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Quantum Simulation of Light-Front Quantum Field Theory

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ILCAC Seminar

Outline

1. What quantum computers are?
2. How can QCs accelerate solving LF QFT?
3. Peculiarities of simulating LF QFT on QCs.
4. Past, Current, & Future work.

1. What is a QC?

What is a QC?

- Wiki:

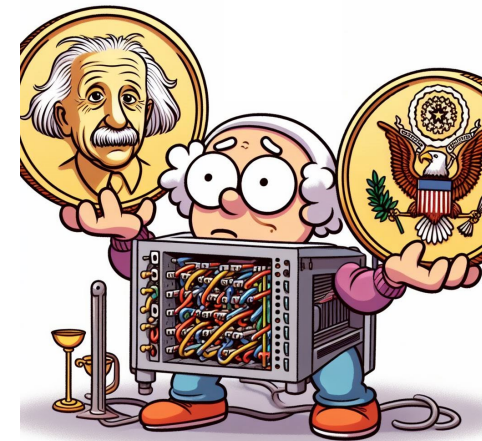
“A quantum computer is a computer that takes advantage of quantum mechanical phenomena.”

- What I see:

“A system capable of naturally storing superpositions of states which can be manipulated and measured with sufficient precision to enable its utility as a computer.”



Google Gemini AI



Microsoft Designer AI

What is a QC?

User perspective

- Wiki:

“A classical computer can solve the same computational problems as a quantum computer, given enough time. Quantum advantage comes in the form of time complexity rather than computability.”




- What I see:

“Quantum computers are (potentially?) powerful calculators for (highly?) specific tasks.”

What is a QC?

Physicist perspective

Why would a physicist be interested in QC?

- Running previously inaccessible calculations / simulations 
- Thinking differently 
 - Quantum-inspired classical algorithms (lin. alg.; q. chem; etc.);
 - Hamiltonian approaches to QFT*
(ET Hamiltonian QFT offers many lessons for LF QFT);
 - Investigations of computational complexity.
- Scientific cross-pollination 

Digital QCs

- States of (deterministic) classical and quantum computers:

Type	Initial State	Intermediate State	Final State
Classical	0	0, 1	0, 1
	00	00, 01, 10, 11	00, 01, 10, 11
	...	$\{0,1\}^{\times n}$...
Quantum	$ 0\rangle$	$c_0 0\rangle + c_1 1\rangle$	0, 1
	$ 00\rangle$	$ \psi\rangle = c_{00} 00\rangle + c_{01} 01\rangle + c_{10} 10\rangle + c_{11} 11\rangle$	00, 01, 10, 11
	...	$\{ 0\rangle, 1\rangle\}^{\otimes n}$...



Reversible computation



Measurement

- QM measurements have probabilistic nature.

Quantum Computation

- Life story of a state in quantum computer:

- Initialization 📱

- Unitary evolution 🧬

- Measurement 📄

$$|\psi_i\rangle = |0\rangle^{\otimes n}$$

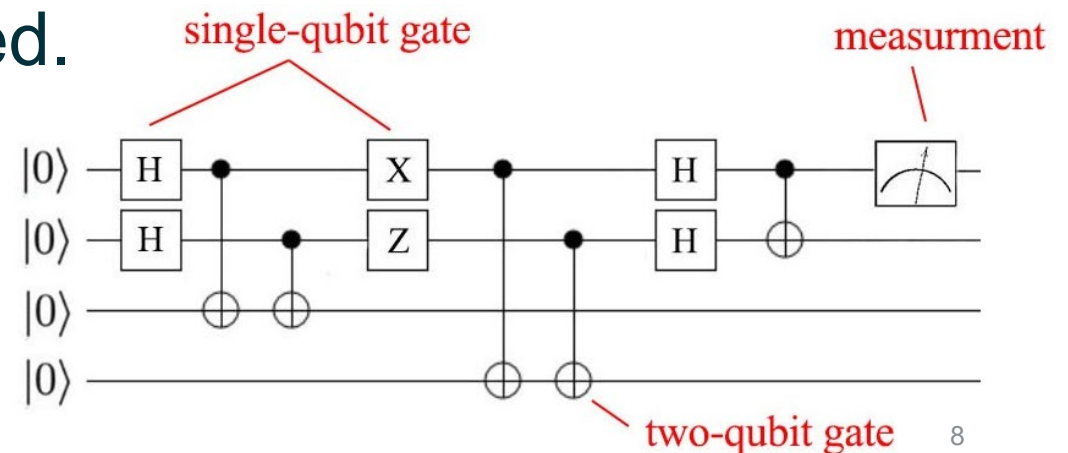
$$|\psi_f\rangle = \dots U_3 U_2 U_1 |\psi_i\rangle$$

$$|\psi\rangle \mapsto \{0, 1\}^{\otimes n}$$

- These stages can be conveniently represented by a circuit diagram:

- 1-, 2-, or many-qubit gates can be used.

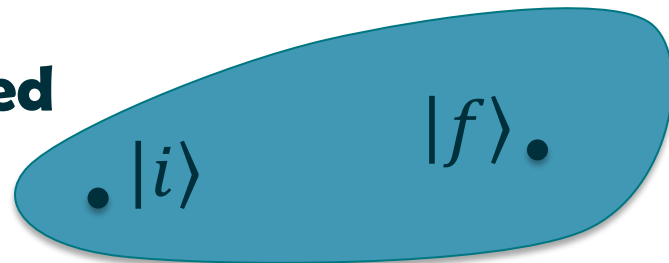
- Efficient algorithm:
#{qubits, gates, circuit runs} are polynomial in problem parameters.



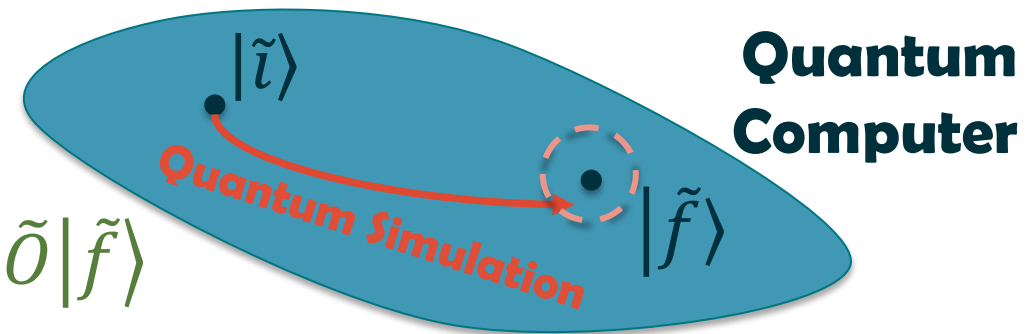
Quantum Simulation

- Choose a discretized model for the system of interest. Contingent on:
 - Observables of interest,
 - Choice of computational tools.
- Mapping physical DOFs onto qubits, e.g., $c_0|\uparrow\rangle + c_1|\downarrow\rangle \rightarrow c_0|0\rangle + c_1|1\rangle$.
- Quantum Computation:

Discretized Model



$$\langle f|O|f\rangle \approx \langle \tilde{f}|\tilde{O}|\tilde{f}\rangle$$



- $|\tilde{i}\rangle \rightarrow |\tilde{f}\rangle$ *can* but does not have to resemble some physical process. (Such as discretized time evolution or adiabatic interaction turn-on.)

Remarks

- “Fault-tolerance”:
Errors in QCs accumulate at an exp rate \Rightarrow long circuits require error correction via quantum analogs of repetition codes, e.g. $|0\rangle \mapsto |000\rangle$.
- We are at an early stage (not even close to punch cards):
Current quantum algorithms address operations on individual qubits.
- In the context of studying quantum physics, exponential advantage is likely to result from simulating **many-body** systems.

Quantum Simulation: Details

Manipulations with the state of a QC can be split into several stages:

- **State preparation** (eigenstates, thermal states, etc.);
- Time evolution (if necessary);
- Measurement (of particular observables—not of all the amplitudes!);

Algorithms achieving these tasks are based on a variety of methods:

- Product formulas ($e^{A+B+C+\dots} \approx e^{\frac{A}{2}} e^{\frac{B}{2}} e^{\frac{C}{2}} \dots e^{\frac{C}{2}} e^{\frac{B}{2}} e^{\frac{A}{2}}$);
- Hybrid quantum-classical methods (variational, subspace, etc.);
- Quantum Signal Processing (implement $|\psi\rangle \mapsto f(H)|\psi\rangle$ for some H).

Preparing Hamiltonian Eigenstates

- **Variational:** $E_0 \approx \min_{\vec{\theta}} \langle \psi(\vec{\theta}) | H | \psi(\vec{\theta}) \rangle$, where $|\psi(\vec{\theta})\rangle = U(\vec{\theta}) |\psi_{\text{init}}\rangle$.
- Subspace methods (\approx Lanczos): $\{|\psi_{\text{init}}\rangle, H|\psi_{\text{init}}\rangle, H^2|\psi_{\text{init}}\rangle, \dots\}$
- **Adiabatic interaction turn-on.**

- **Filtering via Projection:**

$$P_0 |\psi_{\text{init}}\rangle = |\psi_0\rangle,$$

where

$$P_0 = |\psi_0\rangle\langle\psi_0|.$$

- **Thermalization.**

	$ \psi_{\text{init}}\rangle$ optimization	Provably converges	Near- term	Fault- tolerant
Variational	✓	✗	✓	✗
Subspace	✓	✓	✓	✓
Adiabatic	✗	?	✗	✓
Filtering	✓	✓	✗	✓
Thermalization	✗	?	✗	✓

2. How can QCs accelerate solving LF QCD?

Hamiltonian operator

- Assume the following form of the LF Hamiltonian operator:

$$H = \sum_{\xi_1, \xi_2, \dots}^N c_{\xi_1 \xi_2 \dots} q_{\xi_1} q_{\xi_2} ,$$

where

- $\xi = \{n, \vec{n}_\perp, s, \dots\}$ are the single-particle mode indices,
- q_ξ are boson/fermion creation/annihilation operators,
- $c_{\xi_1 \xi_2 \dots}$ are coefficients provided in the form of a lookup table or functional dependence (the latter may give extra advantage).

Exp vs Poly

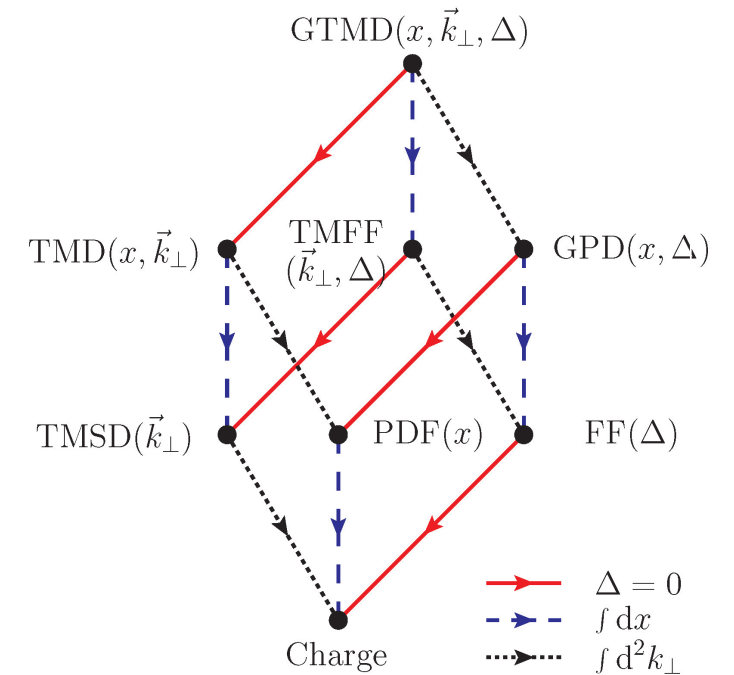
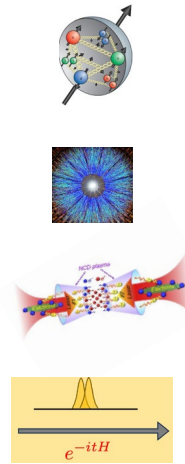
- The number of states of the form $|\xi_1^{w_1}, \xi_2^{w_2}, \dots, \xi_J^{w_J}\rangle$ scales as $\sim N^J$ for $J \ll N$, where N is the total number of modes.
- Classical simulation has $\exp(J)$ cost (memory, operations).
- Quantum simulation can be performed with $\text{poly}(J)$ cost:
 - number of qubits: those used for storing the state + “ancillas”,
 - number of quantum gates in a quantum circuit,
 - number of algorithm iterations.

Prerequisites

- Providing as input matrix elements of H in the basis of $|\xi_1^{W_1}, \xi_2^{W_2}, \dots\rangle$ eliminates the possibility of exponential advantage.
- For an efficient simulation, it is crucial that the input is provided in the form of $c_{\xi_1 \xi_2 \dots}$ in $H = \sum_{\xi_1, \xi_2, \dots}^N c_{\xi_1 \xi_2 \dots} q_{\xi_1} q_{\xi_2}$.
- Therefore, a (properly renormalized) Hamiltonian operator acting on multi-particle states is required.

Scenarios

- Spectra and static observables,
- Thermodynamics,
- Background fields,
- Dynamics.
- ...



Pasquini, Lorcé 2012

- Anything that can be calculated as $\langle \psi_2 | \hat{A} | \psi_1 \rangle$, where
 - \hat{A} is a known operator,
 - $|\psi_{1,2}\rangle$ are states which one knows* how to prepare.

3. Peculiarities of simulating LF QFT on QCs.

What makes LF QCD an appealing application of QC?

- Resources: (hadron structure from LF) $<$ (*ab initio* lattice dynamics):
 - Box size argument: fewer DOFs (N),
 - Harmonic resolution > 0 : fewer occupied modes (J),
 - Fewer fermionic DOFs per spinor.
- Similarity with non-relativistic many-body physics:
 - Trivial vacuum*,
 - Straightforward treatment of fermions and bosons,
 - Easy to get observables from the wavefunction.
- Simple structure of Hamiltonian operator coefficients*.

Toy example: ϕ^4 theory

Mapping onto qubits

- Binary encoding of an integer
(e.g., occupancy or momentum of a single bosonic DOF):
 $|0\rangle \mapsto |\dots 00\rangle$, $|1\rangle \mapsto |\dots 01\rangle$, $|2\rangle \mapsto |\dots 10\rangle$, $|3\rangle \mapsto |\dots 11\rangle$, etc.
- Fock state encoding in ϕ^4 theory in 1+1D:

Encoding	Qubit Representation of $ 1^5; 2^3; 3\rangle \equiv 1^5; 2^3; 3^1; 4^0; 5^0; \dots\rangle$	Qubit # Scaling	Qubit # Scaling for QCD
Direct	$ \underbrace{101}_{w_1=5}; \underbrace{011}_{w_2=3}; \underbrace{001}_{w_3=1}; \underbrace{000}_{w_4=0}; \underbrace{000}_{w_5=0}; \dots; \underbrace{000}_{w_K=0}\rangle$	$O(K \log K)$	$O(K \Lambda_{\perp}^2 \log(K + \Lambda_{\perp}))$
Compact	$ \{\underbrace{001}_{n=1}, \underbrace{101}_{w_1=5}\}; \{\underbrace{010}_{n=2}, \underbrace{011}_{w_2=3}\}; \{\underbrace{010}_{n=3}, \underbrace{001}_{w_3=1}\}; 000 \dots\rangle$	$O(\sqrt{K} \log K)$	$O(K \log(K + \Lambda_{\perp}))$

- Not all the simulation techniques are compatible with Compact.

Toy example: ϕ^4 theory

Measurement

- PDF measurement operator:

$$f(n/K) = \langle \psi | a_n^\dagger a_n | \psi \rangle .$$

- Corresponds to measuring the expectation value of the $a_n^\dagger a_n$ operator for the corresponding qubit register $| \dots ; \underbrace{\dots}_n ; \dots \rangle$.
- Generalizable to GPDs, GTMDs, etc.

What are the complications in QSim of LF QCD?

- What defines simulation cost:
 - Number of elementary operator terms in the Hamiltonian,
 - Locality of Hamiltonian,
 - Hamiltonian norm.
- Number of terms scales as N^3 or N^4 as compared to N in ET LGT.
- LF Hamiltonian is non-local
(which is true for almost any basis other than real space lattice).
- A **tight** bound on $||H_{LF}||$ is not known (to me).

Wavelets!



4. Past, Current, & Future work.

Existing work

- State preparation, fault-tolerant algorithms:
 - [2002.04016](#), [2105.10941](#), [2211.07826](#).
- State preparation, near-term algorithms (variational):
 - [2009.07885](#), [2011.13443](#).
- Scattering with background field:
 - [2307.09987](#), [2310.13742](#), [2404.00819](#), [2205.07902](#), [2311.18209](#). SFQED!
- Scattering of dynamical observables:
 - [2310.13742](#), [2401.04496](#).

Ongoing work: resource estimation for state preparation

- Motivation: estimating the size of a QC required for *ab initio* state preparation of the Sexaquark (*uuddss*) state, a SM DM candidate.
- DLCQ: $H = \dots + \sum_{\xi_1 \xi_2 \xi_3 \xi_4}^N h_{\xi_1 \xi_2 \xi_3 \xi_4} b_{\xi_1}^\dagger b_{\xi_2} a_{\xi_3}^\dagger a_{\xi_4} \delta(p_1 + p_3 - p_3 - p_4) + \dots$
- *State preparation algorithm* implements $|\psi\rangle \rightarrow P_0 |\psi\rangle$.
 - P_0 is constructed using calls to the *Block Encoding* subroutine U_H .
 - U_H provides info about H by implementing $|\psi\rangle \rightarrow H|\psi\rangle$.

- State prep. cost = $\underbrace{\text{Cost}(U_H)}_{\sim N \sqrt{}} \times \underbrace{\text{Cost}(\|H\|)}_{\sim N^3 \text{???}} \times \frac{1}{\Delta} \times \frac{1}{\underbrace{\gamma}_{\text{???}}} \times \log \frac{1}{\epsilon}$.

$$\Delta = E_1 - E_0 \ \& \ \gamma = |\langle \psi_{\text{init}} | \psi_0 \rangle| \ \& \ \epsilon = 1 - |\langle \psi_{\text{init}} | \psi_0 \rangle|^2$$

Future work

- Deriving renormalized Hamiltonian operators.
- Reference state choice: utilizing best classically obtained solutions.
- Systematic analysis: uncertainties from truncation, $||H||$, etc.
- Applying state-of-the-art simulation algorithms, e.g.:
 - Subspace methods; state preparation via thermalization; etc.
- Transverse LF lattice:
 - Taking advantage of locality in the transverse directions,
 - Suitable for scattering of dynamical composite particles?

Thank You!