Non-diagonal GPDs and the structure of hadrons

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3D Structure of the Nucleon via Generalized Parton Distributions,

Howard Johnson Incheon Airport Hotel, Incheon, June 25-28, 2024

A plan for today

1. Introduction and (some) general motivation
2. Kinematics of non-diagonal DVCS
3. A simple example: $N \to \Delta$ non-diagonal DVCS
4. $N \to \pi N$ transition GPDs
5. Some lessons from $\pi \to \pi\pi$ transition GPDs;
6. Omnes solution for dispersion relation;
7. Use of the Gribov-Froissart projection;
8. Conclusions and Outlook.

In collaboration with S. Diehl, H.-D. Son, S. Son, P. Sznajder and M. Vanderhaeghen

See the talk by K. Joo.

arXiv:2405.15386v1 [hep-ph] 24 May 2024

Exploring Baryon Resonances with Transition Generalized Parton Distributions: Status and Perspectives

S. Diehl[1,2,a,b], K. Joo[2,a], K. Semenov-Tian-Shansky[3,a], C. Weiss[4,a], V. Braun[5], W.C. Chang[6], P. Chatagnon[4], M. Constantinou[7], Y. Guo[8], P. T. P. Hutauruk[9], H.-S. Jo[3], A. Kim[2], J.-Y. Kim[4], P. Kroll[10], S. Kumano[11], C.-H. Lee[12], S. Liuti[13], R. McNulty[14], H.-D. Son[15], P. Sznajder[16], A. Usman[17], C. Van Hulse[18], M. Vanderhaeghen[19], M. Winn[20]

What is non-diagonal DVCS/DVMP?

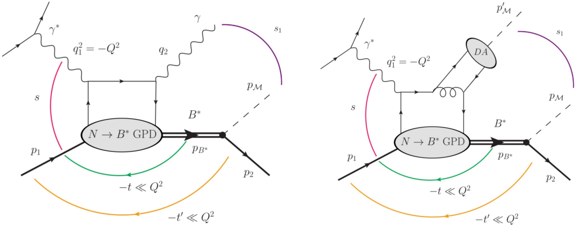
$$\gamma^*(q_1) + N(p_1) \to \left\{ \begin{matrix} \gamma^*(q_2) \\ \mathcal{M}'(p'_{\mathcal{M}}) \end{matrix} \right\} + \left[\mathcal{M}(p_{\mathcal{M}}) N(p') \right]; \; \mathcal{M} = \pi, \, \eta, \rho, \, \omega \, \ldots$$

- Factorized description in terms of $N \to B^*$ GPDs in the generalized Bjorken kinematics:

$$-q_1^2; \; s = (p_1 + q_1)^2; \; s_1 = (p_{\mathcal{M}} + q_2)^2 - \text{ large}; \quad x_B = \frac{-q_1^2}{2p_1 \cdot q_1} - \text{ fixed};$$

$$-t = -(p_{B^*} - p_1)^2; \quad -t' = -(p_2 - p_1)^2; \quad W^2_{\mathcal{M}N} = (p_1 + p_{\mathcal{M}})^2 \quad \text{of hadronic scale.}$$

- Meson-nucleon system resonates at $W_{\mathcal{M}N} = M_{B^*}$.



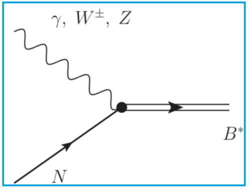
- Status of factorization: same as for the DVCS&DVMP: X. Ji et al.'98, J. Collins et al.'97,99.

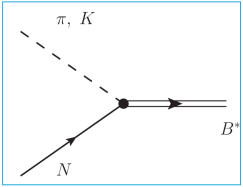
Some motivation

- Main goal is to understand B^* in terms of q, $\bar{q}$ and gluons.
- Available probes and their QCD structure:

$$\text{E.m./weak probe :} \qquad \text{QCD structure :}$$
$$\gamma \quad \Leftrightarrow \quad \langle B^* | \bar{q} \hat{Q}_{\rm e.m.} \gamma_\mu q | N \rangle$$
$$W^\pm,\ Z^0 \quad \Leftrightarrow \quad \langle B^* | \bar{q} \hat{Q}_{w} \gamma_\mu (1-\gamma_5) q | N \rangle$$

- Only $C=-1$ probe;
- Local in space-time;
- No direct access to gluon d.o.f.





$$\text{Hadronic probe :} \qquad \text{QCD structure :}$$
$$\pi,\ K \quad \Leftrightarrow \quad \langle B^* |\, ???\, | N \rangle$$

- QCD structure of the probe unknown;

Graviton probe and QCD Energy-Momentum Tensor

- Graviproduction of resonances I. Kobzarev and L. Okun'62

SOVIET PHYSICS JETP VOLUME 16, NUMBER 5 MAY, 1963

GRAVITATIONAL INTERACTION OF FERMIONS

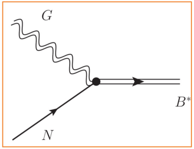
I. Yu. KOBZAREV and L. B. OKUN'

Institute of Theoretical and Experimental Physics, Academy of Sciences, U.S.S.R.

Submitted to JETP editor June 14, 1962

J. Exptl. Theoret. Phys. (U.S.S.R.) **43**, 1904-1909 (November, 1962)

Gravitational interaction of spin-1/2 particles is considered in the linear approximation. It is shown that if gravitational interaction is taken into account, the question whether a free neutrino is two- or four-component acquires a physical meaning. The vertex part for the interaction between fermions and the gravitational field is shown to possess properties analogous to those of the electrodynamic vertex described by the Ward theorem. Observable effects due to spins are considered.



- Gluon d.o.f. enter explicitly!
- No good source of G (:

G probe : QCD structure :

$$G \quad \Leftrightarrow \quad \langle B^* | \underbrace{\bar{q}\gamma_\mu(\partial_\nu - A_\nu)q + \frac{1}{4}F^a_{\mu\alpha}F^a_{\nu\alpha}}_{\text{QCD EMT}} |N\rangle$$

$$\frac{\text{Rate of } GN \to B^*}{\text{Rate of } \gamma N \to B^*} \simeq \frac{m_N}{M_{\rm Pl}} \frac{1}{\alpha_{\rm em}} \simeq 10^{-17}$$

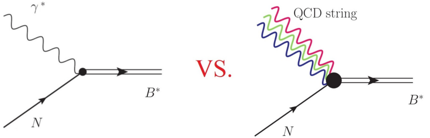
- Instead, can be accessed in hard exclusive reactions.

Some remarks

- Short distance part of the process creates a low-energy QCD string = a tower of local probes (γ, $G, \ldots$);
- Spin J expansion of the QCD string operator:

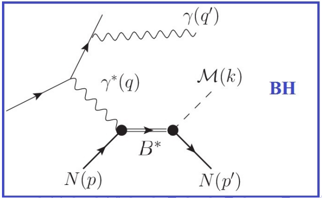
$$\bar{\Psi}(n) P \exp\left(i \int_{-n}^{n} dz^{\mu} A_{\mu}(z)\right) \Psi(-n) =$$

- Although non-diagonal DVCS is a hard process it probes a soft B^* excitation by low-energy QCD string;
- More analogous to B^* photoexcitation rather than hard electroproduction (qualitatively different physics);



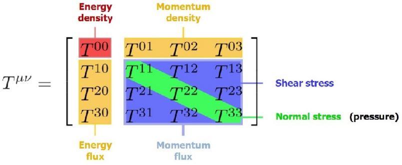
Feasibility:

- Rates are the same order as in usual DVCS/DVMP;
- In case of DVCS: interference with the Bethe-Heitler process provides enhancement of signal;



Physical contents I

Gravitational FFs of the proton, see e.g. V.D. Burkert et al. 2303.08347 Burkert, Elouadrhiri, Girod, Nature 557(2018)



M. Polyakov' 03:

$$T^{ij}(\vec{r}) = \left(\frac{r^i r^j}{r^2} - \frac{1}{3}\delta^{ij}\right) s(r) + \delta^{ij}\, p(r)$$

$$s(r) = -\frac{1}{4M_N}\, r \frac{d}{dr}\frac{1}{r}\frac{d}{dr}\widetilde{D}(r)$$

$$p(r) = \frac{1}{6M_N}\frac{1}{r^2}\frac{d}{dr} r^2 \frac{d}{dr}\widetilde{D}(r)$$

$$\widetilde{D}(r) = \int \frac{d^3\Delta}{(2\pi)^3} e^{-i\vec{\Delta}\cdot\vec{r}} D(-\vec{\Delta}^2).$$

- Study of QCD EMT $N \to B^*$ transition matrix elements complements the studies of e.m. transition FFs;
- Possible access to transition spin contents (for $N \to N^*$, Δ), pressure and shear forces (for $N \to N^*$) and new insight for resonance formation;
- Studies underway. *Cf.* transition angular momentum $N \to \Delta$,

arXiv:2304.08575v1 [hep-ph] 17 Apr 2023

QCD angular momentum in $N \to \Delta$ transitions

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Physical contents II: a unique option for baryon spectroscopy

Important advantages with respect to the usual electroproduction:

1. Excitation of resonances by non-local QCD quark light-cone operators:

$$\left\langle N^* \left| \bar{\psi}_\alpha(0) \mathrm{P} e^{ig \int_0^z dx_\mu \lambda^c A_c^\mu} \psi_\beta(z) \right| N \right\rangle$$

   ★ excitation by probes of arbitrary spin (not just $J = 1$);

2. Possible generalization to the gluon light-cone operators. ★ explicit access to the gluonic DOFs:

$$\langle N^* | G^a_{\alpha\beta}(0) \; [0, z]^{ab} \; G^b_{\mu\nu}(z) | N \rangle$$

3. Direct access to Im (spin asymmetry) and Re (charge asymmetry) of the amplitude $A^{\mathrm{DVCS}}_{N \to B^*}$. Without complicated PWA!

2008 White paper

Baryon spectroscopy in non-diagonal DVCS

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Physical contents III: baryon spectroscopy: hunt for exotics

- Possible access to non-usual spin-flavor configurations: *e.g.* SU(6) $[20, 1^+]$: $N = 2$ orbital excitation of the SU(6) 20-plet.
- SU(3) classification: $3 \otimes 3 \otimes 3 = \underbrace{10}_{S} \oplus \underbrace{8}_{X} \oplus \underbrace{8}_{X} \oplus \underbrace{1}_{A}$

$$\mathrm{SU}(6) \quad |S\rangle = \underbrace{{}^4 10}_{S\cdot S}, \underbrace{{}^2 8}_{X\cdot X}: \quad 56 \text{ states}; \qquad |A\rangle = \underbrace{{}^4 1}_{A\cdot S}, \underbrace{{}^2 8}_{X\cdot X}: \quad 20 \text{ states};$$

$$|X\rangle = \underbrace{{}^2 1}_{A\cdot X}, \underbrace{{}^2 8}_{X\cdot X}, \underbrace{{}^4 8}_{X\cdot S}, \underbrace{{}^2 10}_{S\cdot X} \quad 70 \text{ states};$$

- How to combine with internal orbital motion to make completely symmetrical state?
  - $N = 0$: usual $[56, 0^+]$
  - $N = 1$ (orbital excitation has X-symmetry): $|S\rangle = X \cdot \underbrace{X}_{70} = [70, 1^-]$
  - $N = 2$ (two orbital excitation has X-symmetry make S, X, A with total angular momentum $2, 1, 0$)
    - 20-plet antisymmetric in SU(6) indices ($J = S + L$, $L = -1, 0, 1$):

$$[20, 1^+] =^4 1_{\frac{1}{2}}; \quad {}^4 1_{\frac{3}{2}}; \quad {}^4 1_{\frac{5}{2}}; \quad {}^2 8_{\frac{1}{2}}; \quad {}^2 8_{\frac{1}{2}};$$

- Symmetry argument by R. Feynman'1972: *"Two quark at least must have their motion changed to get to the* $[20, 1^+]$ *from the fundamental* $[56, 0^+]$*."*

Physical contents III: Chiral dynamics in gravitational interaction

- More general description: $N \to \pi N$ transition GPDs, M. Polyakov and S. Stratmann, arXiv:hep-ph/0609045.
- A new test ground for χPT - low energy EFT of QCD, First principle calculations!

PHYSICAL REVIEW D **102,** 076023 (2020)

Chiral theory of nucleons and pions in the presence of an external gravitational field

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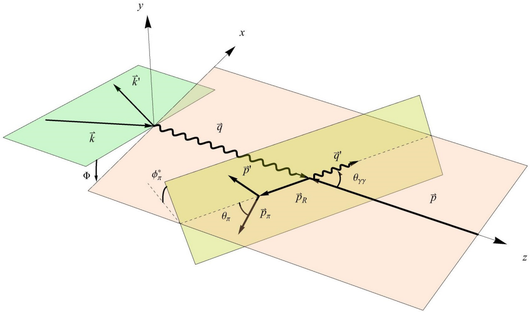
We extend the standard second order effective chiral Lagrangian of pions and nucleons by considering the coupling to an external gravitational field. As an application we calculate one-loop corrections to the one-nucleon matrix element of the energy-momentum tensor to fourth order in chiral counting, and next-to-leading order tree-level amplitude of the pion-production in an external gravitational field. We discuss the relation of the obtained results to experimentally measurable observables. Our expressions for the chiral corrections to the nucleon gravitational form factors differ from those in the literature. That might require to revisit the chiral extrapolation of the lattice data on the nucleon gravitational form factors obtained in the past.

Kinematics and decay angular distribution

$$e(k) + N(p_N) \to e'(k') + \gamma^*(q) + N(p_N) \to e'(k') + \gamma(q') + \pi(p_\pi) + N'(p'_N)$$

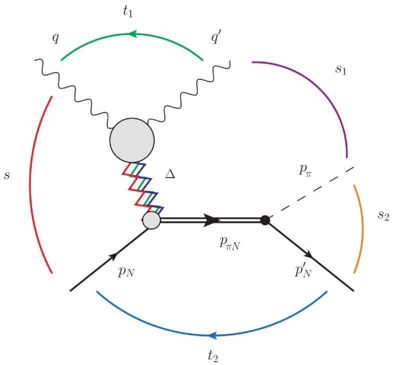


- $\gamma^* N \to B^* \gamma$: $\gamma^* N$ CMS;
- $B^* \to \pi N'$: $\pi N'$ CMS $\equiv$ $(\pi N')$ at rest;

$$\frac{d^7\sigma}{\underbrace{dQ^2 dx_B}_{\text{lepton side}} \underbrace{dt\, d\Phi}_{\gamma^* N \to \gamma B^*} \underbrace{dW^2_{\pi N} d\Omega^*_\pi}_{B^* \to \pi N}}$$

Kinematics: invariants

- Invariant variables for $\gamma^* N \to \gamma\pi N'$



In addition to $s = (p_N + q)^2 \equiv W^2$ and $t_1 = (q - q')^2 \equiv \Delta^2$:

- $\gamma\pi$ invariant mass: $s_1 = (p_\pi + q)^2 \Leftrightarrow \cos\theta^*_\pi$;
- πN invariant mass: $s_2 = (p_\pi + p'_N)^2 \equiv W^2_{\pi N}$;
- $t_2 = (p'_N - p_N)^2 \Leftrightarrow \cos\varphi^*_\pi$;

A test ground: $N \to \Delta(1232)$ DVCS

$$\gamma^*(q) + N^p(p_N) \to \gamma(q') + \Delta^+(p_\Delta) \to \gamma(q') + \pi^0(p_\pi) + N^p(p'_N)$$

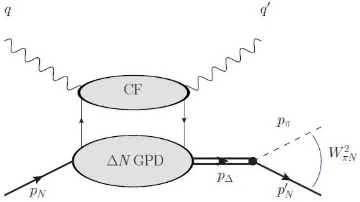
K. Goeke, M.Polyakov and M. Vanderhaeghen'01:

- 3 +1 unpolarized+4 polarized leading twist $N \to \Delta$ GPDs;
- 1 + 2 relevant in the large N_c limit;
- Early analysis: P. Guichon, L. Mossé and M. Vanderhaeghen'03;

A. Belitsky and A. Radyushkin'05:

- 4 unpolarized+4 polarized leading twist $N \to \Delta$ GPDs;

K.S. and M. Vanderhaeghen, PRD 108 (2023)



- Important goal: work out of angular dependencies of $|\text{DVCS}|^2$, $|\text{BH}|^2$ and interference term.
- Implications for experiment: necessary coverage in the cm angle of the final πN state.

$N \to \Delta$ GPDs I

- Leading twist-2: 4 unpolarized and 4 polarized GPDs;
- Unpolarized isovector $N \to \Delta$ GPDs (K. Goeke et al.2001):

$$\frac{1}{2\pi}\int dy^- e^{ixP^+y^-} \left\langle \Delta(p_\Delta) \left| \bar{\psi}(-y/2)\gamma\cdot n\tau_3\psi(y/2) \right| N(p_N) \right\rangle \Big|_{y^+=\vec{y}_\perp=0}$$

$$= \sqrt{\frac{2}{3}}\bar{\mathcal{U}}^\beta(p_\Delta)\Big\{ H_M(x,\xi,t)\left(-\mathcal{K}^M_{\beta\mu}\right)n^\mu + H_E(x,\xi,t)\left(-\mathcal{K}^E_{\beta\mu}\right)n^\mu$$

$$+H_C(x,\xi,t)\left(-\mathcal{K}^C_{\beta\mu}\right)n^\mu + H_4(x,\xi,t)\underbrace{\left(\Gamma^4_{\beta\mu}\right)}_{\text{omitted structure}} n^\mu\Big\} u(p_N),$$

Jones-Scadron covariants ($\bar{P} = \frac{p_N+p_\Delta}{2} = p_\Delta - \frac{\Delta}{2}$, $\Delta = p_\Delta - p_N$, $t \equiv \Delta^2$):

$$\mathcal{K}^M_{\beta\mu} = -i\frac{3(m_\Delta+m_N)}{2m_N\left((m_\Delta+m_N)^2-t\right)}\varepsilon_{\beta\mu\lambda\sigma}\bar{P}^\lambda\Delta^\sigma;$$

$$\mathcal{K}^E_{\beta\mu} = -\mathcal{K}^M_{\beta\mu} - \frac{6(m_\Delta+m_N)}{m_N Z(t)}\varepsilon_{\beta\sigma\lambda\rho}\bar{P}^\lambda\Delta^\rho\varepsilon^\sigma_{\mu\kappa\delta}\bar{P}^\kappa\Delta^\delta\gamma^5;$$

$$\mathcal{K}^C_{\beta\mu} = \not{-}\not{i}\frac{3(m_\Delta+m_N)}{m_N Z(t)}\Delta_\beta\left(t\bar{P}_\mu - \Delta\cdot\bar{P}\Delta_\mu\right)\gamma^5;$$

$$\Gamma^4_{\beta\mu} = \frac{1}{m_N m_\Delta}\left[\Delta_\beta - \frac{(\Delta\cdot p_\Delta)}{p^2_\Delta}p_{\Delta\beta}\right]\Delta_\mu\gamma_5.$$

$N \to \Delta$ GPDs II

- Polarized $N \to \Delta$ GPDs:

$$\frac{1}{2\pi}\int dy^- e^{ixP^+y^-}\left\langle \Delta(p_\Delta)\left|\bar{\psi}(-y/2)\gamma\cdot n\gamma^5\tau^3\psi(y/2)\right| N(p_N)\right\rangle =$$

$$\sqrt{\frac{2}{3}}\bar{\mathcal{U}}^\beta(p_\Delta)\left[C_1(x,\xi,t)g_{\beta\mu}n^\mu + C_2(x,\xi,t)\frac{\Delta_\beta\Delta_\mu}{m_N^2}n^\mu + C_3(x,\xi,t)\frac{1}{m_N}\left[g_{\beta\mu}\not{\Delta} - \Delta_\beta\gamma_\mu\right]n^\mu\right.$$

$$\left. + C_4(x,\xi,t)\frac{2}{m_N^2}\left[\bar{P}\cdot\Delta g_{\beta\mu} - \Delta_\beta\bar{P}_\mu\right]n^\mu\right]u(p_N).$$

Relation to form factors

- Unpolarized GPDs are related to e.m. form factors Jones and Scadron'73:

$$\int_{-1}^{1} dx H_{M,E,C}(x,\xi,t) = 2G^*_{M,E,C}(t); \quad \int_{-1}^{1} dx H_4(x,\xi,t) = 0;$$

- Polarized transition GPDs are related to axial form factors Adler'75;
- These FFs can be accessed in neutrino-production reactions;

$$\int_{-1}^{1} dx C_{1,2,3,4}(x,\xi,t) = 2C^A_{5,6,3,4}(t).$$

Large N_c relations and sum rule

- Large N_c relations for octet-to-decuplet transition GPDs, Goeke et al.'01:

$$H_M(x,\xi,t) = \frac{2}{\sqrt{3}} \left[E^u(x,\xi,t) - E^d(x,\xi,t) \right];$$
$$C_1(x,\xi,t) = \sqrt{3} \left[\tilde{H}^u(x,\xi,t) - \tilde{H}^d(x,\xi,t) \right];$$
$$C_2(x,\xi,t) = \frac{\sqrt{3}}{4} \left[\tilde{E}^u(x,\xi,t) - \tilde{E}^d(x,\xi,t) \right];$$

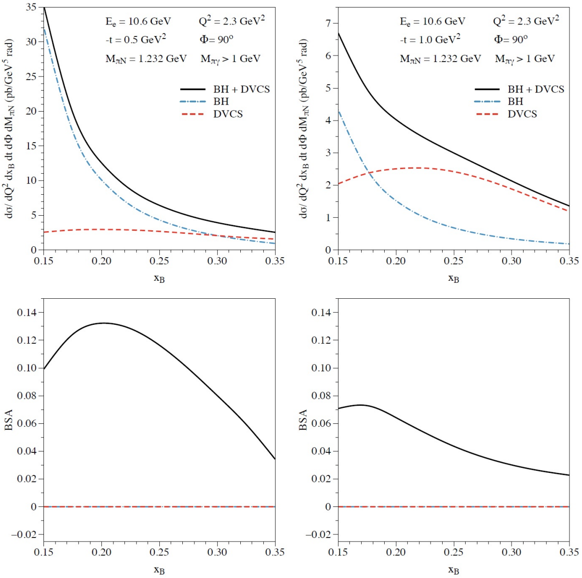
- Pion pole contribution into C_2:

$$\lim_{t \to m_\pi^2} C_2(x,\xi,t) = \sqrt{3} \frac{g_A m_N^2}{m_\pi^2 - t} \theta[\xi - |x|] \frac{1}{\xi} \Phi_\pi \left(\frac{x}{\xi} \right);$$

- Angular momentum sum rule:

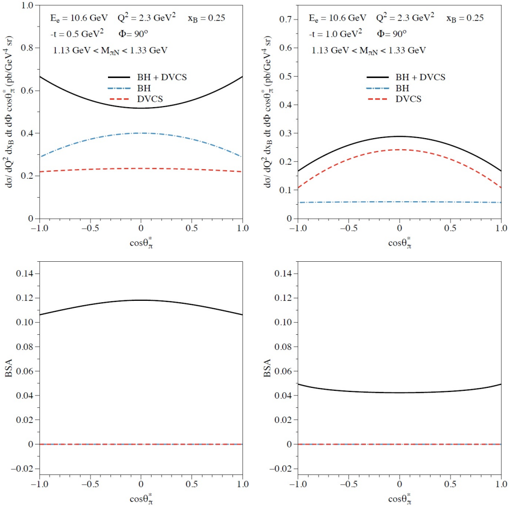
$$\lim_{t \to 0, N_c \to \infty} \int_{-1}^{1} dx x H_M(x,\xi,t) = \frac{2}{\sqrt{3}} \left[2 \left(J^u - J^d \right) - M_2^u + M_2^d \right].$$

Cross sections and BSA for JLab@12 GeV I



Cross sections and BSA for JLab@12 GeV II

- Δ in helicity $\pm 1/2$ state: $\frac{1}{4}\left(1 + 3\cos^2\theta^*_\pi\right)$
- Δ in helicity $\pm 3/2$ state: $\frac{3}{4}\sin^2\theta^*_\pi$



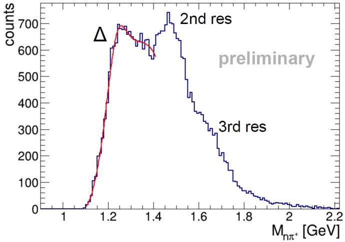
Experimental status I: resonance spectrum for $N^* \to n\pi^+$

Stefan Diehl, CLAS collaboration, preliminary

$en\pi^+\gamma$

$\langle Q^2 \rangle = 2.3$ GeV^2

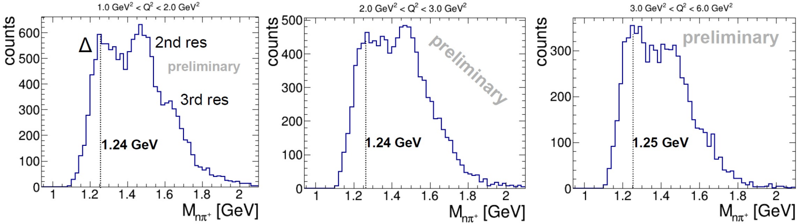
$\langle x_B \rangle = 0.25$



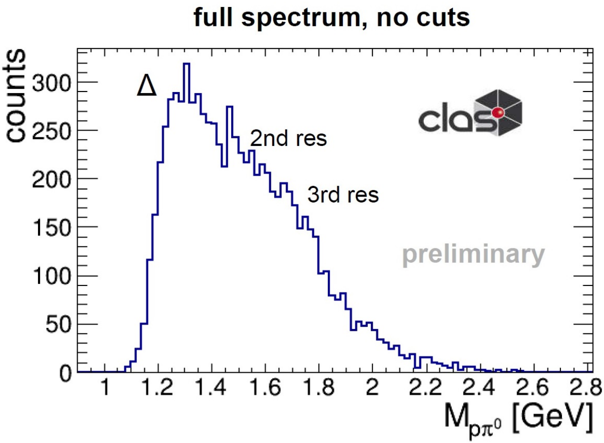
Δ-fit: Breit-Wigner + polyn. backgr.

$\mu = 1.235$ GeV
$\Gamma = 0.15$ GeV

Q^2 dependence:

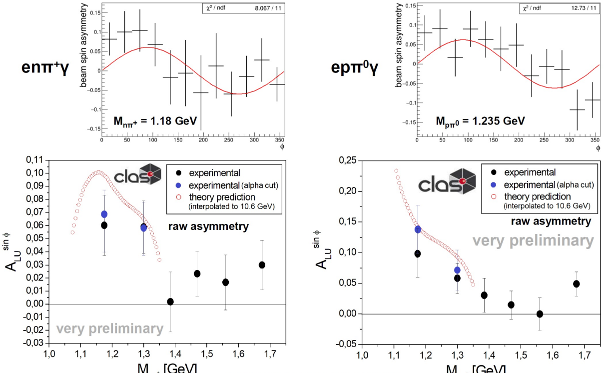


Experimental status II: resonance spectrum for $N^* \to p\pi^0$



Experimental status III: Beam Spin Asymmetry

$$\mathrm{A} = \frac{1}{P}\frac{N^+ - N^-}{N^+ + N^-} \approx A_{LU}^{\sin\phi}\sin\phi$$



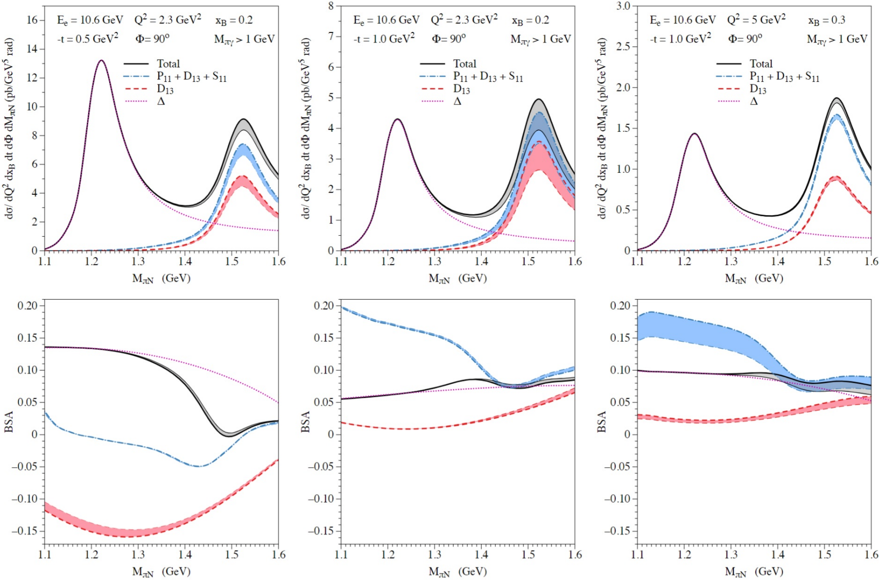
theory prediction: Phys. Rev. D68 (2003)

- BSA $\sim T^{\mathrm{BH}} \times \mathrm{Im}\, T^{N\Delta\,\mathrm{DVCS}}$.

Going to the 2nd resonance region

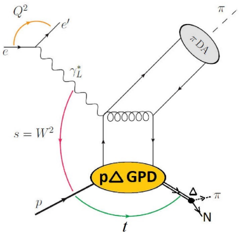
- Formalism extended to $N \to N^*$ DVCS for $N^* = \mathrm{P}_{11}(1440)$, $\mathrm{D}_{13}(1520)$, $\mathrm{S}_{11}(1535)$:
  - for spin-$\frac{1}{2}$ resonances at twist-2: 2 unpolarized GPDs (vector operator), 2 polarized GPDs (axial-vector operator);
  - for spin-$\frac{3}{2}$ resonances at twist-2: 4 unpolarized GPDs (vector operator), 4 polarized GPDs (axial-vector operator);
- t-dependence of GPDs (first moments):
  - - unpolarized GPDs: first moments constrained by data on e.m. transition FFs (CLAS@6 GeV)
  - - polarized GPDs: 2 dominant axial FFs constrained using PCAC + pion pole dominance:
    - normalization at $t = 0$ given by $(f_{\pi NN^*}/m_\pi)2f_\pi$;
    - t-dependence: dipole ($M_A = 1$ GeV) and pion-pole $\sim 1/(t - m_\pi^2)$;
    - isoscalar axial FF neglected;
- x & ξ dependence of GPDs: RDDA $b = 1$ and $b = \infty$ with $q(x) \sim x^{-0.5}(1-x)^3$

Cross section and BSA



Hard exclusive $\Delta\pi$ production

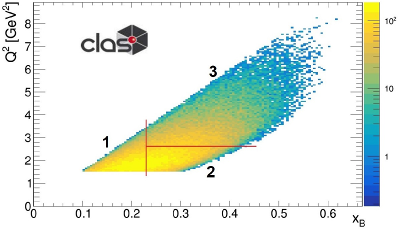
S. Diehl et al. '23

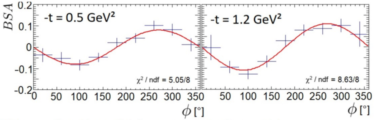


$$ep \to e\Delta^0\pi^+ \to e(p\pi^-)\pi^+$$
$$\to e(n\pi^0)\pi^+$$

$$ep \to e\Delta^+\pi^0 \to e(n\pi^+)\pi^0$$
$$\to e(p\pi^0)\pi^0$$

$$ep \to e\Delta^{++}\pi^- \to ep\pi^+\pi^-$$





BSA as a function of ϕ for representative $-t$ bins (Q^2= 2.48 GeV^2, x_B = 0.27). The red line shows the $\sin\phi$ fit.

- Amplitude involves polarized GPDs $C_{1,2,3,4}(x, \xi, \Delta^2)$;
- BSA is a twist-3 effect;

Experimental perspectives

- $N\Delta$ DVCS and $\pi\Delta$ can be measured at CLAS. Analysis underway.
- Present status: 3-4 bins in $-t$. With extra angular variables 2-3 bins in each variable;
- Statistics increase by a factor 3 in 3-4 years;
- BSA $\pi^-\Delta^{++}$ extracted;
- Possible JLab@20 upgrade: statistics may increase by a factor 100 - 1000;

arXiv:2306.09360v1 [nucl-ex]

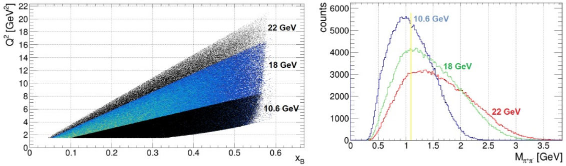


Figure 40: Comparison of the available phase space, accessible with the present CLAS12 setup, in $Q^2 - x_B$ or the $\pi^-\Delta^{++}$ process under forward kinematics ($-t < 1.5$ GeV^2) (left) and for the $\pi^+\pi^-$ invariant mass of the same process, which is used to suppress the dominant ρ production background by the cut on $M(\pi^+\pi^-) > 1.1$ GeV, indicated by the yellow line (right) for a 10.6 GeV, 18 GeV and 22 GeV electron beam.

- Can we get access to the complete angular distribution of $N\Delta$ DVCS/DVMP and $\pi\Delta$ production cross section?
- A sizable $\pi^-\Delta^{++}$ BSA a challenge for theory: twist-3 observable;
- Extension to small-x_B and studies for the EIC conditions necessary;

$N \to \pi N$ transition GPDs

M. Polyakov and S. Stratmann, arXiv:hep-ph/0609045

- Unpolarized $N \to \pi N$ GPDs:

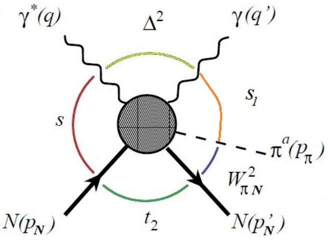
$$\int \frac{d\lambda}{2\pi} e^{i\lambda x \bar{P}\cdot n} \langle N(p'_N)\, \pi^a(p_\pi) | \bar{\psi}(-\lambda n/2) \not{n} \psi(\lambda n/2) | N(p_N) \rangle = \frac{i g_A}{m_N f_\pi} \sum_{i=1}^{4} \bar{U}(p'_N)\, \Gamma_i \tau^a H_i^{(0)} U(p_N)$$

$$\Gamma_1 = \gamma_5; \quad \Gamma_2 = \frac{m_N \not{n}}{n \cdot \bar{P}} \gamma_5; \quad \Gamma_3 = \frac{\not{k}}{m_N} \gamma_5; \quad \Gamma_4 = \frac{\not{k} \not{n}}{m_N} \gamma_5; \quad \left(\bar{P} = \frac{p'_N + p_N + p_\pi}{2}\right)$$

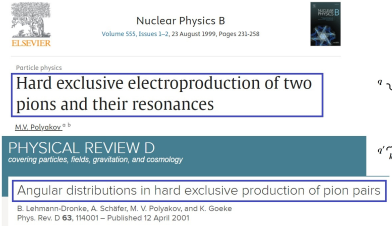
A guide to the kinematical variables of $H_i^{(0)}(x, \xi, \Delta^2; W^2_{\pi N}, \alpha, t_2)$:

- πN invariant mass $W^2_{\pi N} = (p' + p_\pi)^2$
- $t_1 = (p'_N + p_\pi - p_N)^2 = (q - q')^2 \equiv \Delta^2$
- $t_2 = (p'_N - p_N)^2$
- Skewness $\xi = -\frac{n \cdot \Delta}{2 n \cdot \bar{P}}$
- Relative pion longitudinal momentum of the πN system:

$$\alpha = \frac{n \cdot p_\pi}{n \cdot (p'_N + p_\pi)}$$



$\pi\pi$ Generalized Distribution Amplitudes (GDAs)


Nuclear Physics B
Volume 555, Issues 1–2, 23 August 1999, Pages 231-258

Particle physics

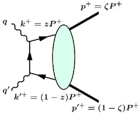
Hard exclusive electroproduction of two pions and their resonances

M.V. Polyakov [a b]

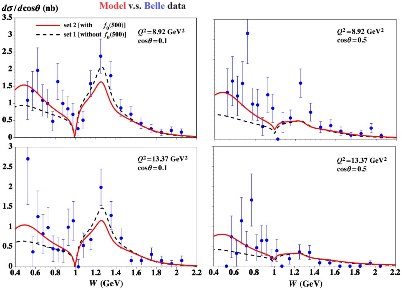
PHYSICAL REVIEW D
covering particles, fields, gravitation, and cosmology

Angular distributions in hard exclusive production of pion pairs

B. Lehmann-Dronke, A. Schäfer, M. V. Polyakov, and K. Goeke
Phys. Rev. D **63**, 114001 – Published 12 April 2001



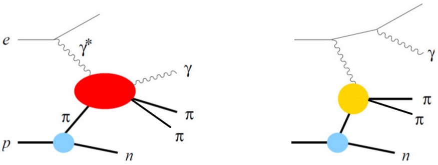
- A general description on 2π system for $\gamma^*\gamma \to \pi\pi$ both at the soft pion threshold and in the $\pi\pi$ resonance region;
- For applications see e.g. S. Kumano, Qin-Tao Song, and O. Teryaev PRD 97 (2018);



A test ground for the formalism: $\pi \to \pi\pi$ ND DVCS

$$e(l) + p(p) \to e(l') + \gamma(q') + \pi^+(k) + n(p') \to e(l') + \gamma(q') + \pi^+(k_1) + \pi^0(k_2) + n(p')$$

- Can be studied through the Sullivan-type process:



- No complications due to spin-$\frac{1}{2}$.
- Access to the meson spectrum: $\rho(770)$, $f_2(1270)$ *etc.*
- An option for the EIC?

Some experimental prospects?

Few-Body Syst (2023) 64:38
https://doi.org/10.1007/s00601-023-01812-1

J. M. Morgado Chávez · V. Bertone · F. De Soto · M. Defurne · C. Mezrag · H. Moutarde · J. Rodríguez Quintero · J. Segovia

Generalized Parton Distributions of Pions at the Forthcoming Electron-Ion Collider

- **N.B.** $\gamma^* N \to \rho N' \to \pi\gamma N'$ a background for $N \to \Delta$ DVCS.

$\pi \to \pi\pi$ transition GPDs

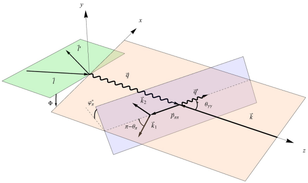
- $\pi \to \pi\pi$ unpolarized transition GPD ($\bar{P} \equiv \frac{k+k_1+k_2}{2}$):

$$\frac{1}{2}\int \frac{d\lambda}{2\pi} e^{i\lambda x \bar{P}\cdot n} \langle \pi(k_1)\pi(k_2)|\bar{\psi}\left(-\frac{\lambda n}{2}\right) \not{n} \psi\left(\frac{\lambda n}{2}\right)|\pi(p_\pi)\rangle$$
$$= \frac{1}{2\bar{P}\cdot n} i\varepsilon(n, \bar{P}, \Delta, k_1) \frac{1}{f_\pi^3} H_{\pi\to\pi\pi}(x, \xi, \Delta^2, W_{\pi\pi}^2, \theta_\pi^*, \varphi_\pi^*);$$

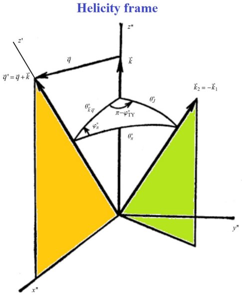
- $\pi \to \pi\pi$ polarized transition GPD:

$$\frac{1}{2}\int \frac{d\lambda}{2\pi} e^{i\lambda x \bar{P}\cdot n} \langle \pi(k_1)\pi(k_2)|\bar{\psi}\left(-\frac{\lambda n}{2}\right) \not{n}\gamma_5 \psi\left(\frac{\lambda n}{2}\right)|\pi(p_\pi)\rangle$$
$$= \frac{1}{2\bar{P}\cdot n} (\bar{P}\cdot n) \frac{1}{f_\pi} \tilde{H}_{\pi\to\pi\pi}(x, \xi, \Delta^2, W_{\pi\pi}^2, \theta_\pi^*, \varphi_\pi^*);$$

- Transition GPD arguments: x, ξ, $\Delta^2 = t$ and of the invariant mass of $\pi\pi$ system $W_{\pi\pi}^2$ and the helicity frame pion decay angles θ_π^*, φ_π^*.



Angles in the helicity frame



- $\cos\theta_\pi^*$ is linear in $s_1 = (q' + p_\pi)^2$;
- $\cos\varphi_\pi^*$ is linear in $t_2 = (p'_N - p_N)^2$;
- Polar and azimuthal angle through the Gram determinants:

$$\cos\theta_\pi^* = \frac{G_2\begin{pmatrix} k_1+k_2, q' \\ k_1+k_2, k_1 \end{pmatrix}}{\left\{\Delta_2\left(k_1+k_2, q'\right)\Delta_2\left(k_1+k_2, k_2\right)\right\}^{\frac{1}{2}}};$$

$$\sin^2\varphi_\pi^* = \frac{\Delta_2\left(k+q, q'\right)\Delta_4\left(k+q, q', k, k_2\right)}{\Delta_3\left(k+q, q', k\right)\Delta_3\left(k+q, q', k_2\right)};$$

- Gram determinants:

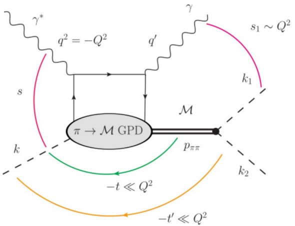
$$\mathrm{G}_n\begin{pmatrix} p_1, \ldots, p_n \\ q_1, \ldots, q_n \end{pmatrix} = \det\left(p_i \cdot q_j\right);$$

- Symmetric Gram determinants:

$$\Delta_n\left(p_1, \ldots p_n\right) = \mathrm{G}_n\begin{pmatrix} p_1, \ldots, p_n \\ p_1, \ldots, p_n \end{pmatrix} = \det\left(p_i \cdot p_j\right)$$

How to treat the angular structure? Real-valued spherical harmonics.

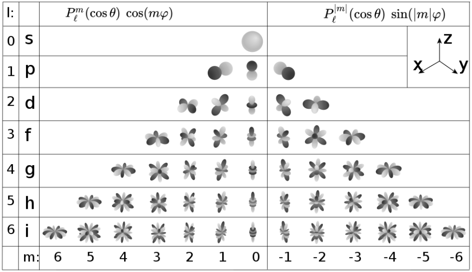
- Partial wave expansion both in $\theta^*_\pi \Leftrightarrow \alpha$ and φ^*_π.



$$Y_\ell^m(\theta^*_\pi, \varphi^*_\pi) = (-1)^m \sqrt{\frac{(2\ell+1)}{4\pi}\frac{(\ell-m)!}{(\ell+m)!}}\, P_\ell^m(\cos\theta - \pi^*) e^{im\varphi^*_\pi}$$

the real-valued spherical harmonics read :

$$Y_\ell^m = \begin{cases} \frac{1}{\sqrt{2}}\left(Y_{\ell,|m|} - i\, Y_{\ell,-|m|}\right) & \text{if } m<0; \\ Y_{\ell,0} & \text{if } m=0; \\ \frac{(-1)^m}{\sqrt{2}}\left(Y_{\ell,|m|} + i\, Y_{\ell,-|m|}\right) & \text{if } m>0; \end{cases}$$



PW expansion of $\pi \to \pi\pi$ GPDs

- PW expansion of CFFs in angles θ^*_π and φ^*_π for unpolarized GPD:

$$H_{\pi\to\pi\pi}(x,\xi,t,W^2_{\pi\pi},\theta^*_\pi,\varphi^*_\pi) = \frac{1}{\sqrt{1-\cos^2\theta^*_\pi}\,|\sin\varphi^*_\pi|}\sum_{\ell=1}^{\infty}\sum_{m=-\ell}^{-1} H^{\ell m}_{\pi\to\pi\pi}(x,\xi,t,W^2_{\pi\pi})\,\mathrm{Y}_{\ell m}(\theta^*_\pi,\varphi^*_\pi);$$

N.B. Spherical harmonics in are odd under $\varphi^*_\pi \to -\varphi^*_\pi$.

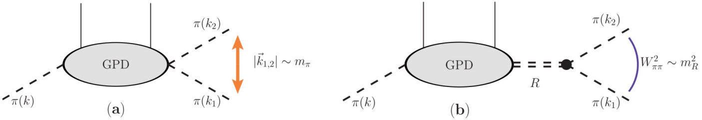
- PW expansion in angles θ^*_π and φ^*_π for polarized GPD:

$$\tilde{H}_{\pi\to\pi\pi}(x,\xi,t,W^2_{\pi\pi},\theta^*_\pi,\varphi^*_\pi) = \sum_{\ell=0}^{\infty}\sum_{m=0}^{\ell} \tilde{H}^{\ell m}_{\pi\to\pi\pi}(x,\xi,t,W^2_{\pi\pi})\,\mathrm{Y}_{\ell m}(\theta^*_\pi,\varphi^*_\pi);$$

N.B. Spherical harmonics in are even under $\varphi^*_\pi \to -\varphi^*_\pi$.

How to go beyond the threshold? I (in collaboration with H. Son)

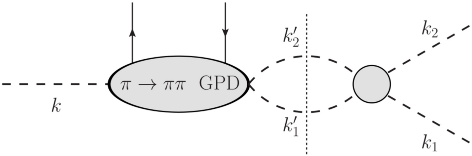
- Two regimes



(a) $\pi \to \pi\pi$ transition GPD in the chiral regime; **(b)** $\pi \to \pi\pi$ transition GPD in the resonance region;

- The Watson'54 final state interaction theorem for $\pi \to \pi\pi$ transition GPD:

$$\text{for} \quad W^2_{\pi\pi} < 16 m_\pi^2 : \quad \text{Im}\, \tilde{H}^I_{\pi \to \pi\pi}(x, \xi, w^2, \theta^*_\pi, \varphi'_\pi)$$
$$= \frac{1}{2!} \int d(\text{phase space}) \left(\tilde{H}^I_{\pi \to \pi\pi}(x, \xi, w^2, \theta'_\pi, \varphi'_\pi) \right)^* A^I_{\pi\pi}(k_1, k_2 | k'_1, k'_2)$$



How to go beyond the threshold? II

- $\pi\pi$-scattering amplitude:

$$A^I_{\pi\pi} = 8\pi W_{\pi\pi} \sum_\ell (2\ell+1) a^I_\ell(W^2_{\pi\pi}) P_\ell \left[\cos\left(\theta_{\rm cm}\right)\right].$$

- Elastic unitarity condition:

$$\text{Im } a^I_\ell(W^2_{\pi\pi}) = |\vec{k}_1||a^I_\ell(W^2_{\pi\pi})|^2;$$

- $\delta^I_\ell\left(W^2_{\pi\pi}\right)$ are the $\pi\pi$ scattering phases:

$$a^I_\ell(W^2_{\pi\pi}) = \frac{1}{|\vec{k}_1|} \sin\left[\delta^I_\ell\left(W^2_{\pi\pi}\right)\right] e^{i\delta^I_\ell\left(W^2_{\pi\pi}\right)}.$$

How to go beyond the threshold? II

- The equation for the expansion coefficients $\tilde{H}^I_{\ell;m}$:

$$\text{Im}\,\tilde{H}^I_{\ell;m}(x,\xi,w^2) = \tan\left[\delta^I_\ell\left(w^2\right)\right]\,\text{Re}\,\tilde{H}^I_{\ell;m}(x,\xi,w^2).$$

- Omnes'58: N-subtracted dispersion relation

$$\tilde{H}^I_{\ell;m}(x,\xi,w^2)$$
$$= \sum_{k=0}^{N-1} \frac{w^{2k}}{k!}\frac{\mathrm{d}^k}{\mathrm{d}w^{2k}}\tilde{H}^I_{\ell;m}(x,\xi,w^2=0) + \frac{w^{2N}}{\pi}\int_{4m_\pi^2}^{\infty}\mathrm{d}\omega\,\frac{\tan\left[\delta^I_\ell(\omega)\right]\,\text{Re}\left\{\tilde{H}^I_{l;m}(x,\xi,\omega)\right\}}{s^N\left(\omega - w^2 - i\epsilon\right)}.$$

- The Omnes solution (for $N=0$):

$$\tilde{H}^I_{\ell;m}(x,\xi,W^2) = \tilde{H}^I_{\ell;m}(x,\xi,W^2=0)\exp\left[\frac{1}{\pi}\int_{4m_\pi^2}^{\infty}\mathrm{d}s\,\frac{\delta^I_\ell(s)}{s-m^2_{\pi\pi}-i\epsilon}\right]$$

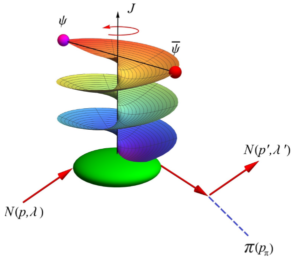
- Transition GPDs are complex functions above threshold!

Can we handle with QCD string for the non-diagonal case?

- Hard part of the DVCS creates a soft QCD string.

$$\bar{q}(z)\gamma_\mu \mathrm{Pexp}\left\{ i\int_0^1 dx^\nu A_\nu(x)\right\} q(0)\Big|_{z\to 0} = z^\nu \underbrace{\bar{q}\gamma_\mu \nabla_\nu q}_{\mathrm{Spin}-2:\ q-\mathrm{part\ of\ EMT}} + z^\nu z^\rho \underbrace{\bar{q}\gamma_\mu \nabla_\nu \nabla_\rho q}_{\mathrm{Spin}-3} + \ldots$$

- How to decompose QCD string into probes of different spin? A tool is provided by the Froissart-Gribov projection ⇔ Abel tomography K.S. and P. Sznajder, PRD 109 (2024).
- Froissart-Grbov projection in the context of DVCS: K. Kumericki, D. Müller, and K. Passek-Kumericki, Eur. Phys. J. C **58**, 193 (2008); M. Polyakov, Phys. Lett. B **659**, 542 (2008);



What is the Froissart- Gribov projection: case of spinless hadrons

Gribov'61, Froissart'61

- t-channel counterpart of the DVCS reaction:

$$\gamma^*(q) + \gamma\left(\bar{q}'\right) \to h\left(p'\right) + \bar{h}(\bar{p});$$

- θ_t: the angle between $\vec{q}$ and $\vec{\bar{p}}$ in the $\gamma^*\gamma$ CMS; in the DVCS kinematics:

$$\cos\theta_t \to -\frac{1}{\xi\sqrt{1-\frac{4m^2}{t}}} + \mathcal{O}\left(1/Q^2\right);$$

- At the moment we set $m=0$. Effect of non-zero target mass: mixing of PWs.

DR for the CFF analytically continued to the t-channel:

$$\mathcal{H}_+(\cos\theta_t, t) = \int_0^1 dx \frac{2x\cos^2\theta_t}{1-x^2\cos^2\theta_t} H_+(x,x,t) + 4D(t).$$

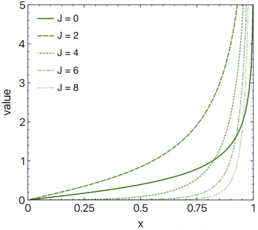
- Cross channel SO(3) PWAs

$$F_J(t) \equiv \frac{2J+1}{2}\int_{-1}^1 d(\cos\theta_t)\, P_J(\cos\theta_t)\mathcal{H}^{(+)}(\cos\theta_t, t).$$

Froissart- Gribov projection II

Neumann's integral representation for the Legendre functions $\mathcal{Q}_J$ ($J \geq 0$, integer):

$$\frac{1}{2}\int_{-1}^{1} dz' \, P_J(z') \frac{1}{z-z'} = \mathcal{Q}_J(z);$$



First weight functions $2(2J+1)\left(\frac{\mathcal{Q}_J(1/x)}{x^2} - \delta_{J0}\frac{1}{x}\right)$

- For even positive J

$$F_{J>0}(t) = 2(2J+1)\int_0^1 dx \frac{\mathcal{Q}_J(1/x)}{x^2} H_+(x,x,t)\,.$$

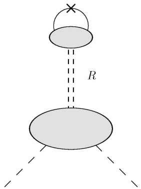
- Special case $J=0$:

$$F_{J=0}(t) = 2\int_0^1 dx \left[\frac{\mathcal{Q}_0(1/x)}{x^2} - \frac{1}{x}\right] H_+(x,x,t) + 4D(t)\,.$$

- N.B. $\frac{\mathcal{Q}_J(1/x)}{x^2} \sim x^{J-1}$ for small x.

t-channel point of view and duality

- Crossing relation between GPD and two particle GDA.
- Duality in the spirit of R. Dolen, D. Horn, C. Schmid'67. GPDs are presented as infinite series of t-channel Regge exchanges M. Polyakov'98:



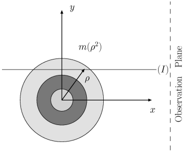
$$\langle \pi(p')| \hat{O} |\pi(p)\rangle \sim \text{ Crossing of } \sum_{R_J} \sum_{\substack{\text{polarization}\\ \text{of } R_J}} \frac{1}{t - M^2_{R_J}}$$

$$\times \underbrace{\langle \pi(p')\pi(-p)|R_J\rangle}_{R_J\pi\pi \text{ effective vertex}} \underbrace{\langle R_J| \hat{O} |0\rangle}_{\text{F.T. of DA of } R_J}$$

- Spin sum of $R_J \sim P_J(\cos\theta_t)$
- Dual parametrization of GPDs

$$H_+(x,\xi,t) = \sum_{\substack{n=1\\ \text{odd}}}^{\infty} \sum_{\substack{l=0\\ \text{even}}}^{n+1} B_{nl}(t)\, \theta\left(1 - \frac{x^2}{\xi^2}\right) \left(1 - \frac{x^2}{\xi^2}\right) C_n^{\frac{3}{2}}\left(\frac{x}{\xi}\right) P_l\left(\frac{1}{\xi}\right)$$

Relation to the Abel transform tomography



The observer at ∞ looking along a line parallel to the x-axis a distance y above the origin sees the projection:

$$a(y^2) = \int_{-\infty}^{\infty} dx\, m(\rho^2) = \int_{y^2}^{\infty} d\rho^2 \frac{m(\rho^2)}{\sqrt{\rho^2 - y^2}}$$

- M. Polyakov'07: the Joukowski conformal map $\frac{1}{w} = \frac{1}{2}\left(y + \frac{1}{y}\right)$ allows to present the relation between $\mathrm{Im}\mathcal{H}(\xi, t)$ and $N(y, t)$ in form of the Abel integral equation.
- The inverse transform for $N(y)$:

$$N(y,t) = \frac{1}{\pi} \frac{y(1-y^2)}{(1+y)^{\frac{3}{2}}} \int_{\frac{2y}{1+y^2}}^{1} \frac{d\xi}{\xi^{\frac{3}{2}}} \frac{1}{\sqrt{\xi - \frac{2y}{1+y^2}}} \left\{ \frac{1}{2} \mathrm{Im}\mathcal{H}^{(+)}(\xi,t) - \xi \frac{d}{d\xi} \mathrm{Im}\mathcal{H}^{(+)}(\xi,t) \right\}.$$

- For massless hadrons:
$\int_0^1 dx x^{J-1} N(x,t) = B_{J-1\,J}(t) + B_{J+1\,J}(t) + B_{J+3\,J}(t) + \ldots \equiv F_J(t)$.

On the physical content of the Froissart-Gribov FFs

An answer we want to have:

"FG FFS $F_J(t)$ quantifies hadron target response on the string-like QCD probe with fixed angular momentum J."

- What we have to support this point of view?
- Shortcomings we need to face:
  - Analytic continuation in t;
  - Target mass corrections;
- Possible benefits of $F_J(t)$ as observable quantities?

Generalization of the FG projections for non-diagonal transitions

- To the leading accuracy in $1/Q^2$,

$$\cos\theta_t = -\frac{1}{\xi}\frac{t+\xi(W_{\pi\pi}^2 - m_\pi^2)}{\Lambda(t, W_{\pi\pi}^2, m_\pi^2)} + \mathcal{O}(\frac{1}{Q^2}) \equiv -\frac{1}{b\xi} + a + \mathcal{O}(\frac{1}{Q^2}).$$

Complicated mixing issue.

- The "topological" tensor structure $\varepsilon(n, \bar{P}, \Delta, k_1)$ makes unpolarized $\pi \to \pi\pi$ GPD expanded in $P'_J(\cos\theta_t)$;
- For the polarized case, for $J \geq 0$ and $0 \leq m \leq \ell$:

$$\tilde{F}_J^{\ell m}(t) = 2\int_0^1 dx \tilde{H}^{\ell,m}(x, x, t, W_{\pi\pi}^2)(2J+1)\frac{\mathcal{Q}_J(1/x)}{x^2}.$$

- The FG projections with $J \geq 1$ and $-\ell \leq m < 0$ of the unpolarized Compton transition FFs is similar to the case of magnetic combination of nucleon GPDs $(H+E)$:

$$F_J^{\ell m}(t) = 2\int_0^1 dx H^{\ell,m}(x, x, t, W_{\pi\pi}^2)\frac{2J+1}{J(J+1)}\frac{(-1)}{x}\sqrt{\frac{1}{x^2}-1}\,\mathcal{Q}_J^1(1/x),$$

where $\mathcal{Q}_J^1(1/x)$ stand for the associated Legendre functions of the second kind.

Summary and Outlook

1. New tool for baryon spectroscopy: arbitrary spin-J probe and PW analysis of excited states;
2. A new bridge between PW analysis and QCD;
3. Access to $N \to N^*$ EMT matrix elements: mechanical properties of resonances;
4. A lab for chiral perturbation theory on the light cone: soft pion theorems and chiral expansion;
5. GPD formalism worked out for $N \to \Delta(1232)$, $P_{11}(1440)$, $D_{13}(1520)$, $S_{11}(1535)$. Can be studied at JLab@12 GeV and an option for JLab@22 GeV;
6. A development for hyperons $N \to \Lambda, \Sigma$ and production of strange mesons?
7. $\pi \to \pi\pi$ and $N \to \pi N$ transition GPDs emerge as a tool to study the spectrum of hadrons;
8. First step: development of the formalism for $\pi \to \pi\pi$ transition GPDs: Abel tomography, threshold theorems and the Omnes dispersion relations;
9. Froissart-Gribov projection as a means to focus on specific J in the cross channel;

Thank you for your attention!