dx \sum_{n=0}^{\infty} H(x) \frac{x^n}{x_B^{n+1}}$$

- DVCS additional problem of analytical continuation of H(x,ξ)
- Solved by using of Double Distributions Radon transform

$$H(z,\xi) = \int_{-1}^{1} dx \int_{|x|-1}^{1-|x|} dy (F(x,y) + \xi G(x,y)) \delta(z - x - \xi y)$$

Double distributions and their integration

- Slope of the integration lineskewness
- Kinematics of DIS: $\xi = 0$ ("forward") vertical line (1)
- Kinematics of DVCS: ξ<1
 line 2
- Line 3: $\xi > 1$ unphysical region required to restore DD by inverse Radon transform: tomography

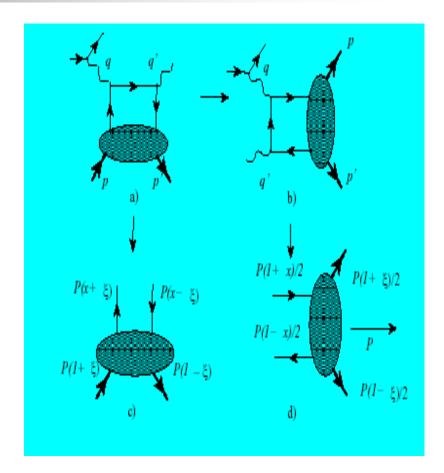


$$f(x,y) = -\frac{1}{2\pi^2} \int_0^\infty \frac{dp}{p^2} \int_0^{2\pi} d\phi |\cos\phi| (H(p/\cos\phi + x + ytg\phi, tg\phi) - H(x + ytg\phi, tg\phi)) =$$

$$= -\frac{1}{2\pi^2} \int_{-\infty}^\infty \frac{dz}{z^2} \int_{-\infty}^\infty d\xi (H(z + x + y\xi, \xi) - H(x + y\xi, \xi))$$

Crossing for DVCS and GPD

- DVCS -> hadron pair production in the collisions of real and virtual photons
- GPD -> GeneralizedDistribution Amplitudes
- Duality between s and t channels (Polyakov,Shuvaev, Guzey, Vanderhaeghen)



GDA -> back to unphysical regions for DIS and DVCS

Recall DIS

$$H(x_B) = -\int_{-1}^{1} dx \sum_{n=0}^{\infty} H(x) \frac{x^n}{x_B^{n+1}}$$

Non-positive powers of X_B

DVCS

$$H(x_B) = -\int_{-1}^{1} dx \sum_{n=0}^{\infty} H(x) \frac{x^n}{x_B^{n+1}} \qquad H(\xi) = -\int_{-1}^{1} dx \sum_{n=0}^{\infty} H(x, \xi) \frac{x^n}{\xi^{n+1}}$$

- Polynomiality (general property of Radon transforms): moments integrals in x weighted with x^n - are polynomials in $1/\xi$ of power n+1
- As a result, analyticity is preserved: only non-positive powers of ξ appear



Holographic property (OT'05)

Factorization Formula

$$\mathcal{H}(\xi) = \int_{-1}^{1} dx \frac{H(x,\xi)}{x - \xi + i\epsilon}$$

$$\mathcal{H}(\xi) = \int_{-1}^{1} dx \frac{H(x,x)}{x - \xi + i\epsilon}$$

$$\Delta \mathcal{H}(\xi) \equiv \int_{-1}^{1} dx \frac{H(x,x) - H(x,\xi)}{x - \xi + i\epsilon}$$

$$= \sum_{n=1}^{\infty} \frac{1}{n!} \frac{\partial^n}{\partial \xi^n} \int_{-1}^1 H(x,\xi) dx (x-\xi)^{n-1} = const$$

Holographic property - II

Directly follows from double distributions

$$H(z,\xi) = \int_{-1}^{1} dx \int_{|x|-1}^{1-|x|} dy (F(x,y) + \xi G(x,y)) \delta(z - x - \xi y)$$

 Constant is the SUBTRACTION one - due to the (generalized) Polyakov-Weiss term G(x,y)

$$\Delta \mathcal{H}(\xi) = \int_{-1}^{1} dx \int_{|x|-1}^{1-|x|} dy \frac{G(x,y)}{1-y}$$

$$= \int_{-\xi}^{\xi} dx \frac{D(x/\xi)}{x - \xi + i\epsilon} = \int_{-1}^{1} dz \frac{D(z)}{z - 1} = const$$

Holographic property - III

- 2-dimensional space -> 1-dimensional section!
- Momentum space: any relation to holography in coordinate space ?!

• Strategy (now adopted) of GPD's studies: start at diagonals (through SSA due to imaginary part of DVCS $x=-\xi$ amplitude) and restore by making use of dispersion relations + subtraction constants

 $x = \mathcal{E}$

Analyticity of Compton amplitudes in energy plane (Anikin, OT'07)

Finite subtraction implied

$$\operatorname{Re}\mathcal{A}(\nu,Q^{2}) = \frac{\nu^{2}}{\pi} \mathcal{P} \int_{\nu_{0}}^{\infty} \frac{d\nu'^{2}}{\nu'^{2}} \frac{\operatorname{Im}\mathcal{A}(\nu',Q^{2})}{(\nu'^{2}-\nu^{2})} + \Delta \qquad \Delta = 2 \int_{-1}^{1} d\beta \frac{D(\beta)}{\beta-1}$$

$$\Delta_{\operatorname{CQM}}^{p}(2) = \Delta_{\operatorname{CQM}}^{n}(2) \approx 4.4, \qquad \Delta_{\operatorname{latt}}^{p} \approx \Delta_{\operatorname{latt}}^{n} \approx 1.1$$

- Numerically close to Thomson term for real proton (but NOT neutron) Compton Scattering!
- Duality (sum of squares vs square of sum; proton: 4/9+4/9+1/9=1)?!

From D-term to pressure

- Inverse -> 1st moment (model)
- Kinematical factor moment of pressure $D \sim -\langle p r^4 \rangle$ $(\langle p r^2 \rangle = 0)$ M.Polyakov (2003)

$$T_{\mu\nu}^{Q}(\vec{r},\vec{s}) = \frac{1}{2E} \int \frac{d^{3}\Delta}{(2\pi)^{3}} e^{i\vec{r}\cdot\vec{\Delta}} \langle p', S'|\hat{T}_{\mu\nu}^{Q}(0)|p, S\rangle$$

$$T_{ij}(\vec{r}) = s(r) \left(\frac{r_i r_j}{r^2} - \frac{1}{3} \delta_{ij} \right) + p(r) \delta_{ij}$$

- Possible justification: Born gravitational scattering
- Stable equilibrium D<0: Holds for quarks (or leptons)
 in photon

Pressure in hadron pairs production

- Back to GDA region
- -> moments of H(x,x) define the coefficients
 of powers of cosine!- 1/ξ
- Higher powers of cosine in t-channel – threshold in s -channel
- Larger for pion than for nucleon pairs because of less fast decrease at x ->1
- $\begin{array}{ccc} \bullet & \mathsf{Large} & \xi & \mathsf{limit} \mathsf{access} \\ \mathsf{to} & \mathsf{D}\text{-term} \end{array}$

$$\mathcal{H}(\xi) = -\int_{-1/\xi}^{1/\xi} dx \sum_{n=0}^{\infty} H(x,\xi) \frac{x^n}{\xi^{n+1}}$$
$$= -\int_{-1/\xi}^{1/\xi} dx \sum_{n=0}^{\infty} H(x,x) \frac{x^n}{\xi^{n+1}} + \Delta \mathcal{H}.$$



Gravitational FFs from Belle data on

S. Kumano, Qin-Tao Song and O. Teryaev, PRD 97 (2018) 014020.

Gravitational FFs are related to twist-2

$$A_{++} = \sum_{q} \frac{e^{2}}{2} \int_{0}^{1} dz \frac{2z - 1}{z(1 - z)} \Phi^{q}(z, \xi, W^{2})$$

 $\begin{aligned} \mathsf{GDAS} \qquad & A_{\lambda_{1}\lambda_{2}} = T_{\mu\nu}\varepsilon^{\mu}(\lambda_{1})\varepsilon^{\nu}(\lambda_{2})/e^{2} \\ A_{++} &= \sum_{q} \frac{e^{2}_{q}}{2} \int_{0}^{1} dz \frac{2z-1}{z(1-z)} \Phi^{q}\left(z,\xi,W^{2}\right) \\ & \left\langle \pi^{0}(p_{1})\pi^{0}(p_{2}) \middle| T^{\mu\nu}(0) \middle| 0 \right\rangle = \frac{1}{2} \left[\left(sg^{\mu\nu} - P^{\mu}P^{\nu} \right) \Theta_{1} + \Delta^{\mu}\Delta^{\nu}\Theta_{2} \right] \end{aligned}$

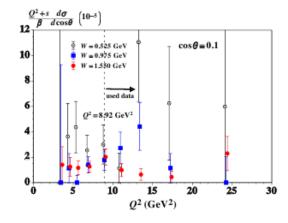
M. Masuda et al. [Belle Collaboration], PRD 93 (2016), 032003.

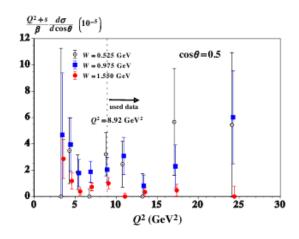
$$P=p_{_1}+p_{_2}$$
 , $\Delta=p_{_1}-p_{_2}$

Belle data and scaling: W=0.525,0.975,

1.55 GeV

$$\frac{(Q^2+s)d\sigma}{\beta d|\cos\theta|} \propto \left|\Phi^{s^0s^0}(z,\cos\theta,W,Q)\right|^2$$





Phase shifts and resonances

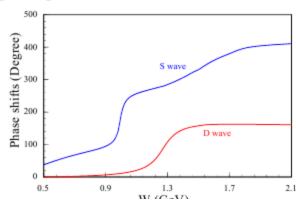
Leading harmonics

$$\sum_{q} \Phi_{q}^{+}(z,\xi,W^{2}) = 18n_{f}z(1-z)(2z-1)[B_{10}(W) + B_{12}(W)P_{2}(2\xi-1)]$$
$$= 18n_{f}z(1-z)(2z-1)[\tilde{B}_{10}(W) + \tilde{B}_{12}(W)P_{2}(\cos\theta)]$$

$$\tilde{B}_{10}(W) = \bar{B}_{10}(W)e^{i\delta_0}$$
, $\tilde{B}_{12}(W) = \bar{B}_{12}(W)e^{i\delta_2}$

S/D shifts

 $f_0(500)$, $f_2(1270)$ contributions



$$\overline{B}_{12}(W) = \beta^2 \frac{10 g_{f_2 \pi \pi} f_{f_2} M_{f_2}^2}{9 \sqrt{2} \sqrt{(M_{f_2}^2 - W^2)^2 - \Gamma_{f_2}^2 M_{f_2}^2}}$$

$$\overline{B}_{10}(W) = \frac{5 g_{f_0 \pi \pi} f_{f_0}}{3 \sqrt{2} \sqrt{(M_{f_0}^2 - W^2)^2 - \Gamma_{f_0}^2 M_{f_0}^2}}$$



Collection

$$\Phi_{q}^{+}(z,\xi,W^{2}) = N_{h}z^{\alpha}(1-z)^{\alpha}(2z-1)[\tilde{B}_{10}(W) + \tilde{B}_{12}(W)P_{2}(\cos\theta)]$$

$$\tilde{B}_{10}(W) = \left[\frac{-3+\beta^{2}}{2}\frac{5R_{\pi}}{9}F_{h}(W^{2}) + \frac{5g_{f_{0}\pi\pi}f_{f_{0}}}{3\sqrt{2}\sqrt{(M_{f_{0}}^{2} - W^{2})^{2} - \Gamma_{f_{0}}^{2}M_{f_{0}}^{2}}}\right]e^{i\delta_{0}}$$

$$\tilde{B}_{12}(W) = \left[\beta^{2}\frac{5R_{\pi}}{9}F_{h}(W^{2}) + \beta^{2}\frac{10g_{f_{2}\pi\pi}f_{f_{2}}M_{f_{2}}^{2}}{9\sqrt{2}\sqrt{(M_{f_{2}}^{2} - W^{2})^{2} - \Gamma_{f_{2}}^{2}M_{f_{2}}^{2}}}\right]e^{i\delta_{2}}$$

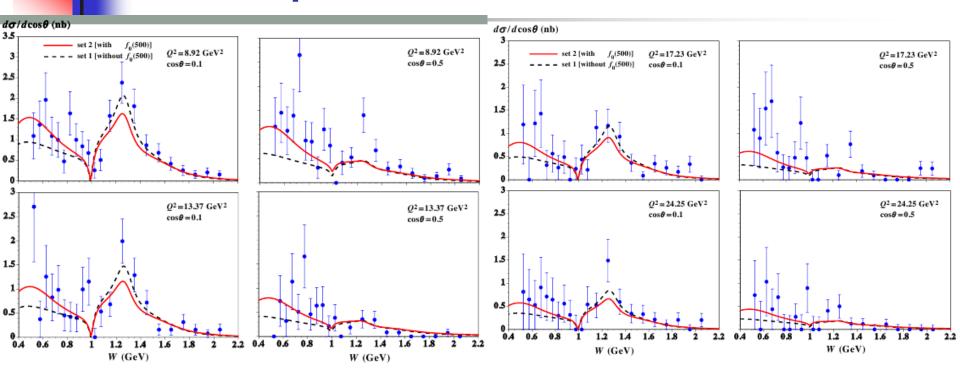
$$F_{h}(W^{2}) = \frac{1}{\left[1 + \frac{W^{2} - 4m_{\pi}^{2}}{\Lambda^{2}}\right]^{n-1}}$$

Best fit with (2) and without (1) f_0

	Set 1	Set 2
α	0.801±0.042	1.157± 0.132
٨	1.602±0.109	1.928±0.213
а	3.878± 0.165	3.800± 0.170
b	0.382± 0.040	0.407± 0.041
f _{f0}		0.0184± 0.034
	2	2

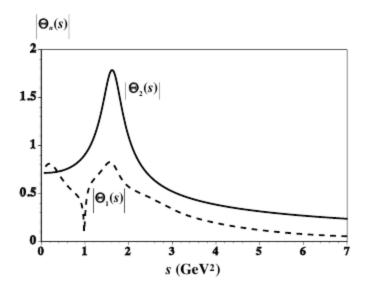
$$\frac{\chi^2}{NOF} = 1.22 \qquad \frac{\chi^2}{NOF} = 1.09$$

Description of data



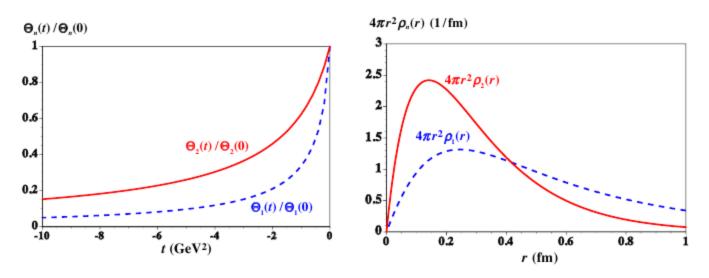
Formfactors

Resonance structure in pressure –related Θ_1



Time-like -> space-like

Dispersion relation and Fourier transform



Mass radius

$$\langle r^2 \rangle = 6 \int_{4m^2}^{\infty} \frac{\text{Im}(F(s))}{s^2}$$

 $\sqrt{\langle r^2 \rangle} = 0.69 \text{ fm for } \Theta_2$

Shear – natural counterpart of pressure (talk of H.-Ch. Kim)

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Forces inside the nucleon on the light front from 3D Breit frame force distributions: Abel tomography case

Julia Yu. Panteleeva¹ and Maxim V. Polyakov⁰, ^{1,2}

¹Ruhr University Bochum, Faculty of Physics and Astronomy, Institute for Theoretical Physics II, D-44870 Bochum, Germany

²Petersburg Nuclear Physics Institute, Gatchina, 188300 St. Petersburg, Russia

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Other traceless structures in EMT? Elastic medium -> liquid (LL v. VII->v. VI)

Shear viscosity in liquid

From spherically symmetric object to fluid (EoS!)

$$T^{\mu\lambda} = (e+p) v^{\mu}v^{\lambda} - p g^{\mu\lambda}$$

 $V^{\mu} = P^{\mu}/M$: correct normalization but no coordinate dependence

Another suggestion:

$$V^{\mu} = (P^{\mu} + a(t) k_{T}^{\mu}) / (M^{2} + a^{2}(t) k_{T}^{2})^{1/2}$$

Viscosity: $\sim E_n p^{[\mu} \Delta^{\lambda]}$

Naïve T-oddness: phases

NO such term in total EMT – violates ExEP (but can be for quarks separately)

Phases <-> dissipation: polarization in pionic superfluidity model (V. I. Zakharov, OT' 17)

Viscosity in GDA channel

Viscosity:will correspond to Exotic JPC=1-+ meson (already studied without reference to viscosity: Anikin, Pire, Szymanowski,OT, Wallon'06)

Spin: related to structure of matrix element: One index of EMT (0th in rest frame) is carried by momentum and other by polarization vector- just what we need for viscosity

No zero-momentum (classical) limit -> quantum

NO for conserved EMT (zero coupling!): violated ExEP

πη pairs observation instead of π π required

Smallness of viscosity: related to smallness of exotic

GDAs and ExEP violation?!

Exotic hybrid meson production

On exotic hybrid meson production in $\gamma^*\gamma$ collisions

I.V. Anikin¹, B. Pire^{2,a}, L. Szymanowski^{3,4,5}, O.V. Teryaev¹, S. Wallon⁵

Eur. Phys. J. C 47, 71–79 (2006) Digital Object Identifier (DOI) 10.1140/epjc/s2006-02533-7

Possible candidate π_1 (1400)

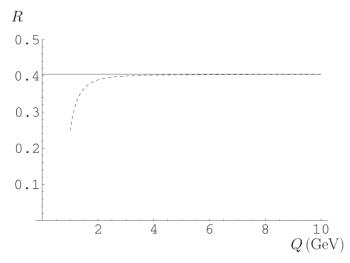


Fig. 2. The ratio $R(Q^2)$ of the squared amplitudes for H and π^0 production in $\gamma^*\gamma$ collisions at leading twist and zero-th order in α_s (solid line) and including twist three contributions in the numerator (dashed line)

Estimate of viscosity

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Terms in EMT:
(e+p) v^{\mu}v^{\lambda} \sim A P^{\mu}P^{\lambda}
 \eta dv^{\mu}/d x_T^{\lambda} \sim E_n p^{[\mu} \Delta^{\lambda]}
TD: e+p \sim Ts
 \eta/s (> 1/(4\pi) \sim E_n T /AM
Correct dependence on Planck constant
   recovered via \Delta^{\lambda}
 (cf K. Trachenko et al.)
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DA vs holographic bound

$$\eta \frac{\partial v^{\nu}}{\partial x_{\mu}} \rightarrow \frac{P^{\nu} \Delta^{\mu}}{M} \sim E(t) P_{\nu} \Delta_{\mu} \qquad v^{\mu} = \frac{P^{\mu} + a(t) \Delta^{\mu}}{\sqrt{M^{2} - a^{2}(t)t}}, \qquad \frac{\partial}{\partial x_{\mu}} \rightarrow i \Delta^{\mu}$$

$$(e + p) v^{\nu} v^{\mu} \rightarrow T s \frac{P^{\nu} P^{\mu}}{M^{2}} \sim A(t) P^{\nu} P^{\mu} \qquad \frac{\eta}{s} \sim \frac{E(t)}{A(t)} \cdot \frac{T}{M} \qquad T \sim \langle K_{T} \rangle$$
Time-like

$$\langle \pi \eta(P, \Delta) | T_i^{\alpha \nu} | 0 \rangle_{\mu^2} = E_i(s, \mu^2) P^{\alpha} \Delta^{\nu}$$

Dimensionful
$$\frac{\eta}{s} \sim \hbar \frac{E(t)}{A(t)} \cdot \frac{T}{M} \sim \frac{\hbar}{k_B} \cdot \frac{E(t)}{A(t)} \cdot \frac{k_B T}{M}$$

$$\hbar \frac{\partial}{\partial x_{\mu}} \rightarrow i \Delta^{\mu}$$

Small bound => small exotic DA, small ExEP violation

Conclusions/Outlook

- ExEP (unique QCD/Gravity relation)
- Time-like Gravitational FFs: may be studied in meson pairs production
- Exotic hybrid mesons: access to shear viscosity and interplay between hadronic and heavy-ion physics
- Holographic bound: related to smallness of exotic GDA and violation of ExEP?
- Shear from asymptotic transition gravitational FFs (Q.-T. Song, OT, work in progress)
- Medium viscosity from GrFFs?