

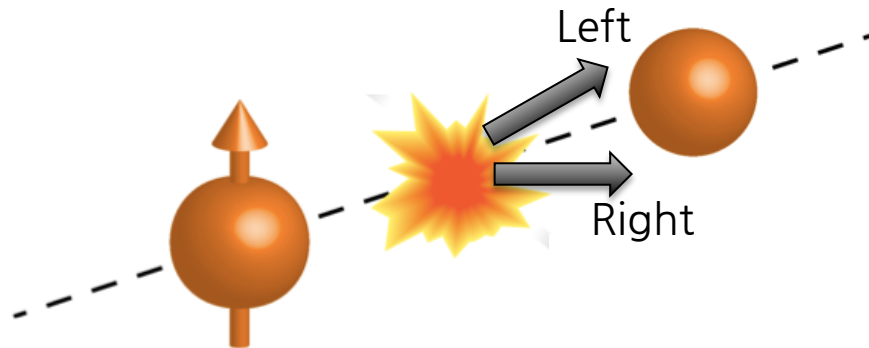
# Transverse single spin asymmetry measurement at the RHICf experiment

Minho Kim (RIKEN)  
on behalf of the RHICf collaboration

December 1, 2023

2nd APCTP Workshop on the Physics of EIC:  
ePIC Physics and Detectors

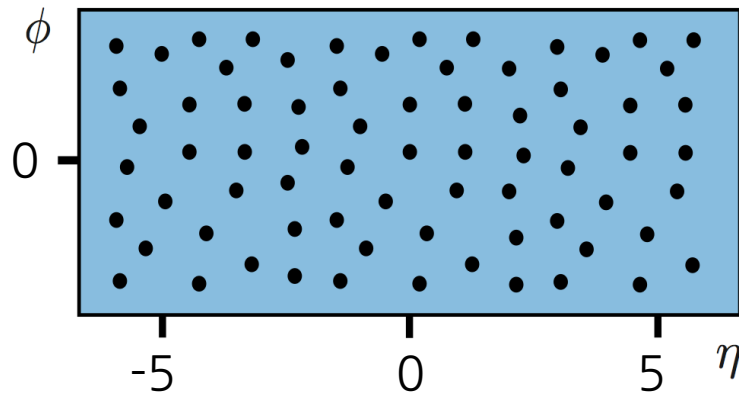
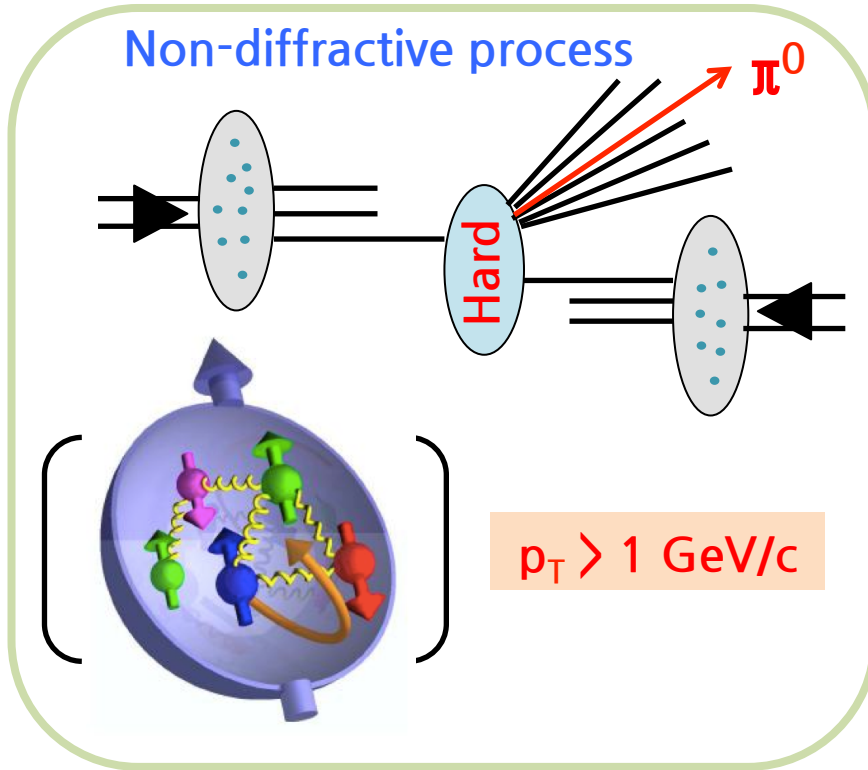
# Transverse single spin asymmetry ( $A_N$ )



$$A_N = \frac{\sigma_L^\uparrow - \sigma_R^\uparrow}{\sigma_L^\uparrow + \sigma_R^\uparrow} = \frac{\sigma_L^\uparrow - \sigma_L^\downarrow}{\sigma_L^\uparrow + \sigma_L^\downarrow}$$

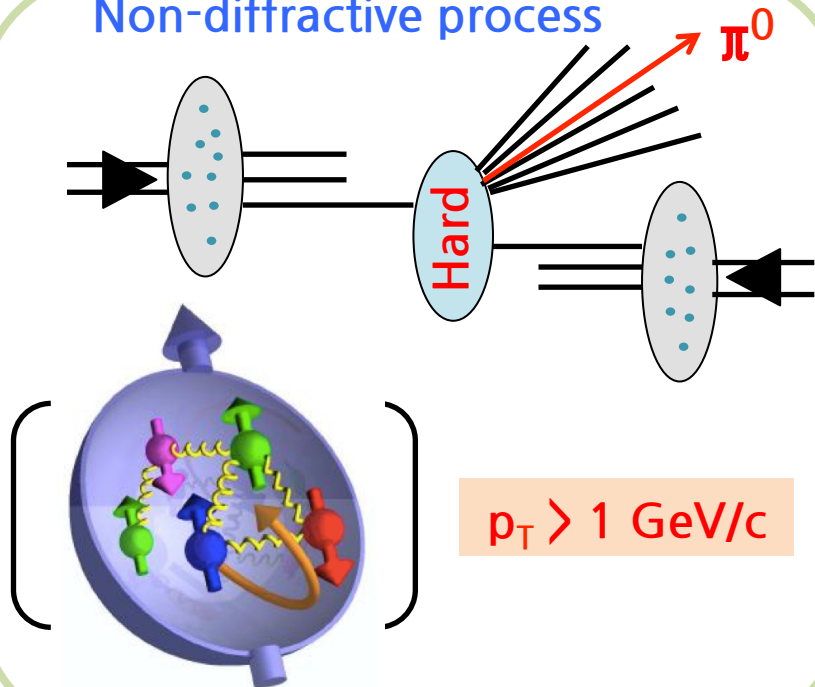
- In polarized p+p collisions, the  $A_N$  is defined by a left-right cross section asymmetry of a specific particle or event.
- Due to the rotational invariance, the left-right asymmetry can also be defined by the spin up-down asymmetry.
- $A_N$ s of the very forward ( $6 < \eta$ ) particles enable us to study the spin-involved diffractive particle production mechanism.

# Non-diffractive Vs. Diffractive process

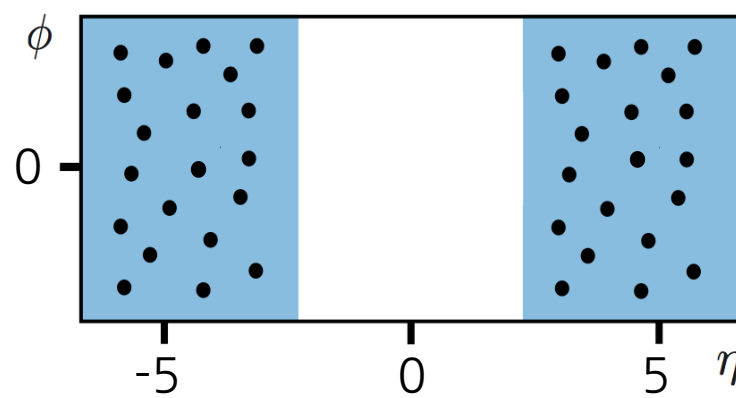
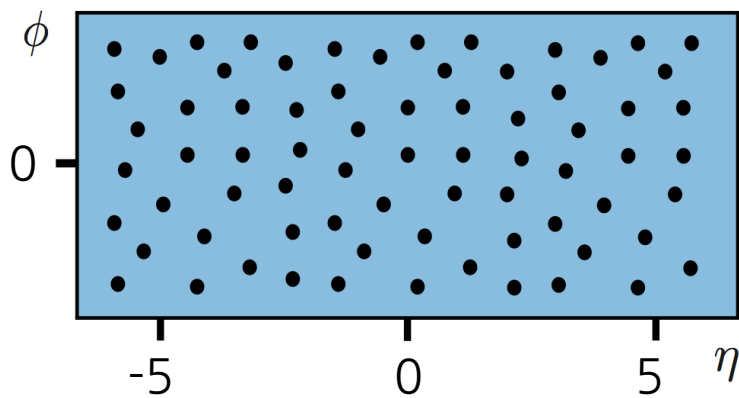
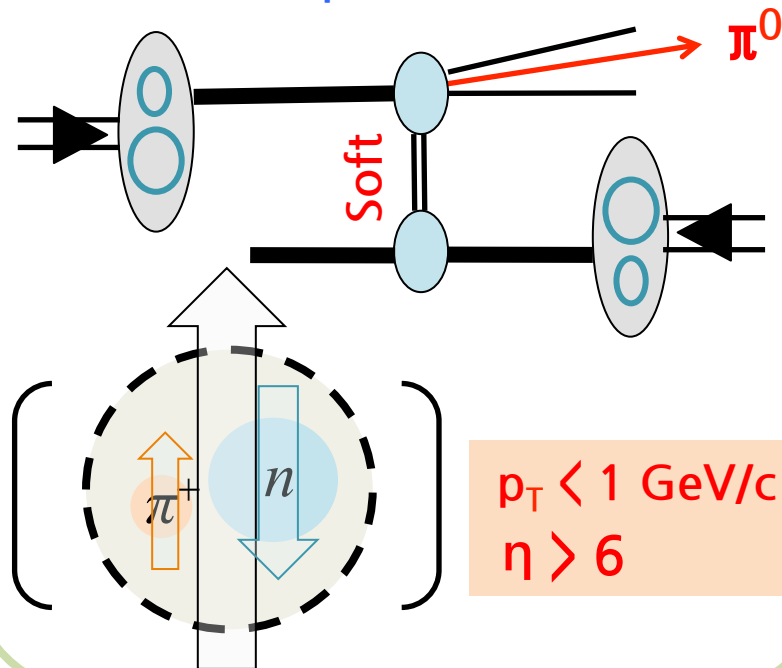


# Non-diffractive Vs. Diffractive process

Non-diffractive process

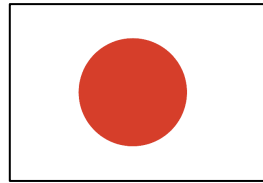
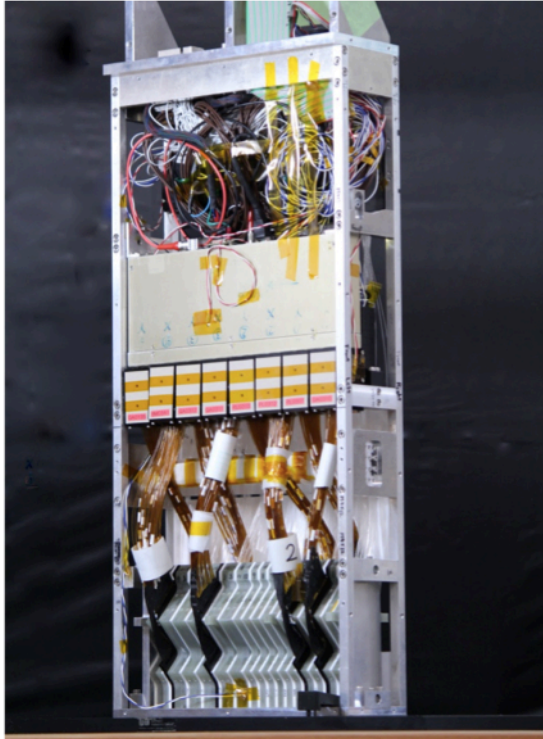


Diffractive process



# RHICf collaboration

## RHICf detector



**RIKEN** (Y. Goto, I. Nakagawa, R. Seidl, M. H. Kim), **Nagoya Univ.** (Y. Itow, H. Menjo, K. Sato, K. Ohashi), **Univ. of Tokyo** (T. Sako), **JAEA** (K. Tanida), **Waseda Univ.** (S. Torii), **Shibaura Inst. of Tech.** (K. Kasahara), **Tokushima Univ.** (N. Sakurai)



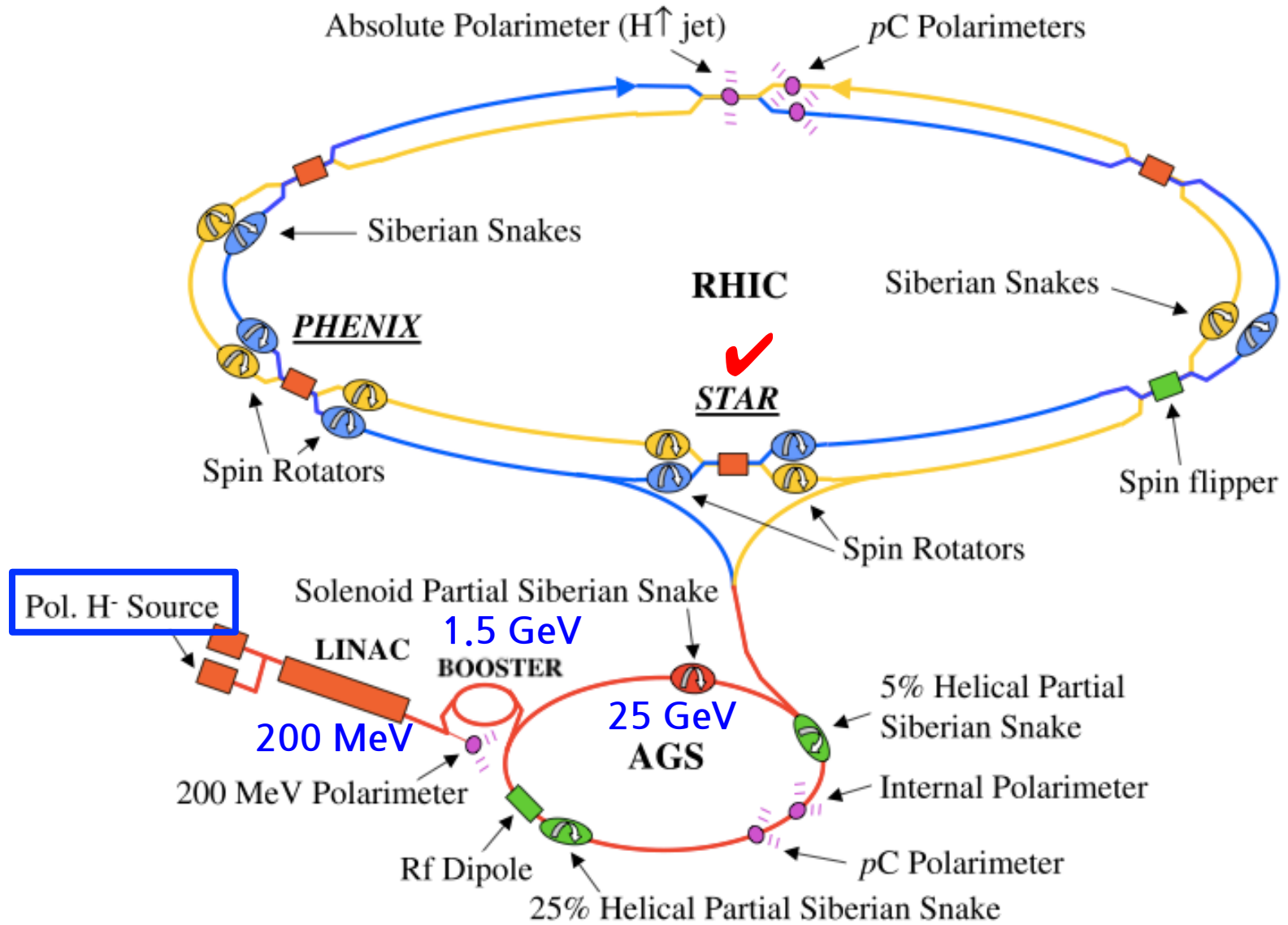
**Korea Univ.** (B. Hong), **Sejong Univ.** (Y. Kim, S. Oh, S. H. Lee)



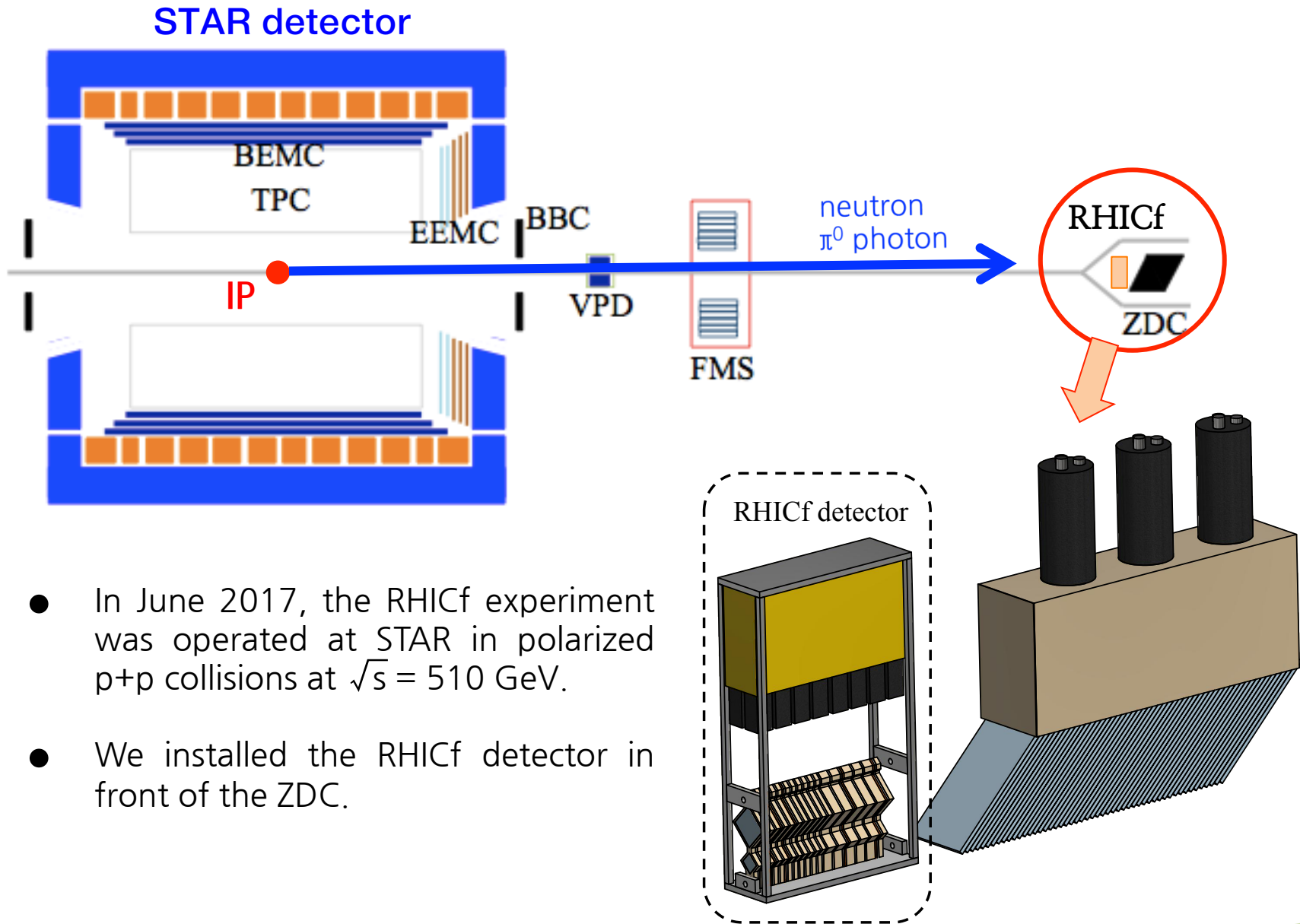
**INFN** (O. Adriani, E. Berti, L. Bonechi, R. D'Alessandro, A. Tricomi)

- **$A_N$**  measurement to study the spin-involved diffractive particle production mechanism.
- **Cross section** measurement to study the origin of the ultra-high energy cosmic ray.

# Relativistic Heavy Ion Collider (RHIC)

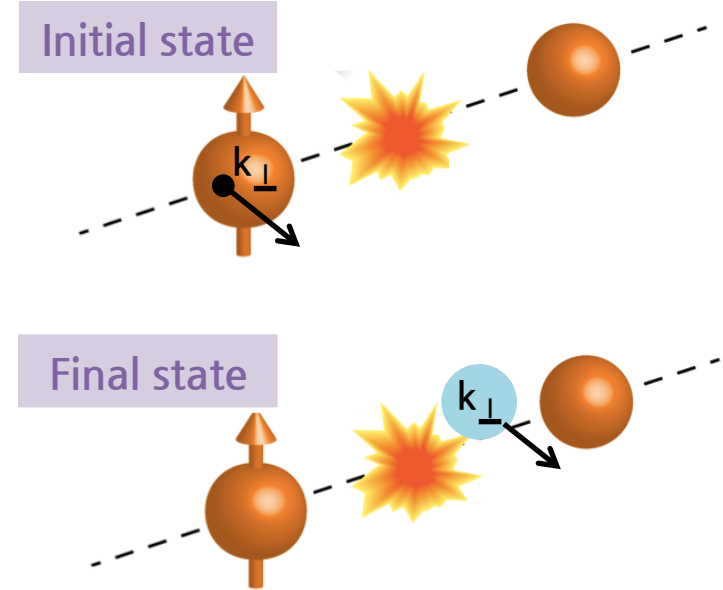
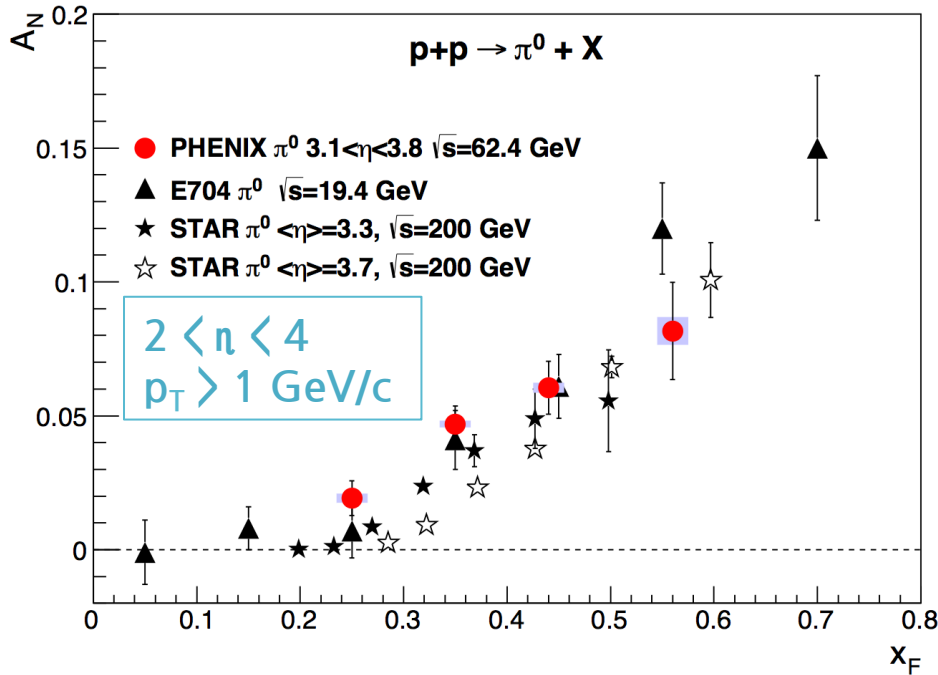


# RHIC forward (RHICf) experiment



# $A_N$ of forward $\pi^0$

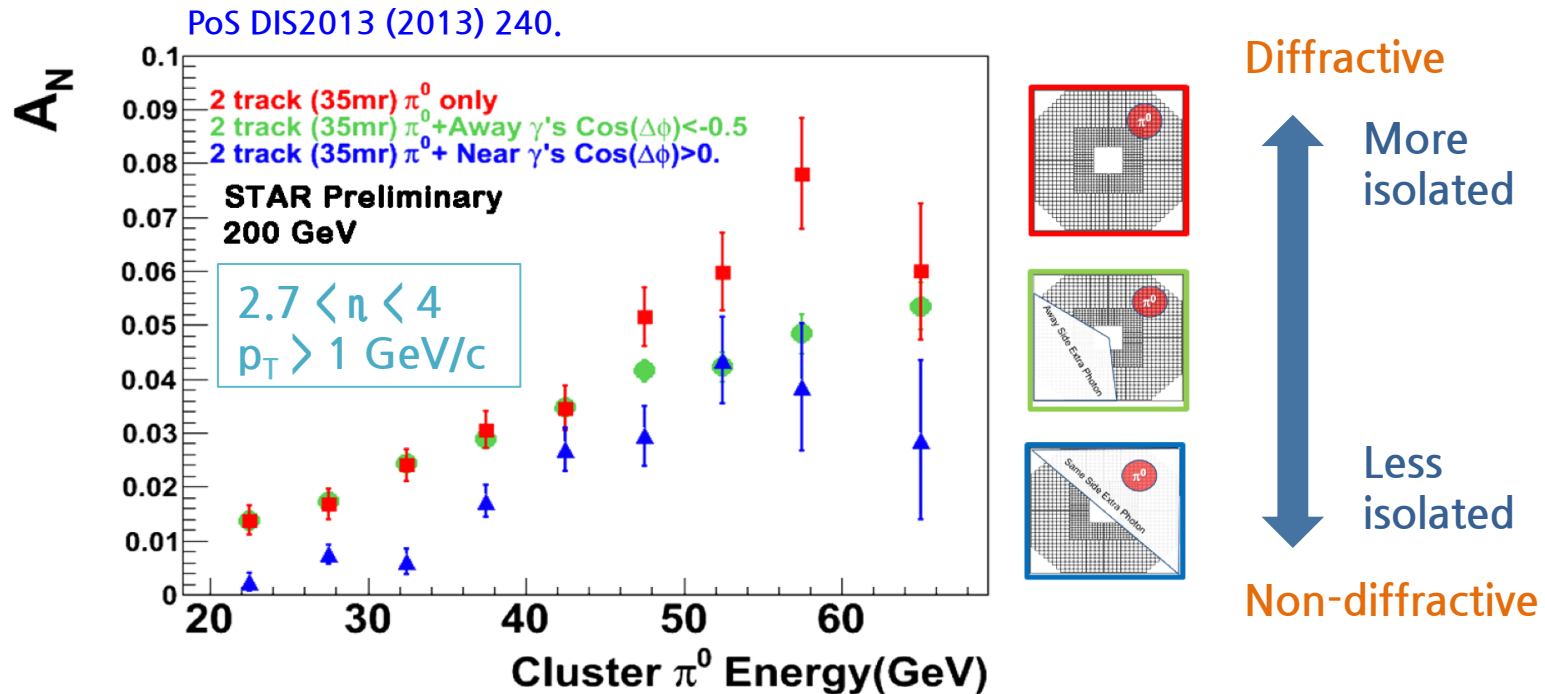
PRD 90, 012006 (2014).



- Before the RHICf experiment,  $A_N$  for  $\pi^0$  production has been mainly measured in the forward ( $2 < \eta < 4$ ) kinematic region.
- The non-zero  $A_N$  has been interpreted based on quarks and gluons' degrees of freedom theoretically.

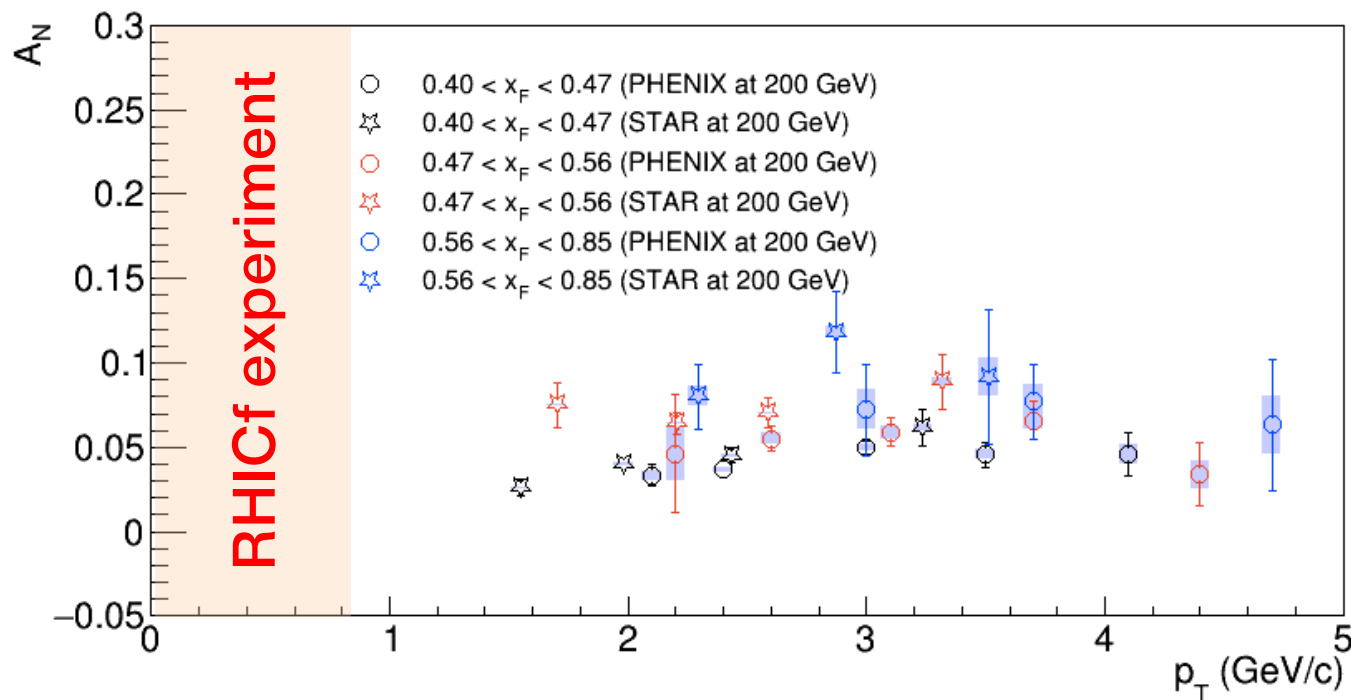


# $A_N$ of forward isolated $\pi^0$



- Larger  $A_N$  was observed by more isolated  $\pi^0$  than less isolated one.
- Diffractive process may have a finite contribution to the  $A_N$  for  $\pi^0$  production as well as the non-diffractive one.

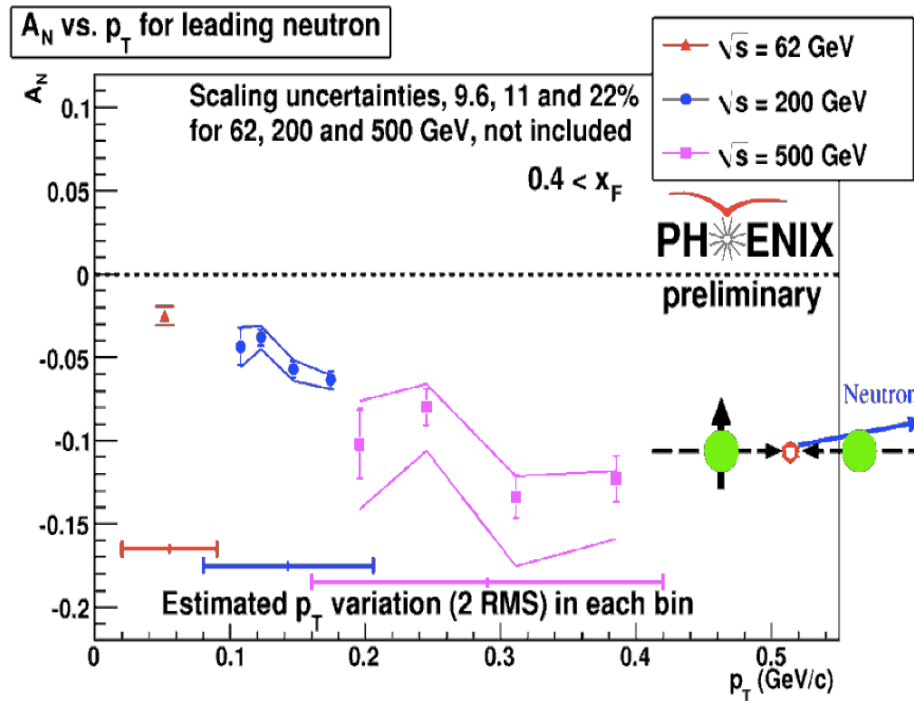
# $\pi^0$ measurement at RHICf



- No experiment has measured the  $\pi^0$  in detail in the range of  $p_T < 1$  GeV/c.
- In order to study a possible diffractive contribution to the  $\pi^0$   $A_N$ , the RHICf experiment firstly measured the  $A_N$  for very forward  $\pi^0$  production.

# $A_N$ of very forward neutron

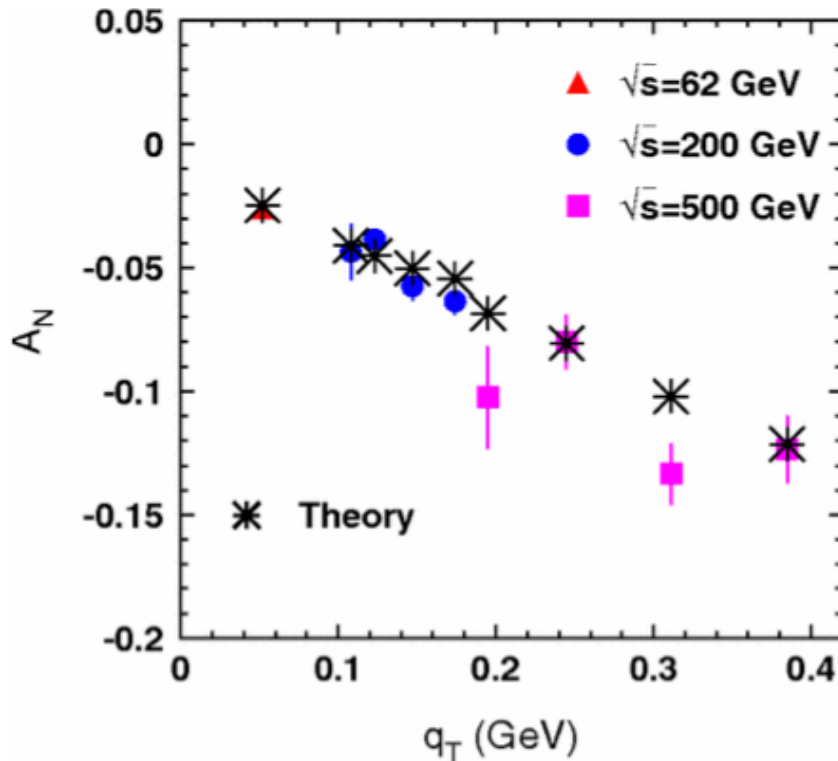
J. Phys. Conf. Ser. 295 (2011) 012097.



- Non-zero  $A_N$  for very forward neutron production was first observed by an experiment called IP12. [PLB 650, 325 \(2007\)](#).
- Afterwards, the PHENIX measured the neutron  $A_N$  as a function of  $p_T$  with three different collision energies.
- The measurement results showed a possible  $p_T$  dependence of the neutron  $A_N$ .

# Theoretical model

PRD 84, 114012 (2011).



$$\begin{aligned}
 A_N &= \frac{d\sigma^\uparrow - d\sigma^\downarrow}{d\sigma^\uparrow + d\sigma^\downarrow} \\
 &= \frac{\sum_X |\langle cX|T|\uparrow\rangle|^2 - \sum_X |\langle cX|T|\downarrow\rangle|^2}{\sum_X |\langle cX|T|\uparrow\rangle|^2 + \sum_X |\langle cX|T|\downarrow\rangle|^2} \\
 &= \frac{-2\text{Im} \sum_X \langle cX|T|-\rangle \langle +|T^\dagger|cX\rangle}{\sum_X |\langle cX|T|+\rangle|^2 + \sum_X |\langle cX|T|-\rangle|^2}
 \end{aligned}$$

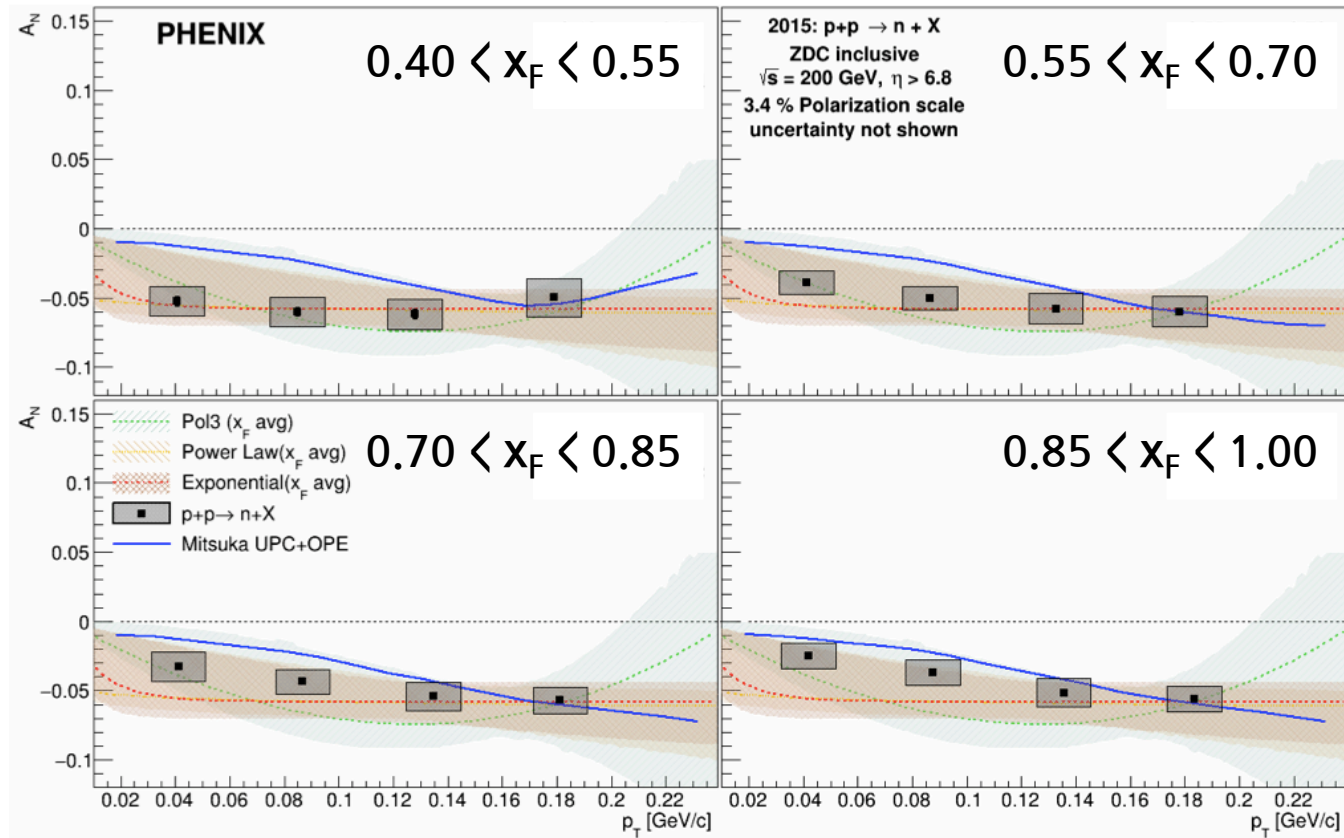
$\pi$  exchange: spin flip

$a_1$  exchange: spin non-flip

- Neutron  $A_N$  was explained by an interference between spin flip and spin non-flip amplitudes with non-zero phase shift.
- The  $\pi$  and  $a_1$  exchange model showed that the neutron  $A_N$  increased in magnitude with increasing  $p_T$  with little  $\sqrt{s}$  dependence.

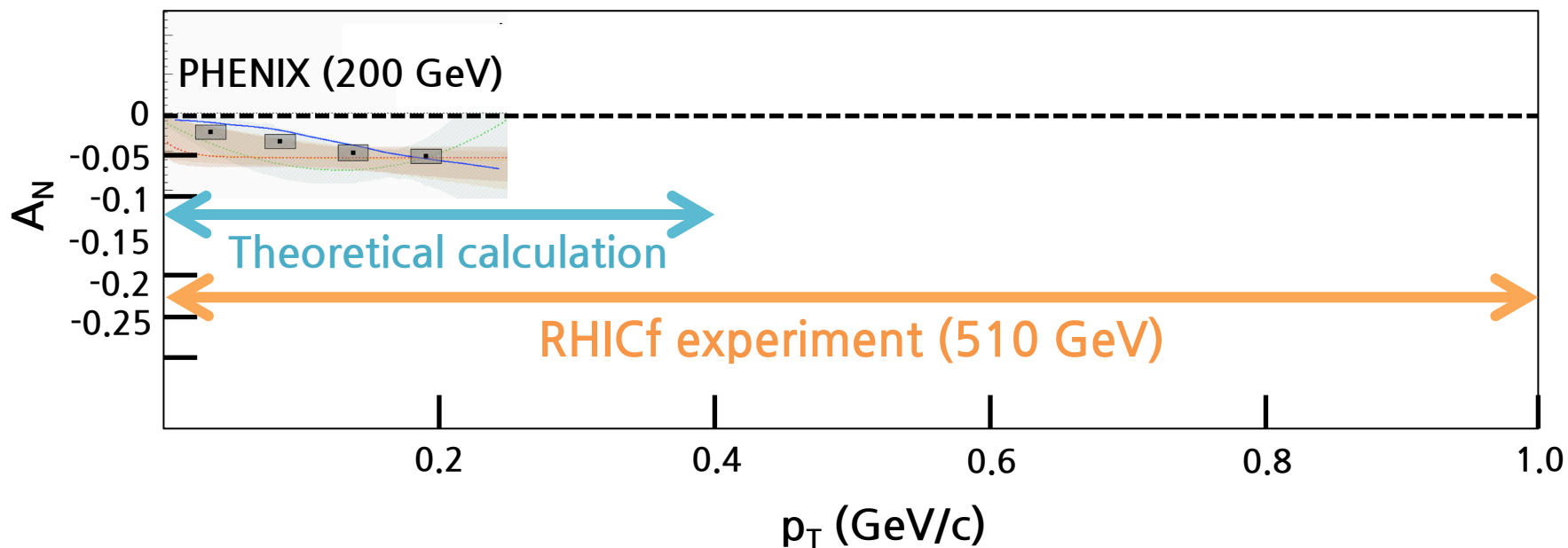
# Unfolded neutron $A_N$ at PHENIX

PRD 105, 032004 (2022).



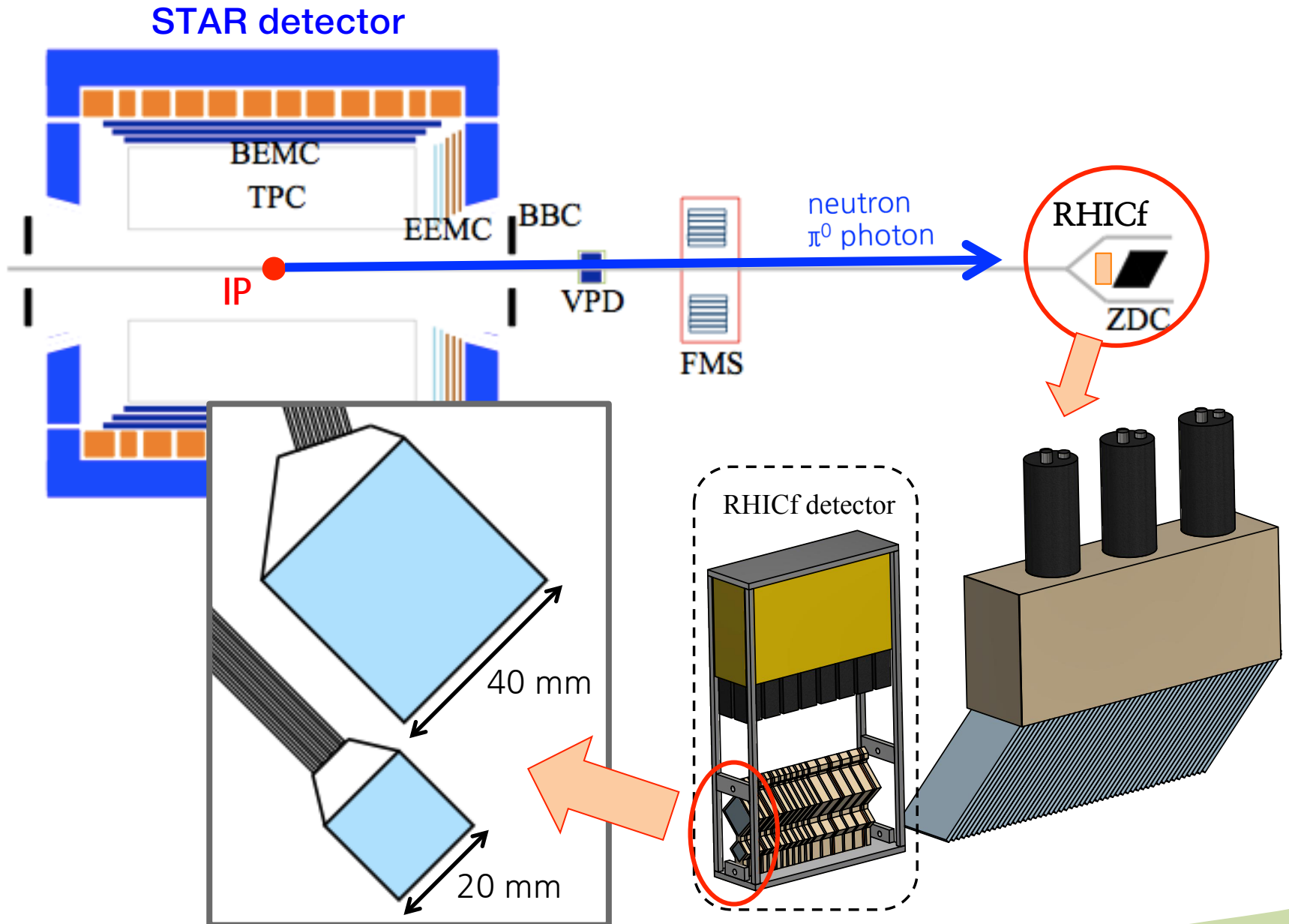
- Recently,  $p_T$  dependence of the PHENIX neutron  $A_N$  at  $\sqrt{s} = 200$  GeV was obtained by unfolding the data.
- The unfolded data showed the same tendency with the model calculations.

# Neutron measurement at RHICf

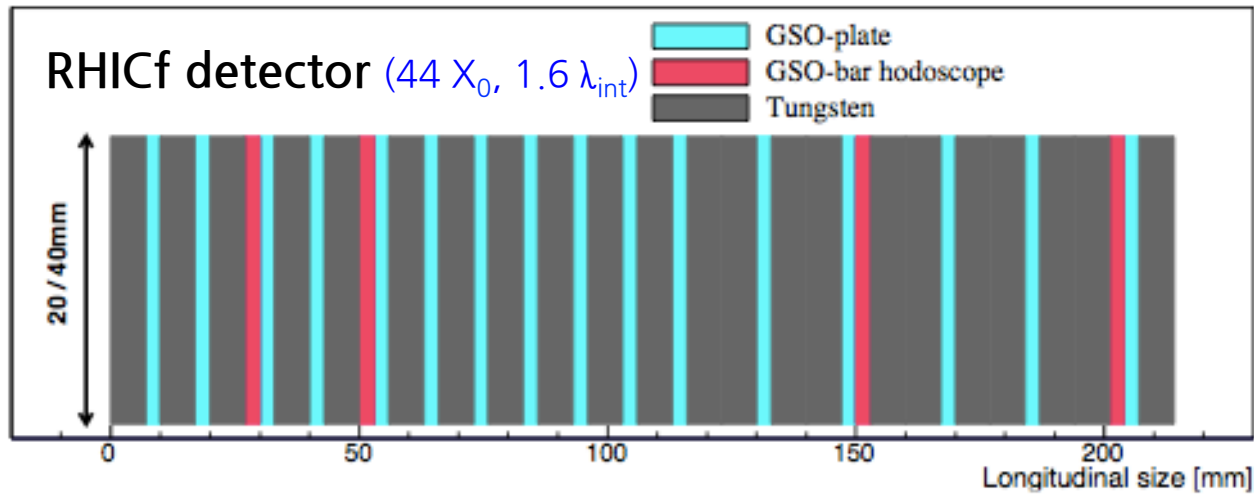


- RHICf the experiment has extended the previous measurements up to 1 GeV/c to study the kinematic dependence of the neutron  $A_N$  in more detail.
- RHICf detector has **one order of better position resolution** (1 cm  $\rightarrow$  1 mm).
- We can also study the  **$\sqrt{s}$  dependence** of the neutron  $A_N$  by comparing the RHICf data with that of the PHENIX.

# RHIC forward (RHICf) experiment

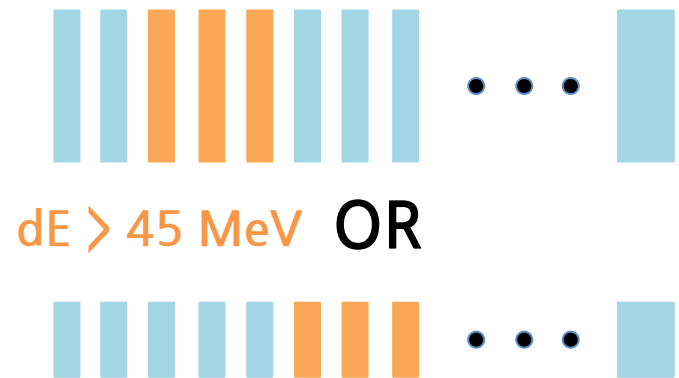


# RHICf detector & neutron measurement



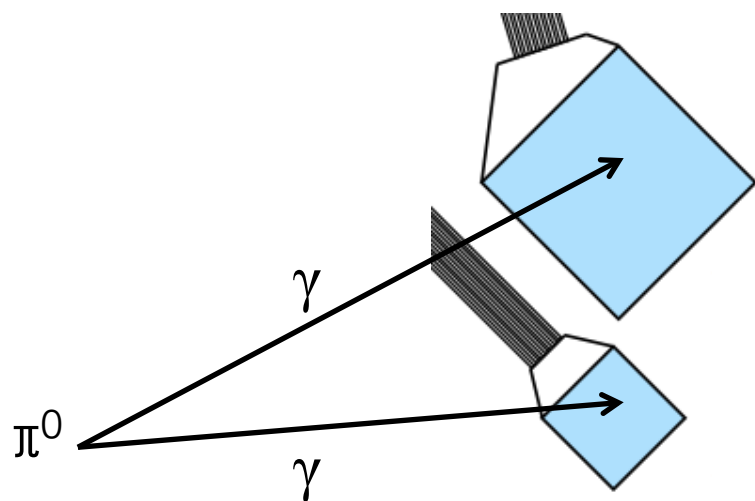
## Shower trigger

- Each tower is composed of 17 layers of tungsten absorbers, 16 layers of GSO plate, and 4 layers of GSO bar hodoscope.
- Shower trigger is operated when the energy deposits of any three successive GSO plate layers are larger than 45 MeV.

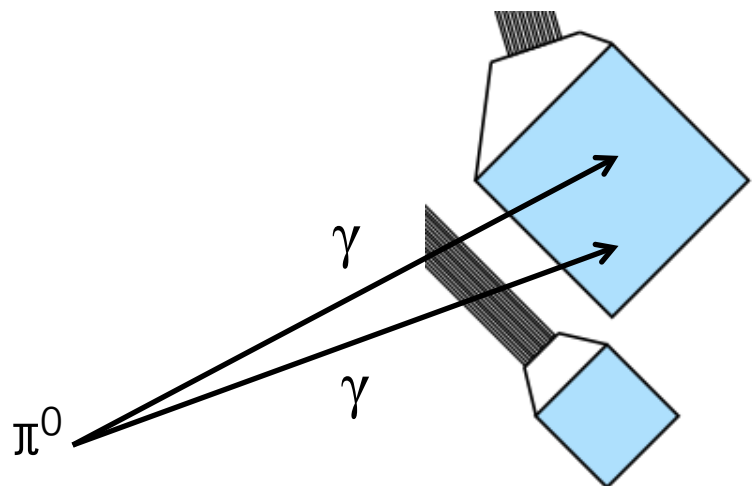
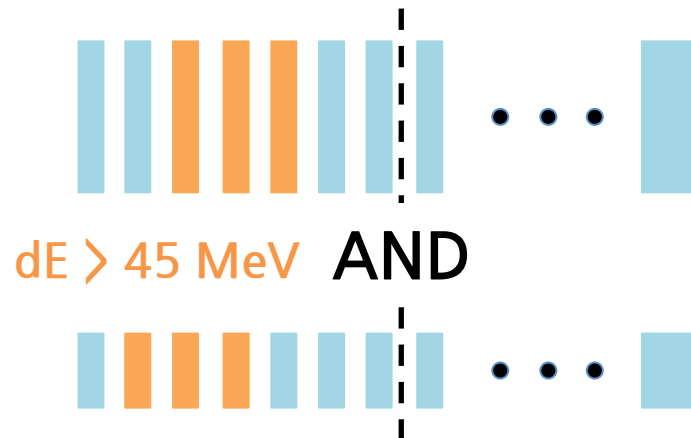




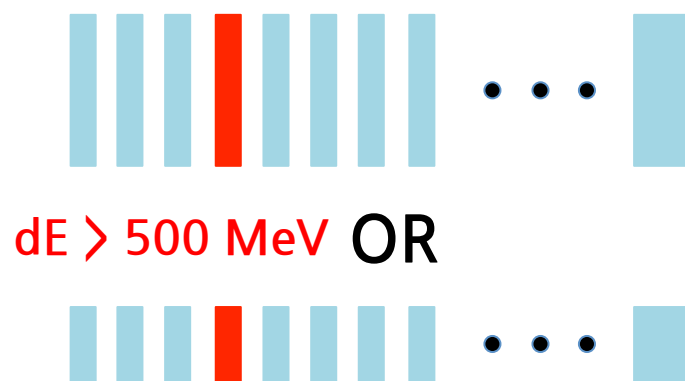
# $\pi^0$ measurement



$\pi^0$ -enhanced trigger

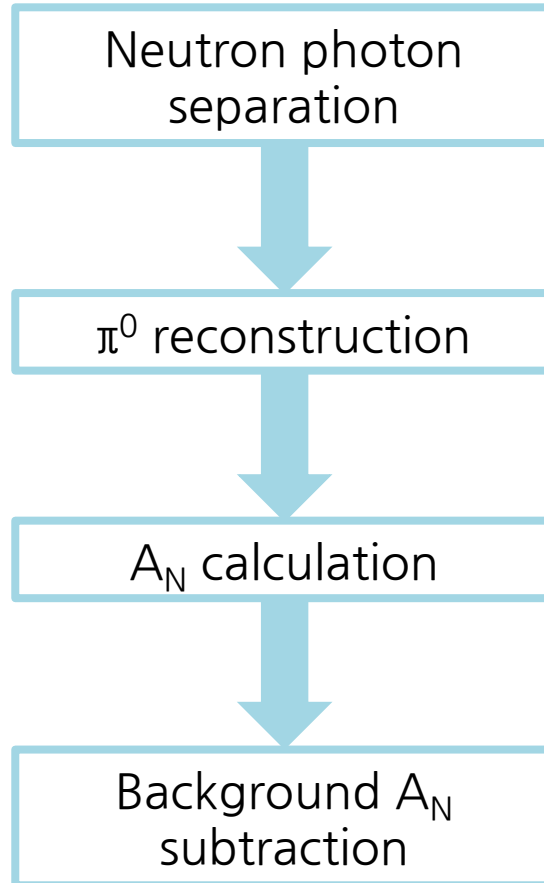


High EM trigger

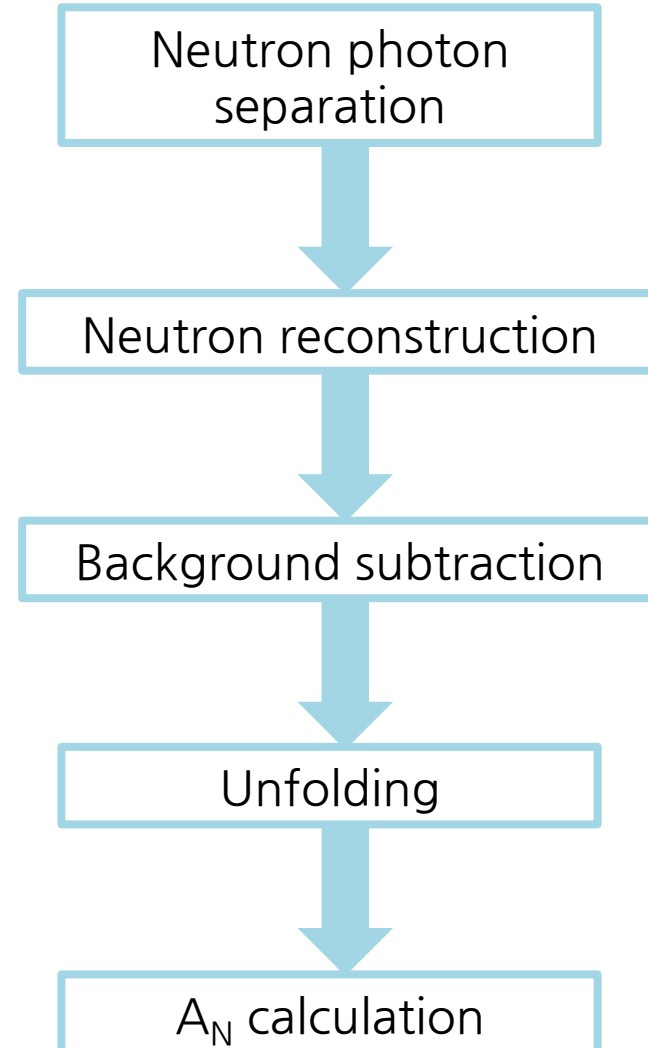


# Analysis flow

$\pi^0$

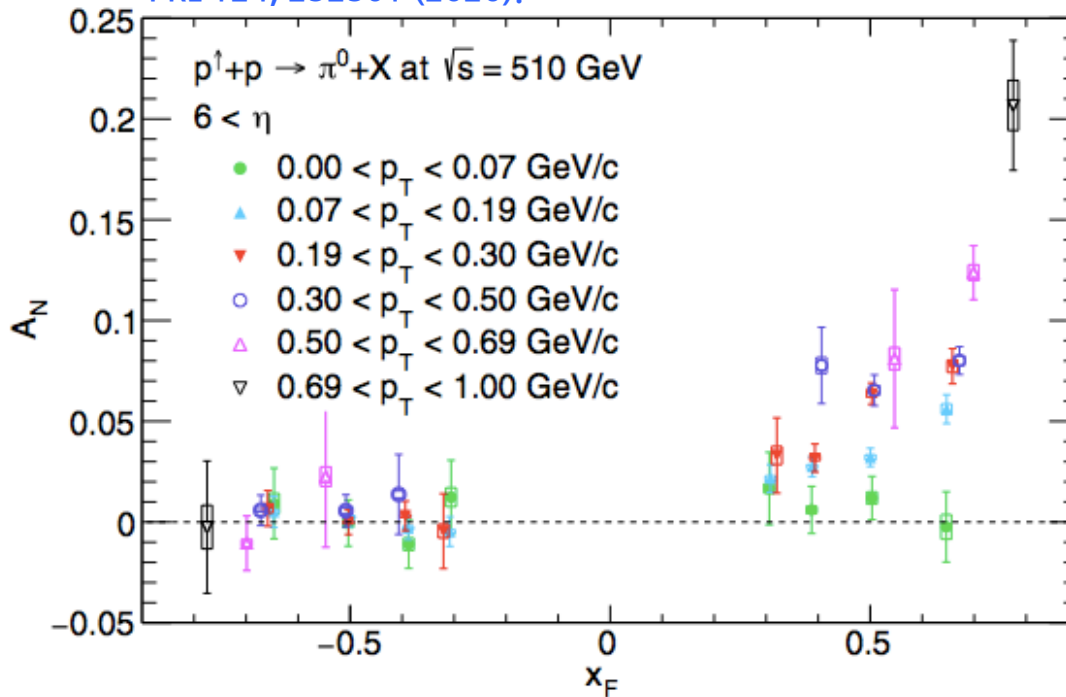


Neutron

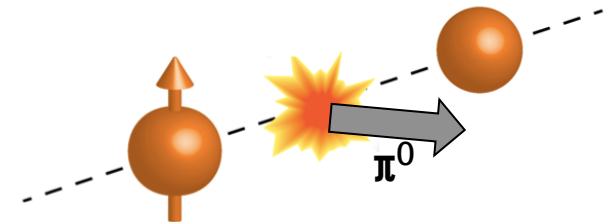


# $A_N$ of very forward $\pi^0$

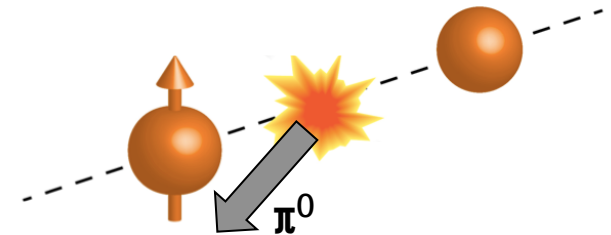
PRL 124, 252501 (2020).



i)  $x_F > 0$  (very forward  $A_N$ )

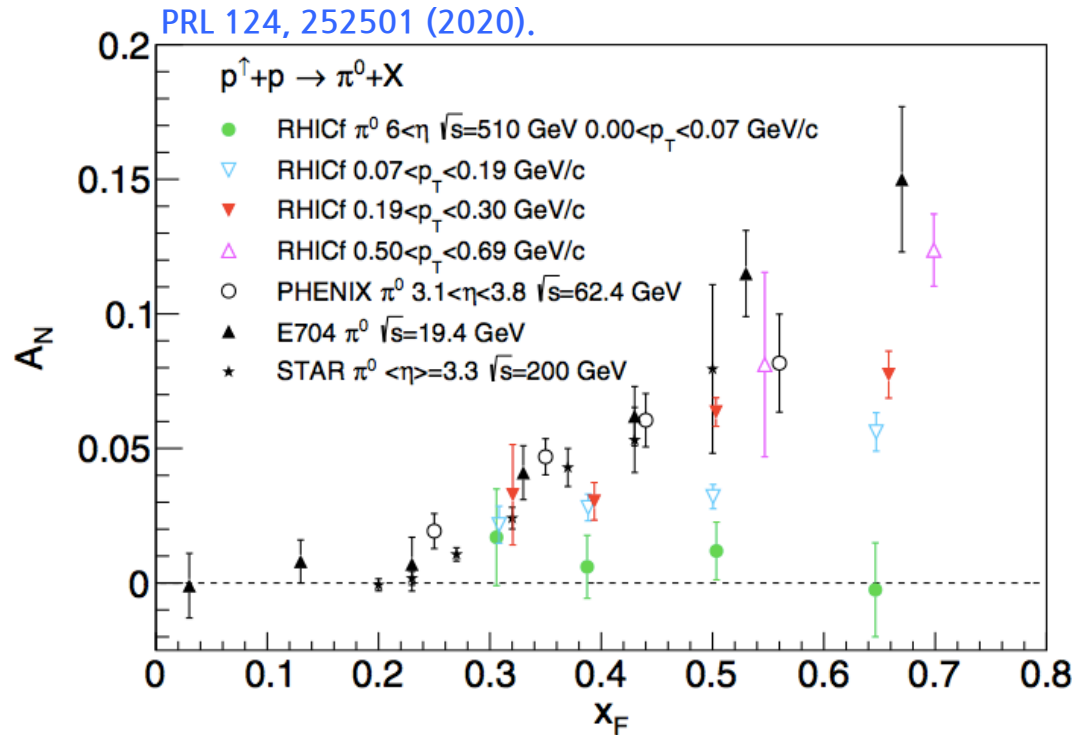


ii)  $x_F < 0$  (very backward  $A_N$ )



- At very low  $p_T < 0.07$  GeV/c, the  $A_N$ s are consistent with zero.
- However, as  $p_T$  increases, the  $A_N$ s also increase as a function of  $x_F$  even though it is expected that the very forward  $\pi^0$  comes from the diffractive process.

# Comparison with the previous measurements

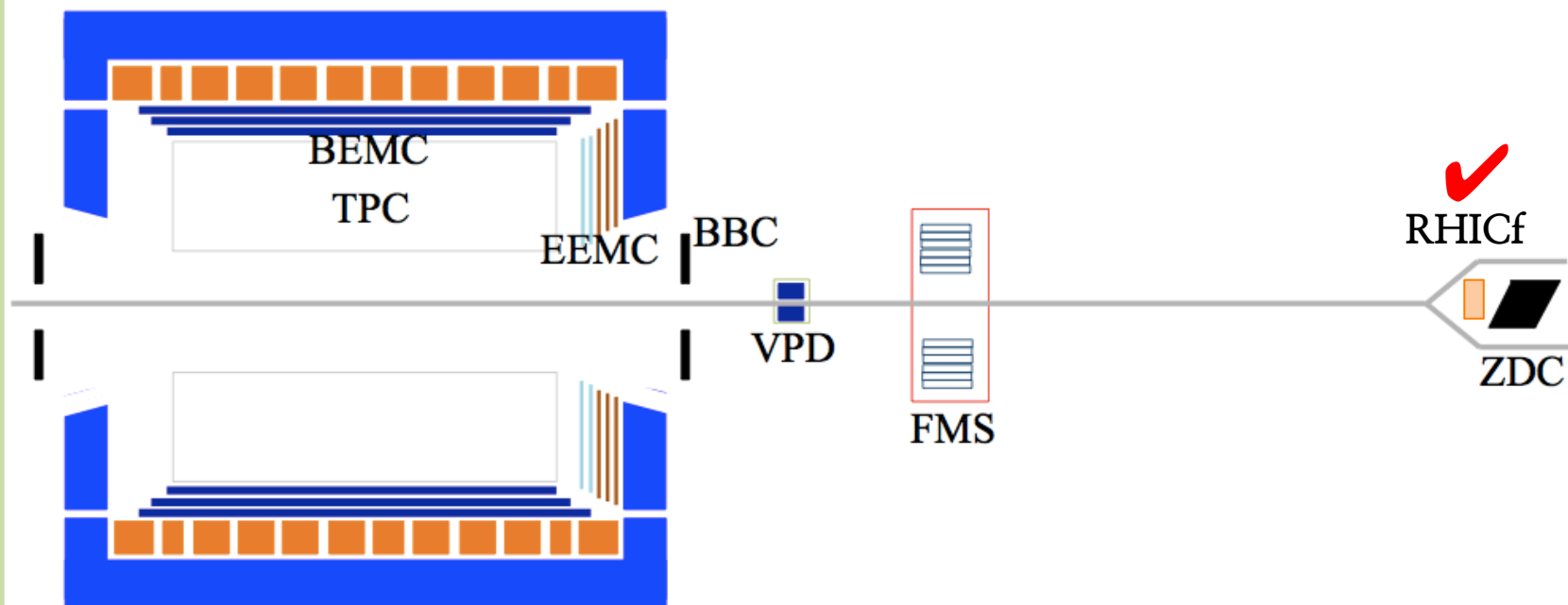


- One very interesting point is that the  $A_N$  of very forward  $\pi^0$  seems to be comparable with that of forward  $\pi^0$ .
- They may share a common underlying production mechanism or have their own ones.

# RHICf-STAR combined analysis

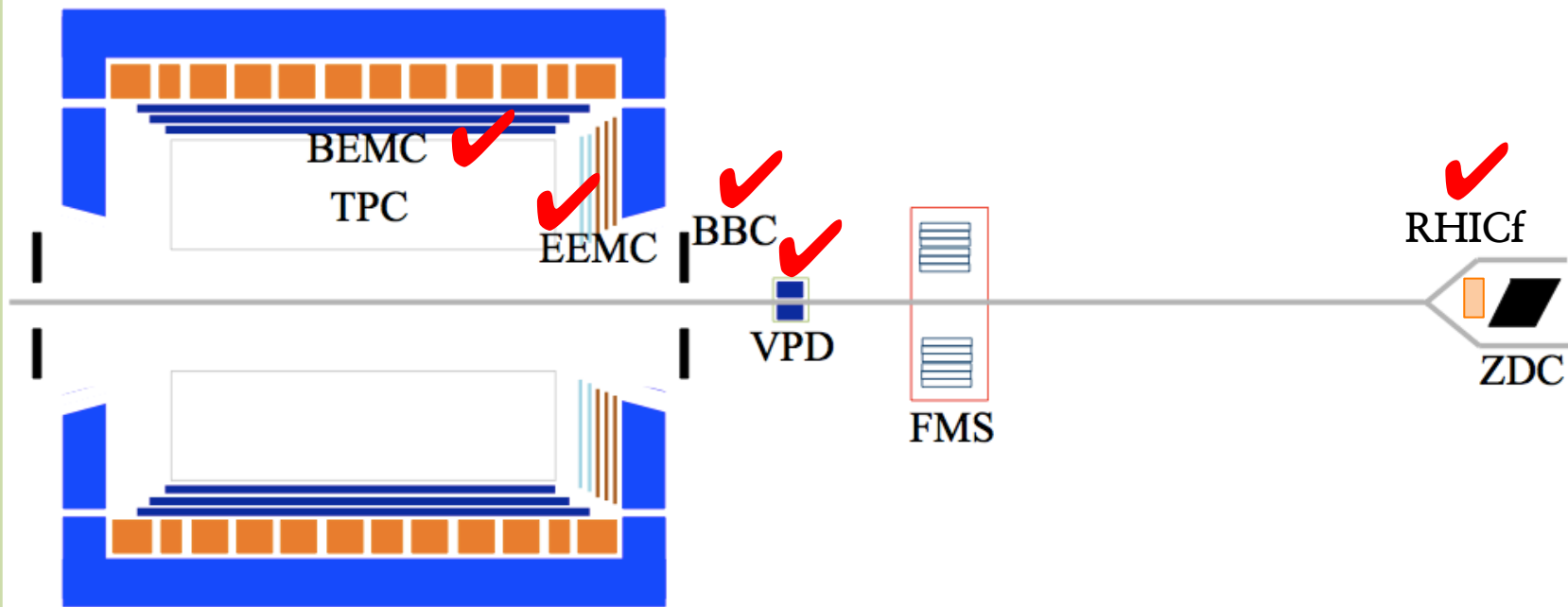
## STAR detector

Central detectors

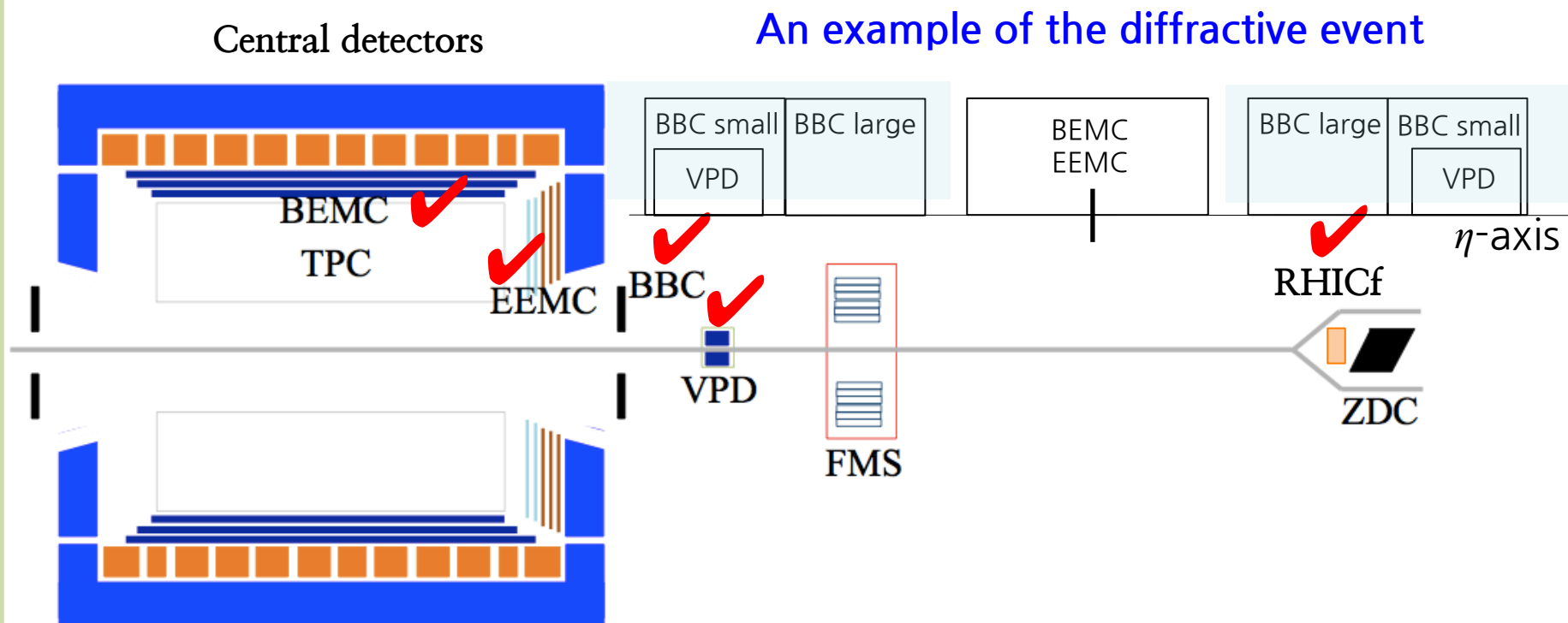


# RHICf-STAR combined analysis

Central detectors

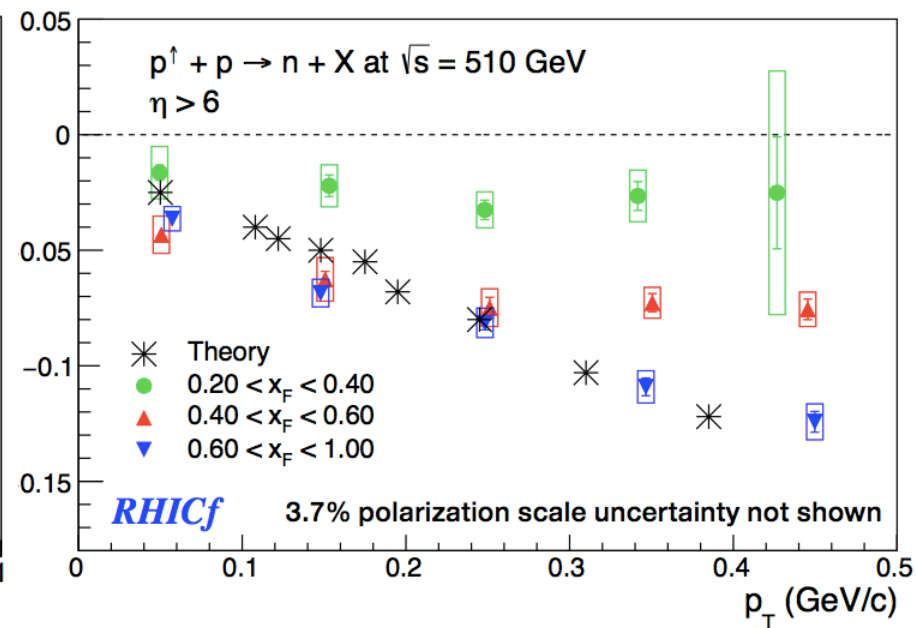
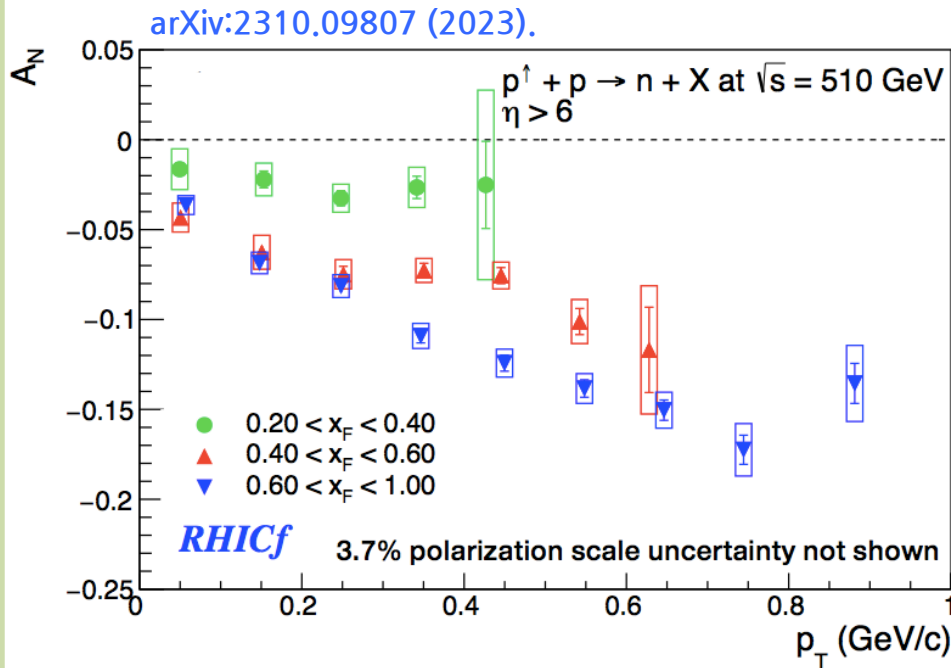


# RHICf-STAR combined analysis



- We're extending the RHICf standalone analysis to a combined analysis with STAR detectors to study the origin of the RHICf  $\pi^0$  results.
- If we use the STAR detectors, we could identify the diffractive and non-diffractive events.

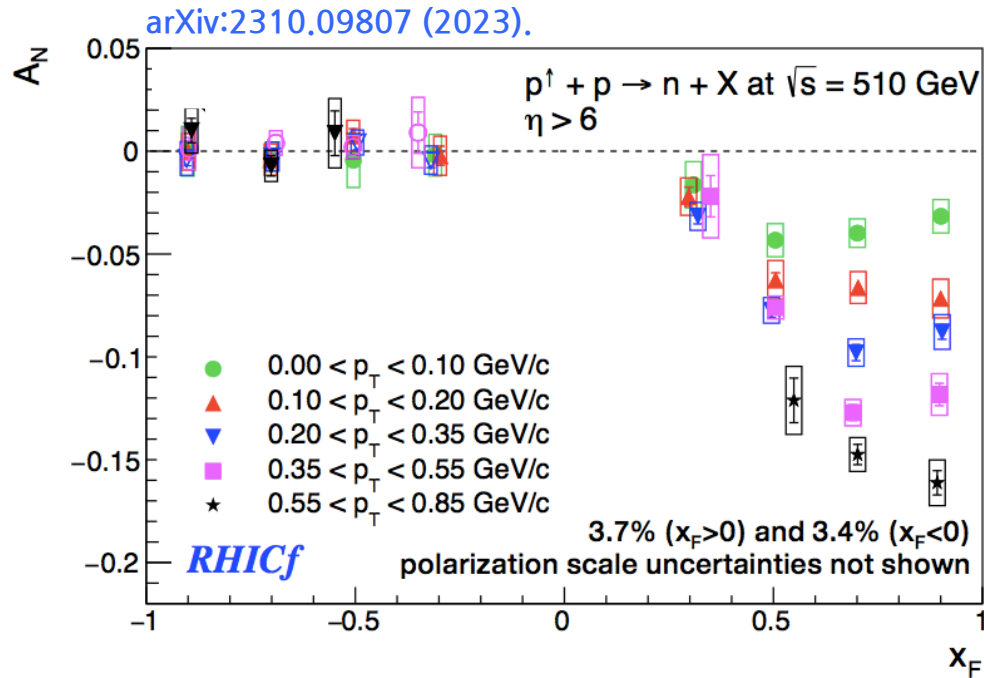
# $A_N$ of very forward neutron



- In the low  $x_F$  range, the neutron  $A_N$  reaches a plateau at low  $p_T$ .
- In the high  $x_F$  range, the  $A_N$  doesn't seem to reach the plateau yet, but we can confirm that the  $A_N$  explicitly increases in magnitude with  $p_T$ .
- The current theoretical calculation only reproduces the  $A_N$  in the high  $x_F$  range.



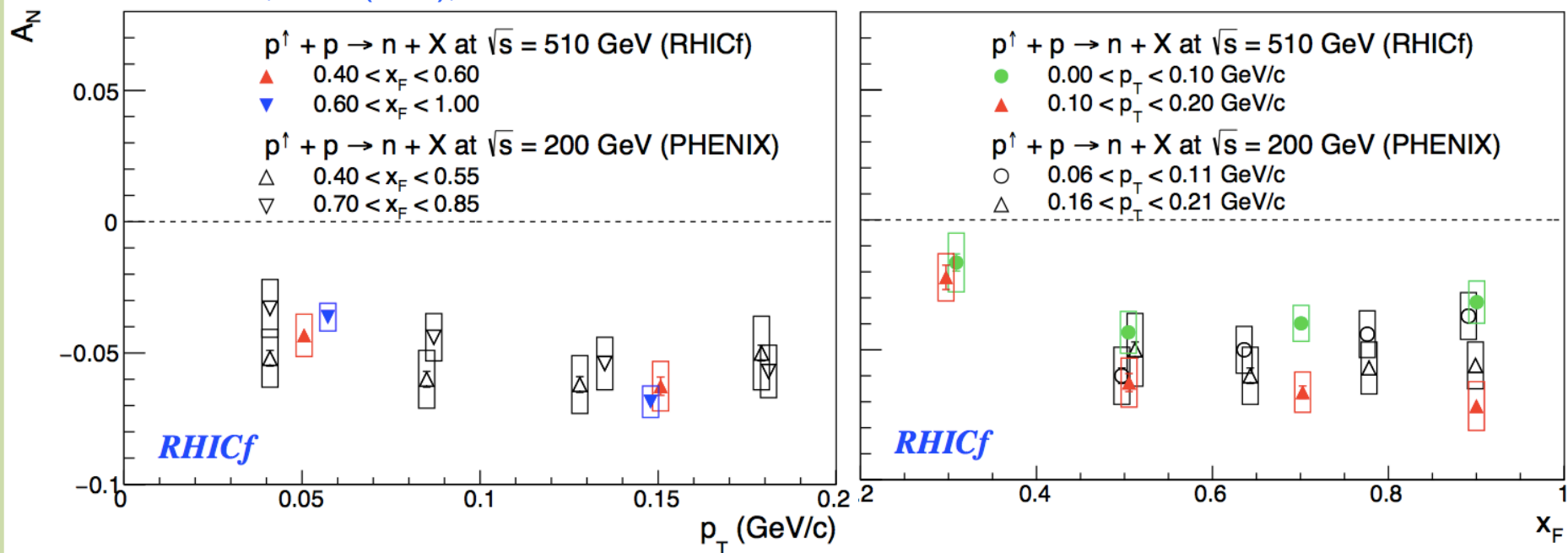
# $A_N$ of very forward neutron



- In the low  $p_T$  range, the  $A_N$  reaches a plateau at low  $x_F$  with little  $x_F$  dependence.
- In the high  $p_T$  range, the  $A_N$  reaches a higher plateau at higher  $x_F$  with **a clear  $x_F$  dependence**.
- More comprehensive theoretical consideration is necessary to explain the present results.

# Comparison with the PHENIX measurements

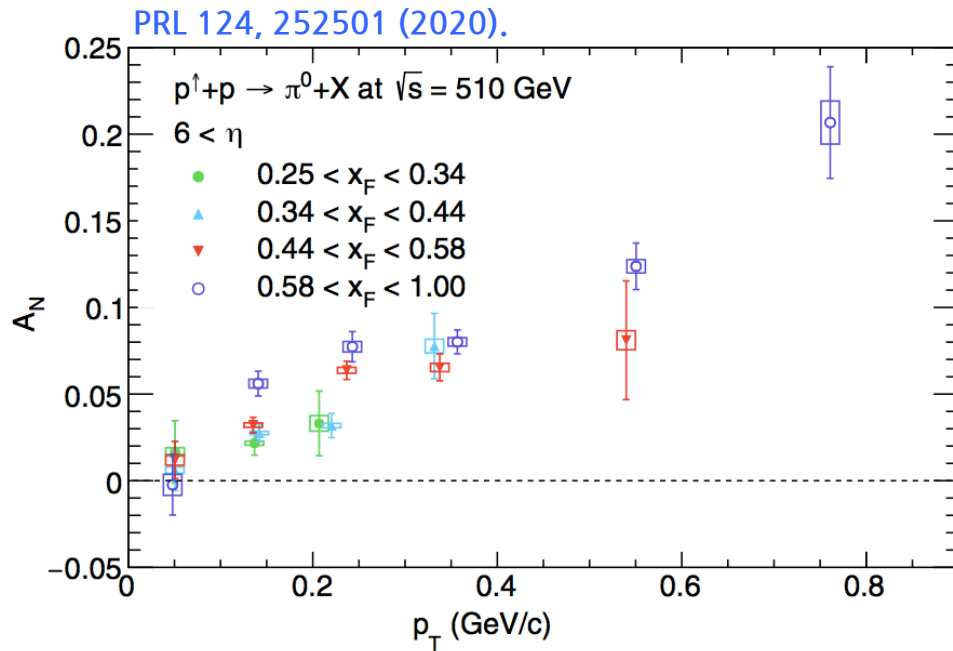
arXiv:2310.09807 (2023).



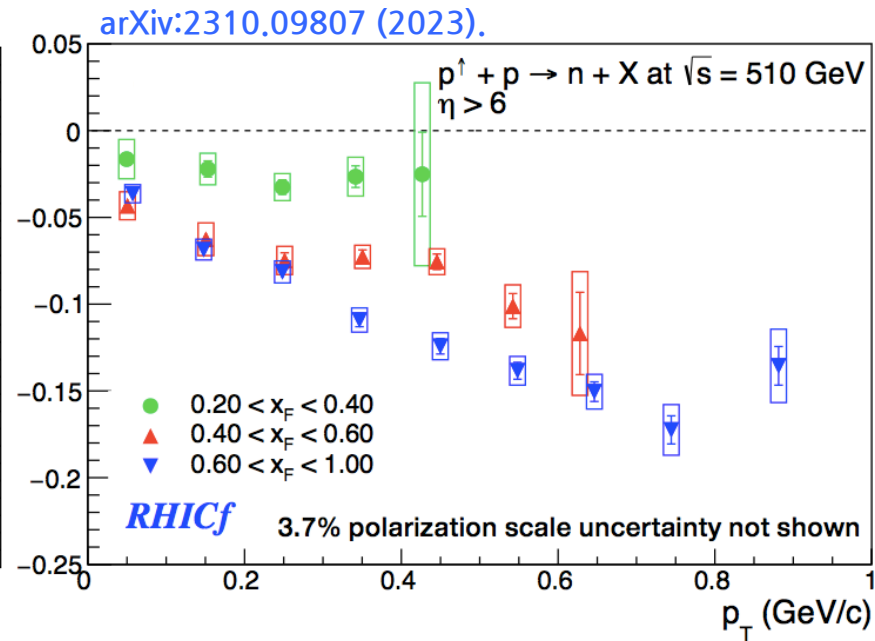
- The RHICf results are consistent with those of PHENIX.
- In the range of  $x_F > 0.4$  and  $p_T < 0.2$  GeV/c, this consistency suggests that there is **no  $\sqrt{s}$  dependence** in the neutron  $A_N$ .

# $A_N$ s of very forward $\pi^0$ and neutron

$\pi^0$

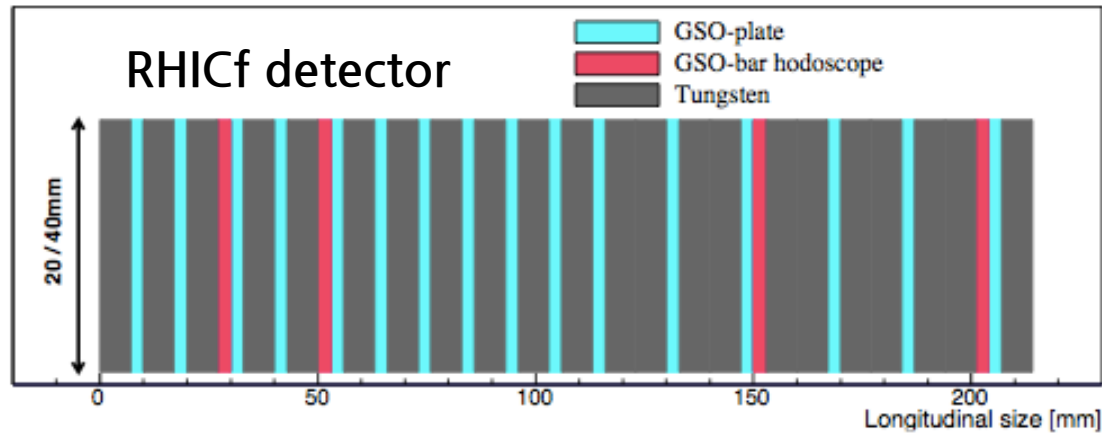


Neutron



- In the very forward region, the **neutron  $A_N$**  is expected to come from the  **$\pi a_1$  interference** and the  **$\pi^0 A_N$**  is expected to come from  **$NN^*$  and  $\Delta\Delta^*$  interferences**.
- Although the  $A_N$ s of  $\pi^0$  and neutron are expected to come from different production mechanisms, they show a couple of common behaviors.
- We may be able to study a correlation between the  $\pi^0$  and neutron  $A_N$ s via that of  $\Lambda$ .

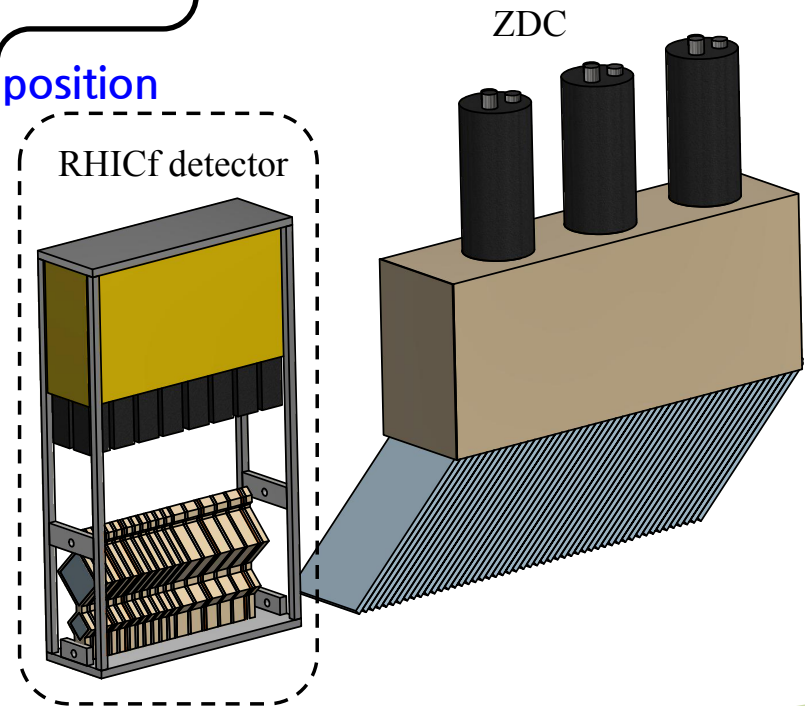
# $\Lambda$ reconstruction



Photon position  
Photon energy

Neutron position

- $A_N$  of  $\Lambda$  may provide us a correlation between the  $\pi^0$  and neutron  $A_N$ s.
- The  $\Lambda$  can be reconstructed using the RHICf detector and ZDC together.

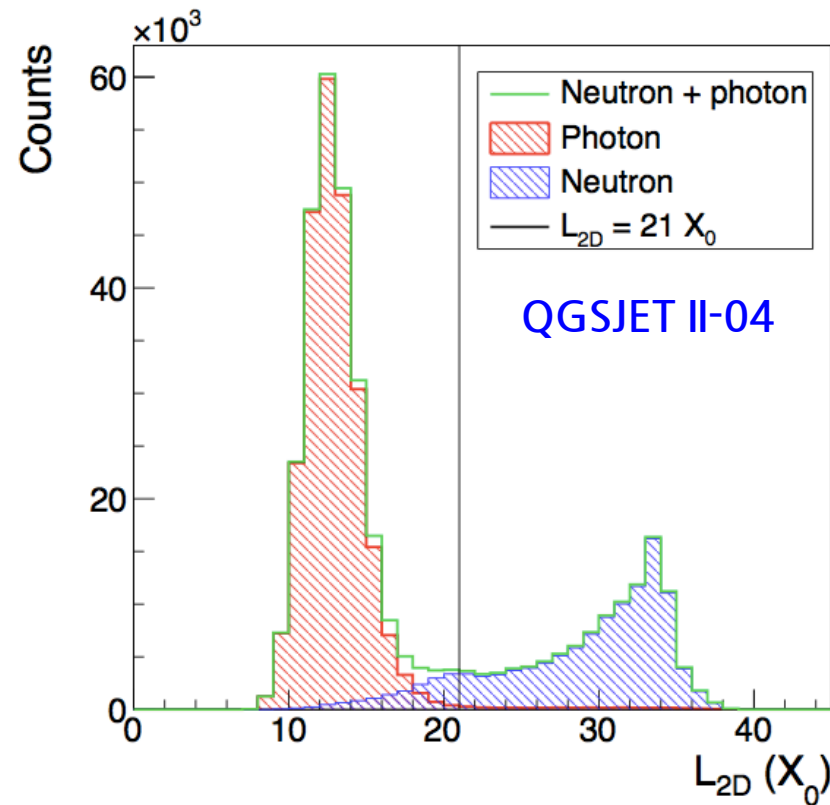


# Summary

- The RHICf experiment has measured  $A_N$ s for very forward  $\pi^0$  and neutron productions in polarized p+p collisions at  $\sqrt{s} = 510$  GeV.
- Non-zero  $A_N$ s were observed even in the  $p_T$  range lower than 1 GeV/c for very forward  $\pi^0$  production.
- To understand the origin of the  $A_N$  for very forward  $\pi^0$  production, we've started the RHICf-STAR combined analysis.
- In the  $x_F$  range higher than 0.6, the neutron  $A_N$  increased in magnitude with  $p_T$  as the model predicted.
- A clear  $x_F$  dependence was observed in the neutron  $A_N$  that has not been predicted by the  $\pi$  and  $a_1$  exchange model.
- There was no  $\sqrt{s}$  dependence in the neutron  $A_N$ .
- We're also reconstructing the  $\Lambda$  to study a possible correlation between the  $\pi^0$  and neutron  $A_N$ s.

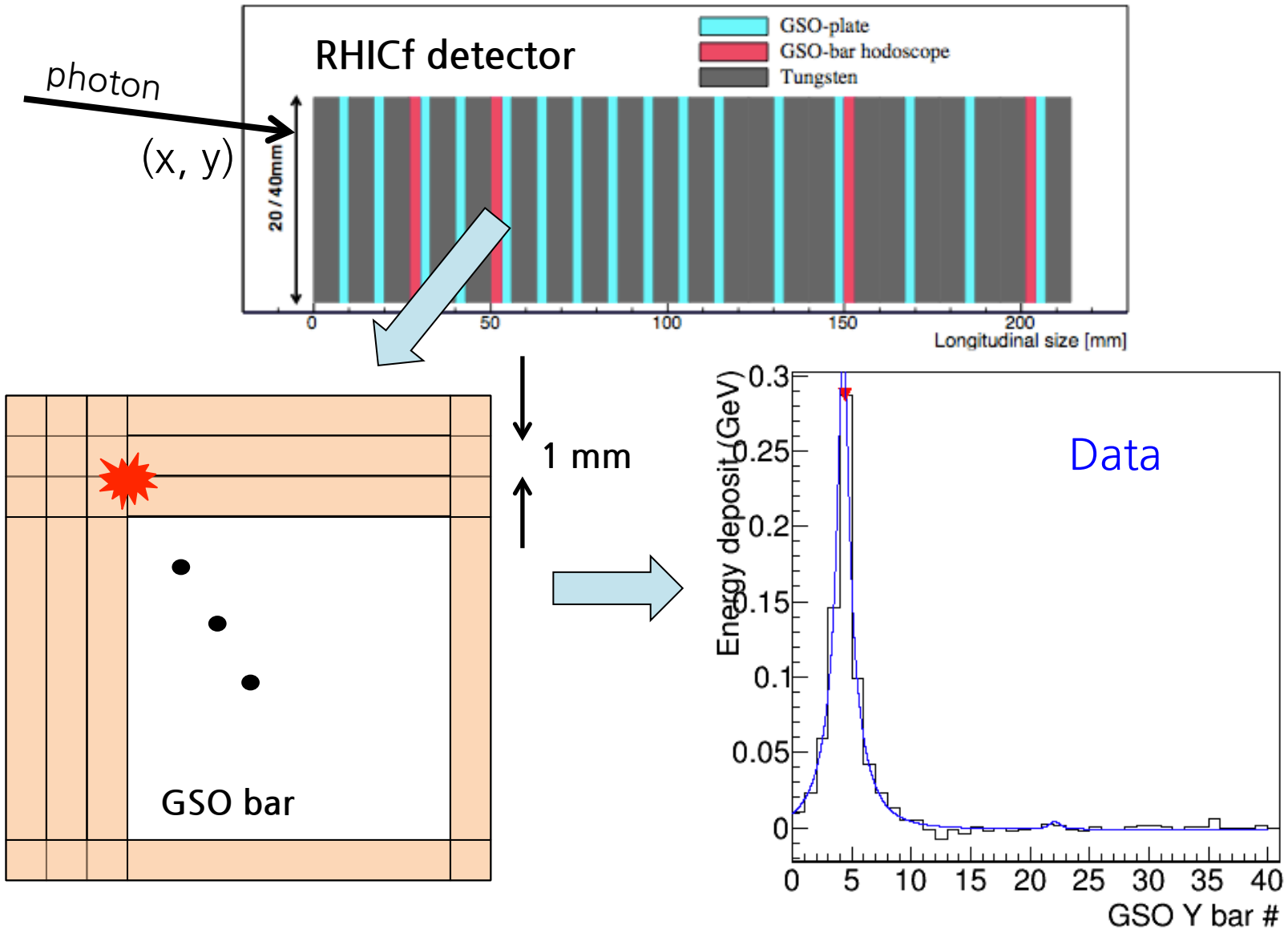
# Backup

# Neutron photon separation



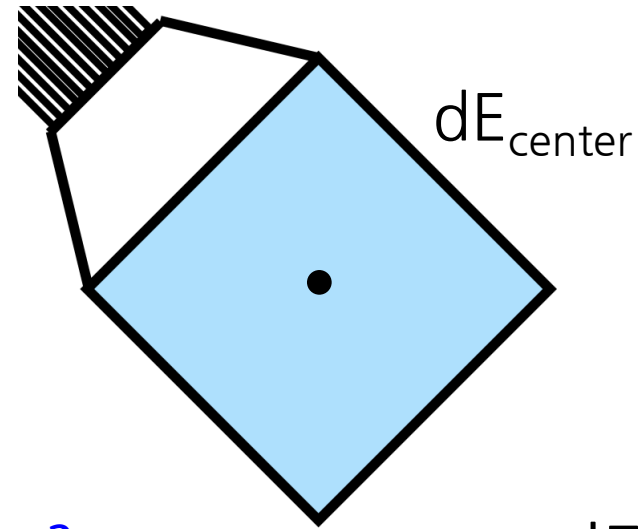
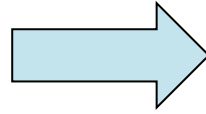
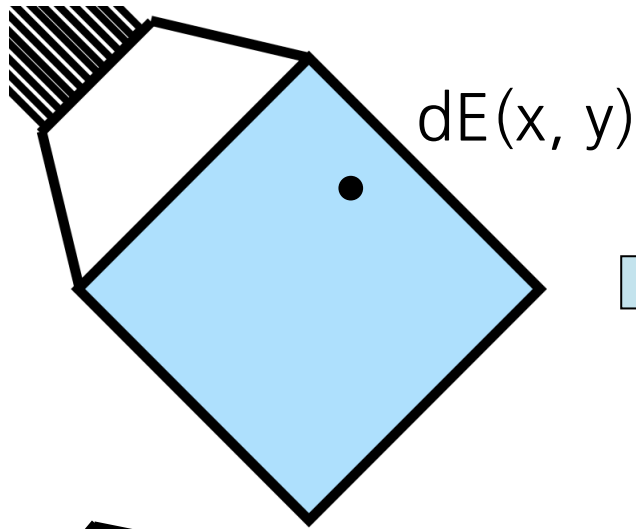
- We used a variable called  $L_{2D}$  for neutron photon separation, which described how early a particle shower was developed in the detector.
- We optimized a  $L_{2D}$  threshold taking into account the particle purity and efficiency.

# Position reconstruction

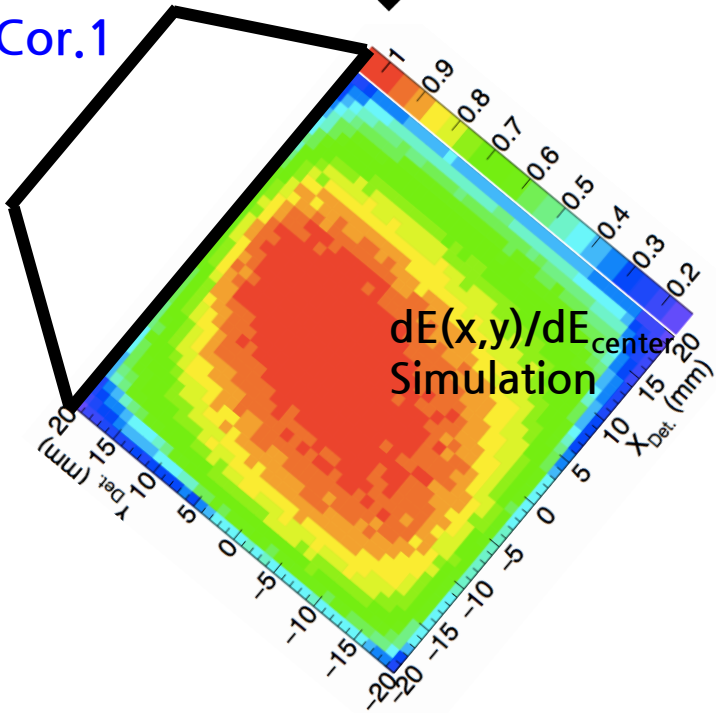




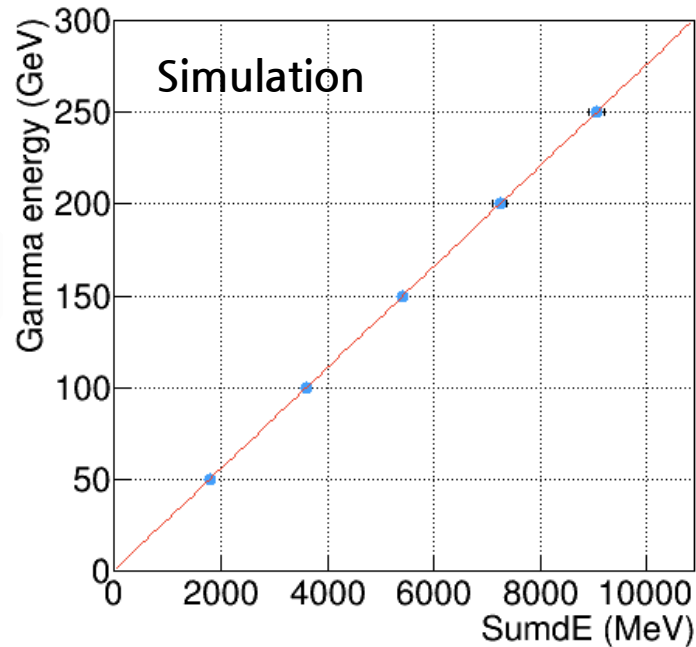
# Energy reconstruction



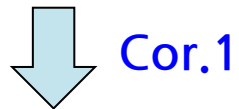
Cor.1



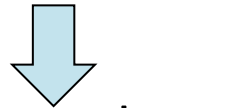
Cor.2



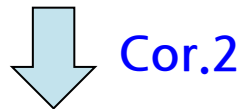
$dE(x, y)$



$dE_{\text{center}}$

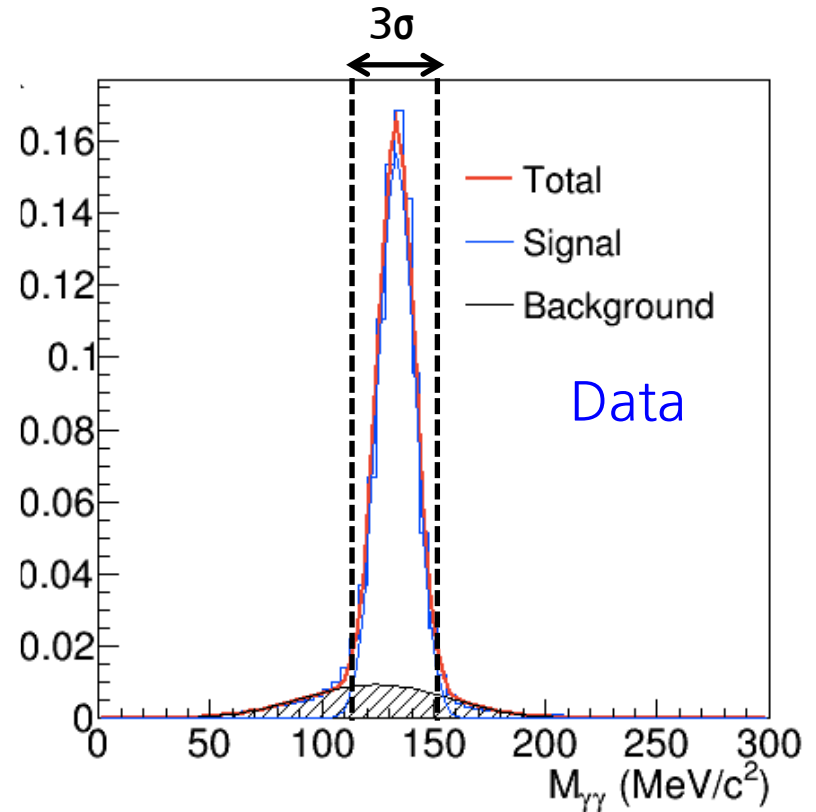
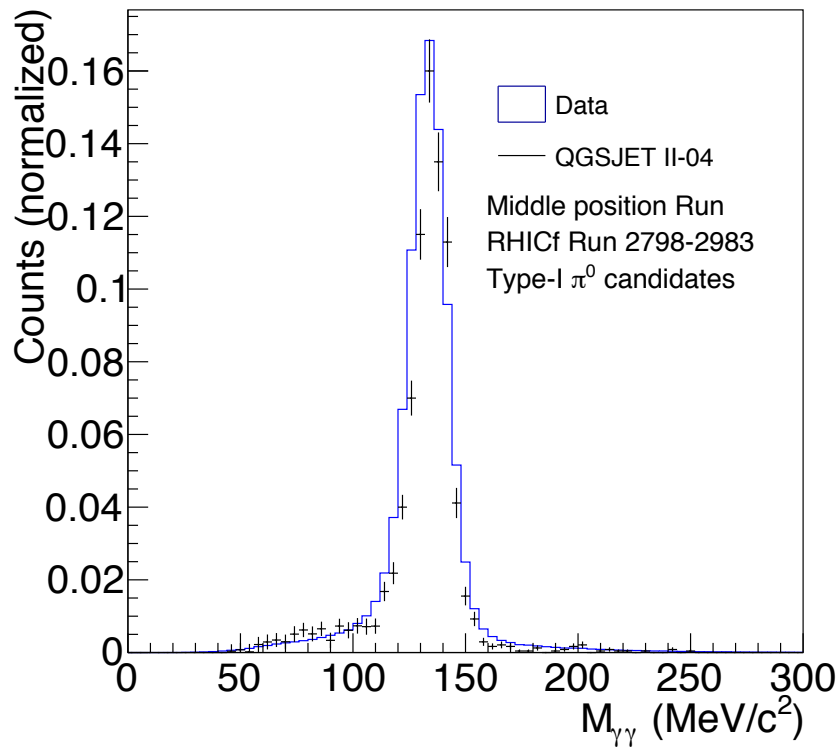


SumdE



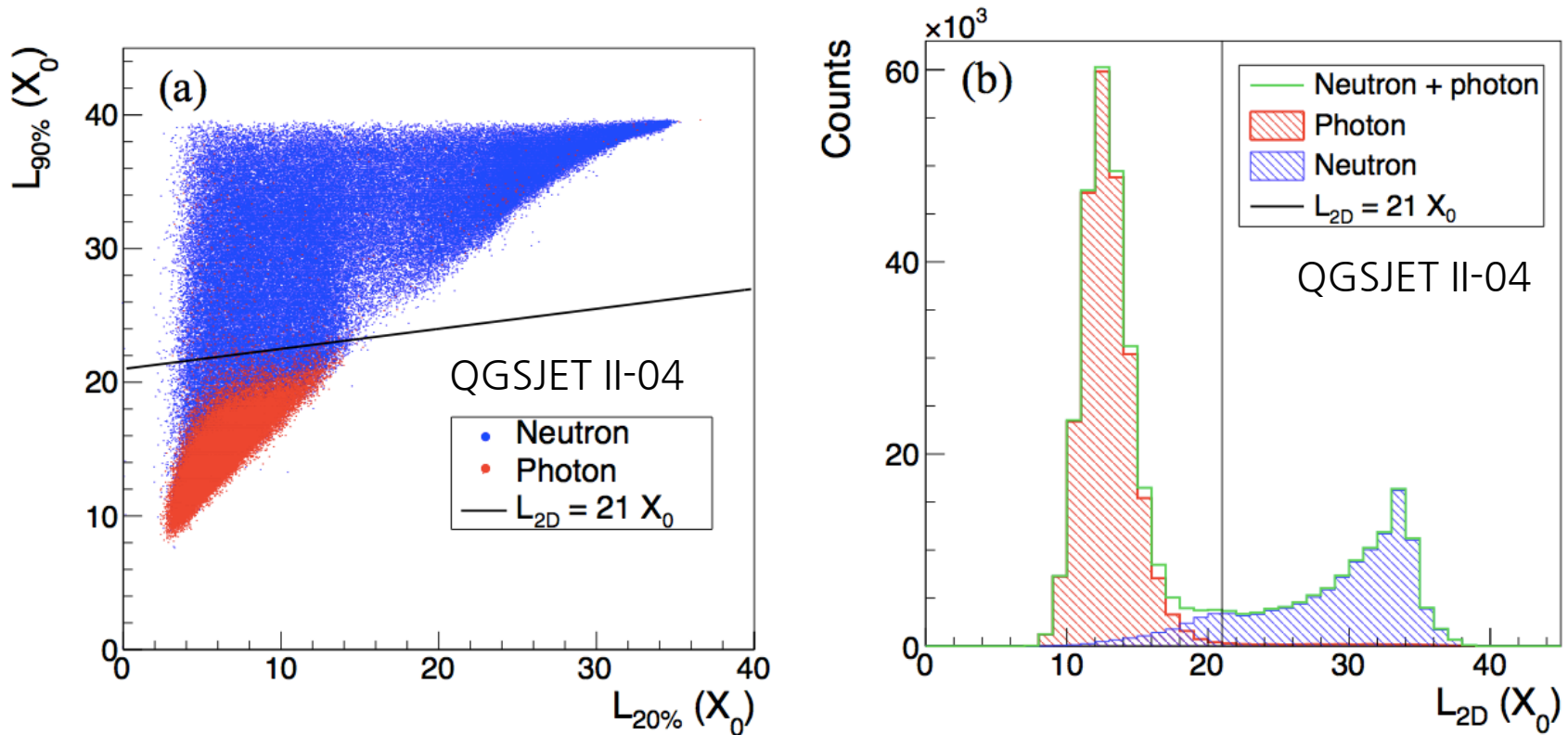
$E_{\text{particle}}$

# $\pi^0$ reconstruction



- Data is well matched with simulation showing a clear  $\pi^0$  peak at 135 MeV/c<sup>2</sup>.
- Invariant mass was fitted by a superposition of polynomial and Gaussian functions, and  $\pi^0$  candidates were selected within  $3\sigma$  tolerance.

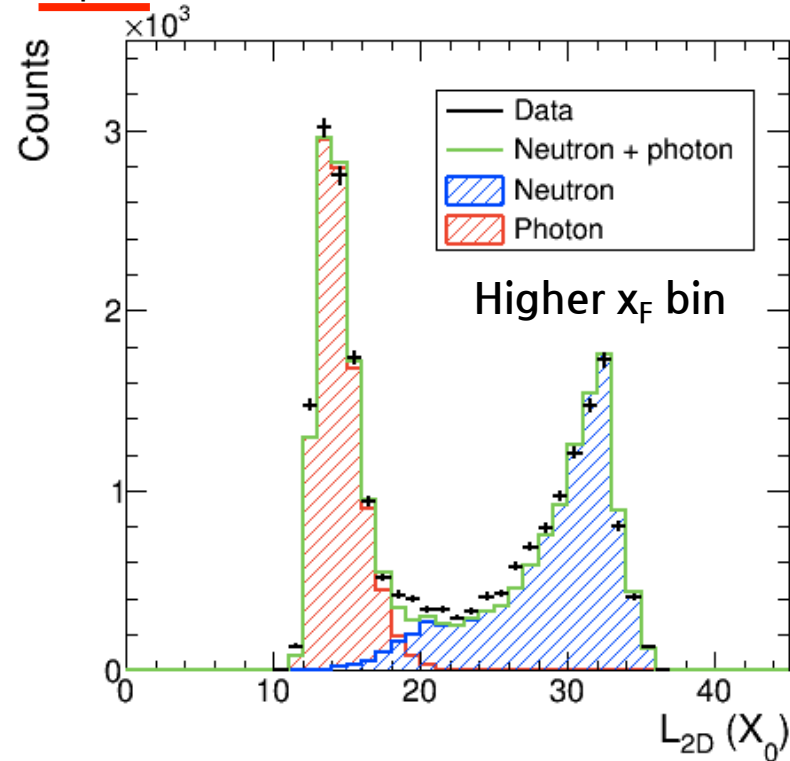
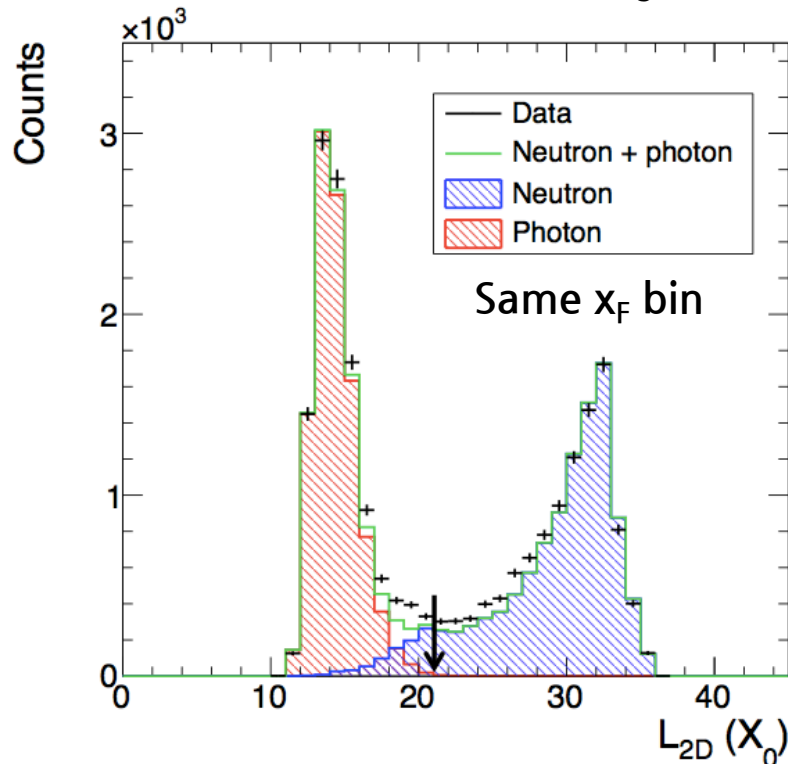
# Neutron photon separation



- An event was considered as a neutron if  $L_{90\%} > aL_{20\%} + b X_0$ .
- Among “a” and “b” values that made the neutron purity higher than 99%, they were optimized so that (purity) x (efficiency) had a maximum value.
- The optimized “a” and “b” are 0.15 and 21, respectively, thereby the  $L_{2D}$  was defined as  $L_{90\%} - 0.15L_{20\%}$ .

# Photon background subtraction

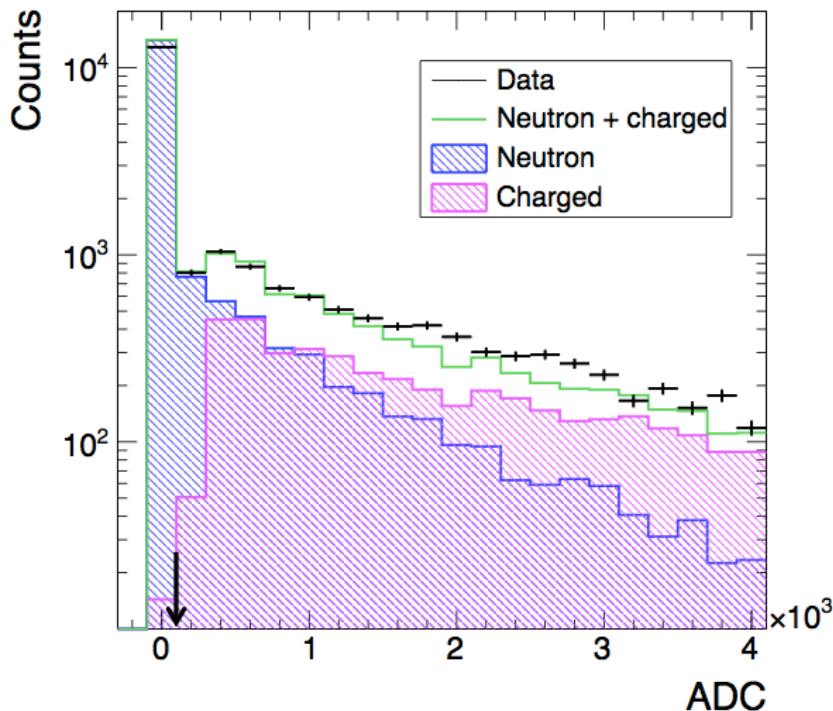
$$N_{\text{trig}} = N_{\text{neu}} + \underline{N_{\text{pho}}} + N_{\text{cha}}$$



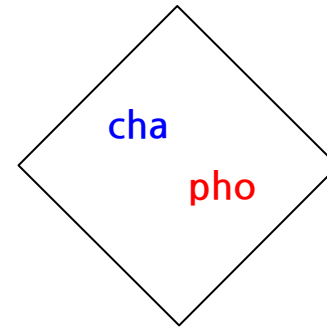
- To estimate and subtract the photon contamination, a template fit was performed to the  $L_{2D}$  distribution.
- To study effect of the discrepancy between the MC and data, the template fit was performed again using the template of the higher  $x_F$  bin.
- $A_N$  difference after unfolding between the two methods was negligible, which was less than 0.0007.  $\rightarrow$  No systematic uncertainty was assigned.

# Charged background subtraction

$$N_{\text{trig}} = N_{\text{neu}} + N_{\text{pho}} + N_{\text{cha}}$$



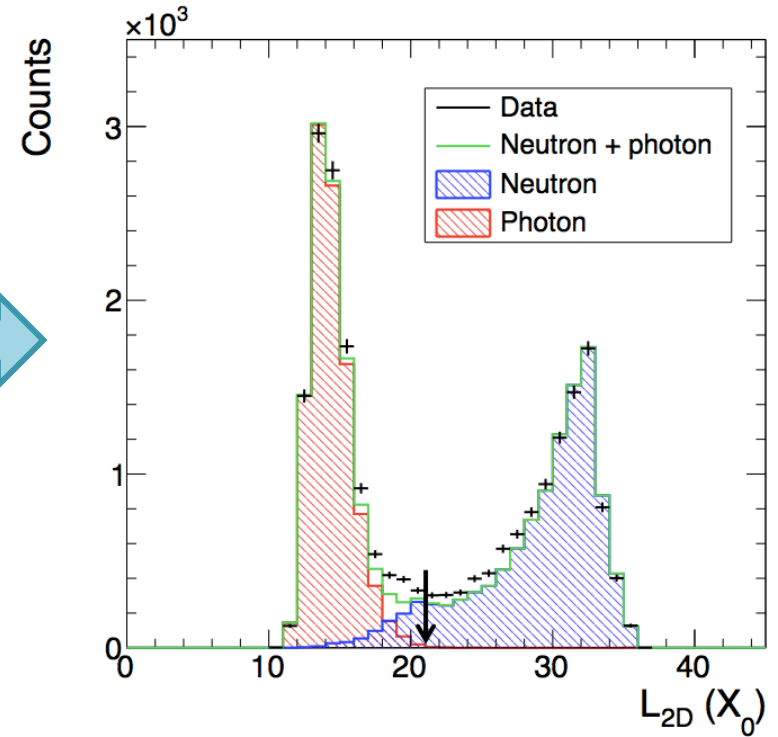
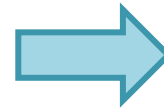
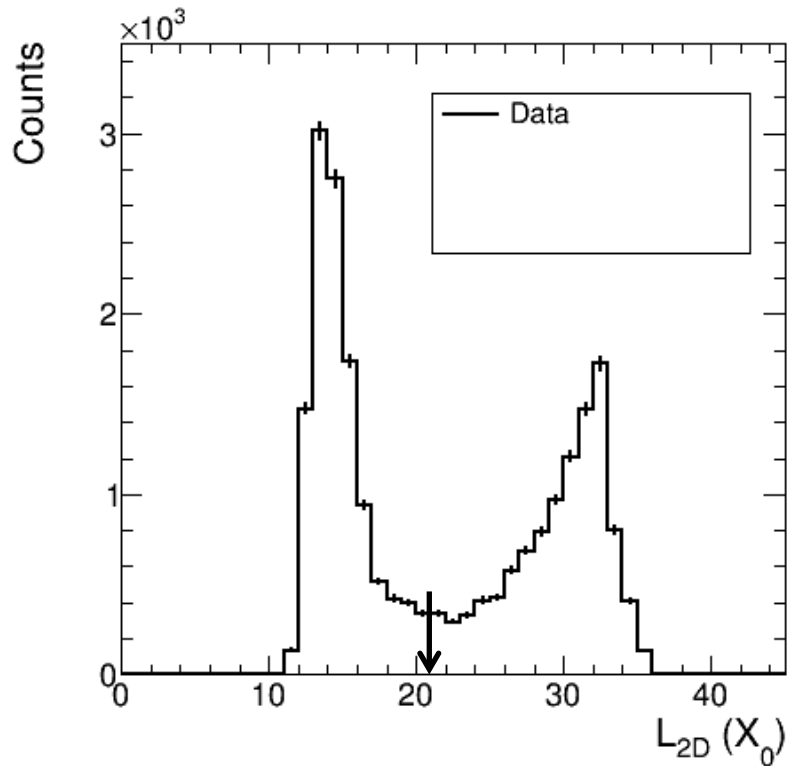
If every charged hadron event includes at least one photon on the detector,



only the photon contamination needs to be subtracted because  $N_{\text{cha}} < N_{\text{pho}}$ .

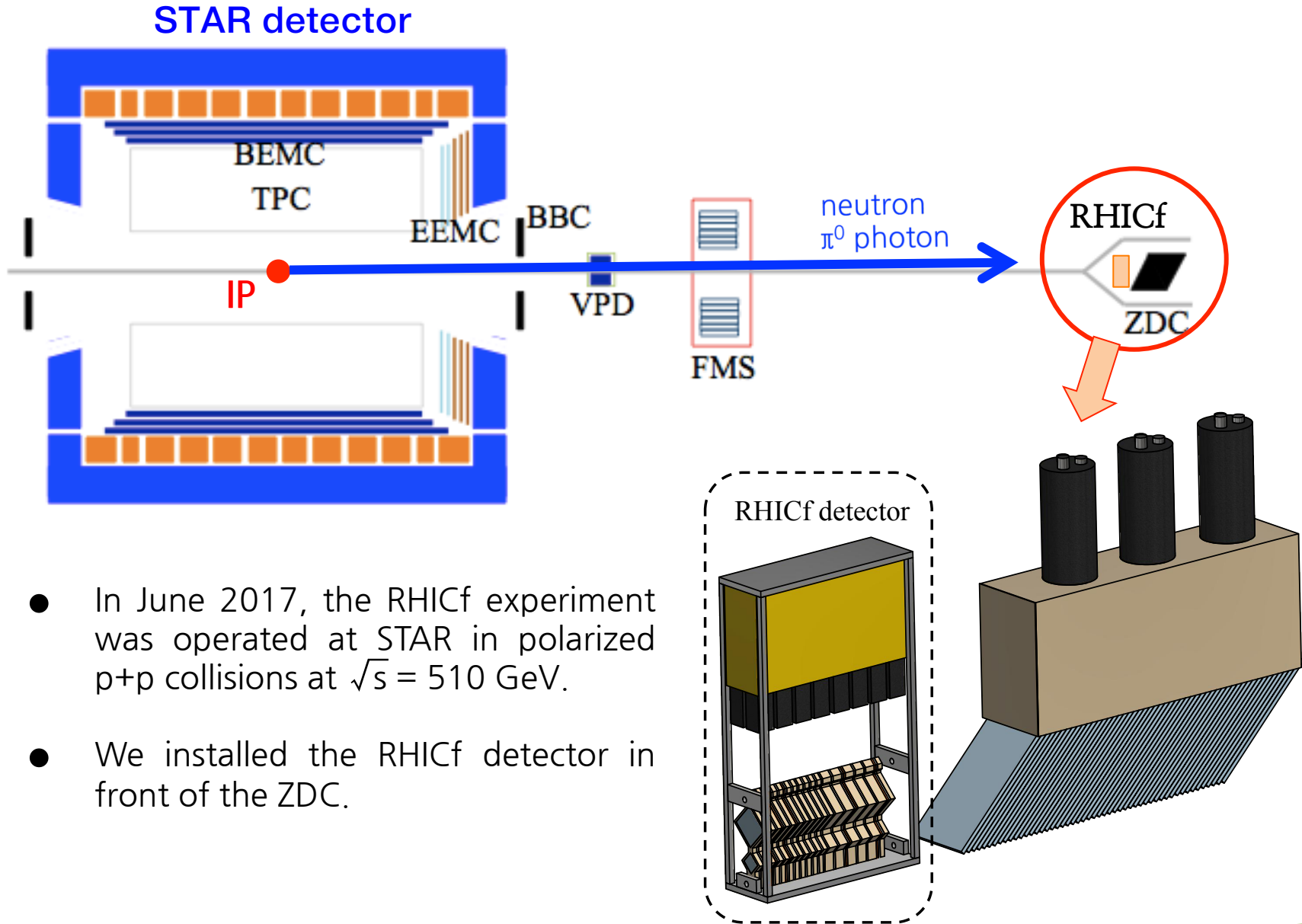
- To estimate and subtract the charged contamination, another template fit was performed to the front counter ADC distribution.
- According to QGSJET II-04, less than 5% of the charged hadron event has photon. → Photon and charged contaminations were subtracted separately.
- There is almost no difference in the resulting  $A_N$  ( $< 0.0004$ ) even if only one contamination was subtracted. → No systematic uncertainty was assigned.

# Photon background

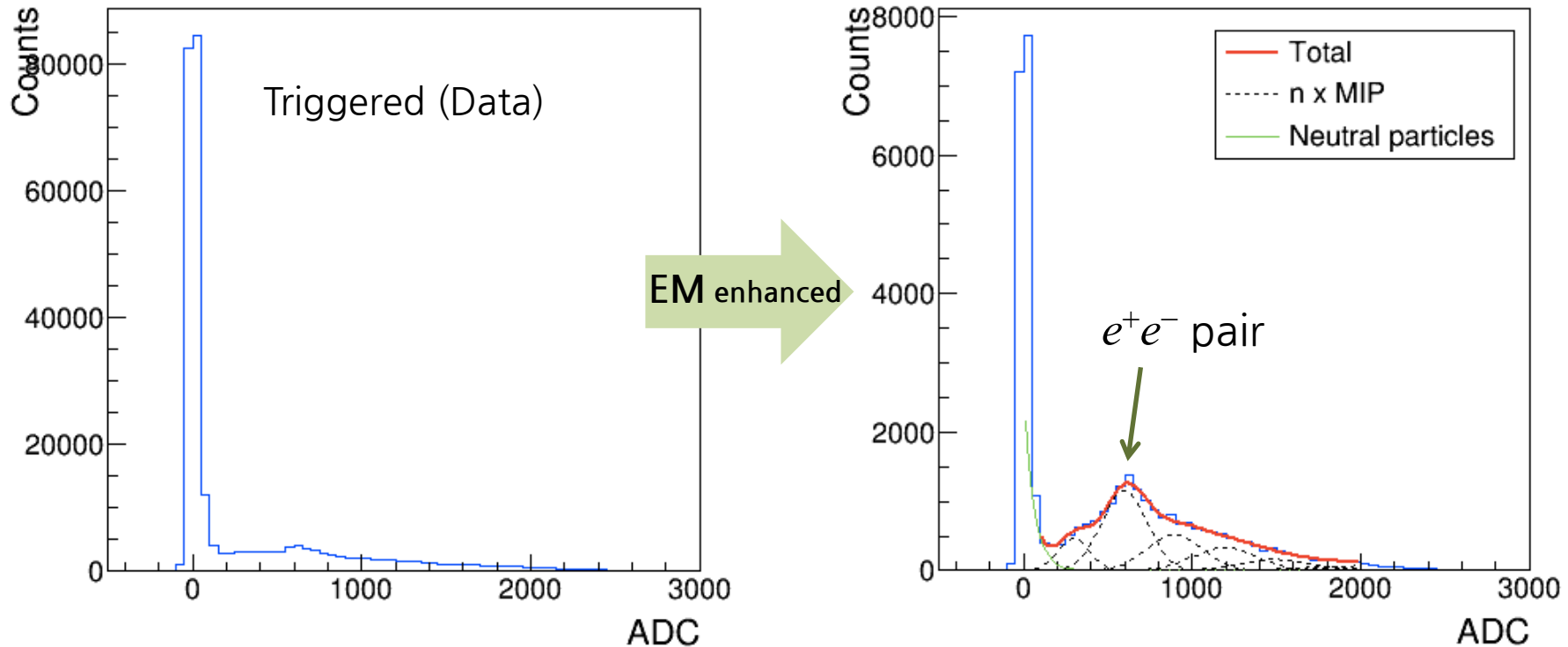


- We performed a template of the  $L_{2D}$  distribution using those of the neutron and photon events obtained from QGSJET II-04 MC sample.
- Photon contamination above the threshold was estimated and subtracted.

# RHIC forward (RHICf) experiment



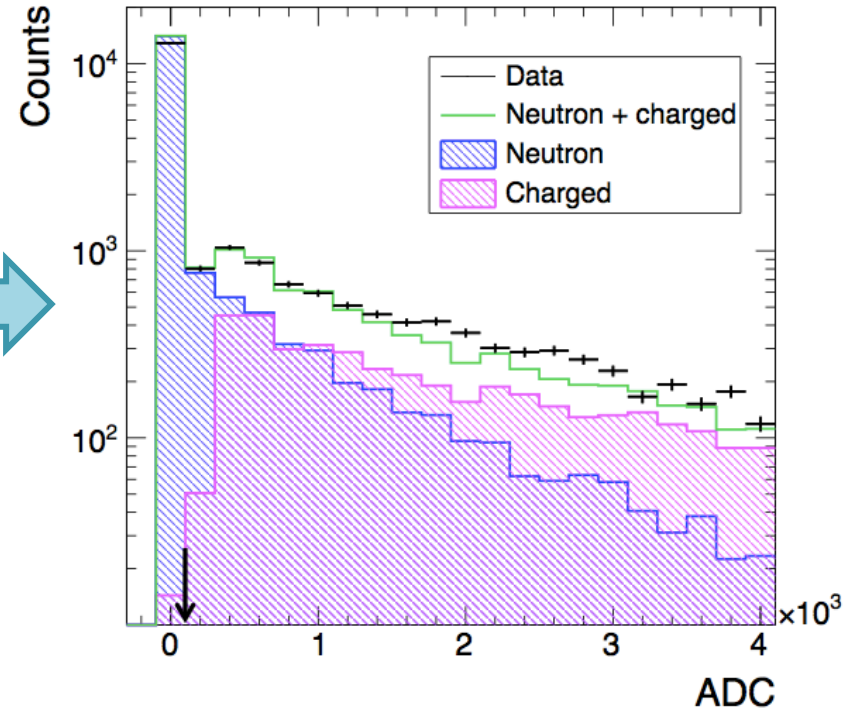
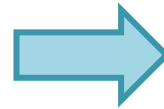
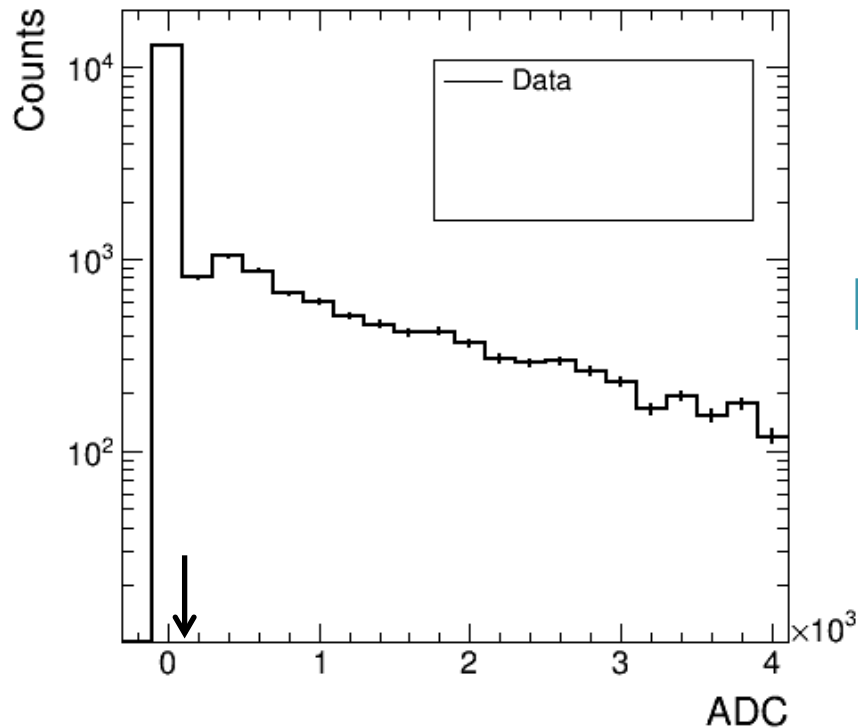
# Reproduction of the front counter response



- To fit the front counter ADC distribution, EM events were enhanced.
- The ADC distribution was fitted by assigning free parameters to mean and sigma of MIP distribution, and number of events of each  $n \times$  MIP distributions.

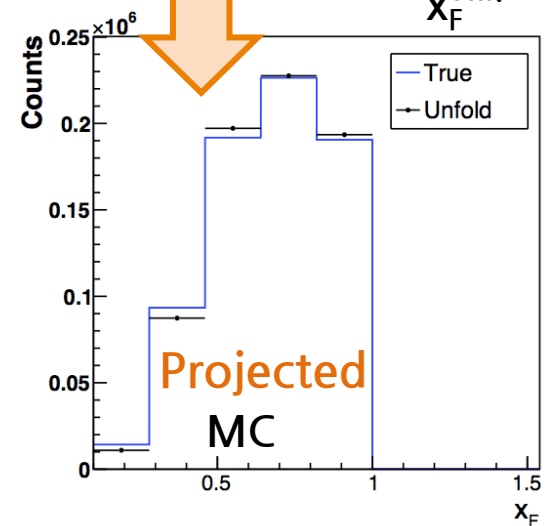
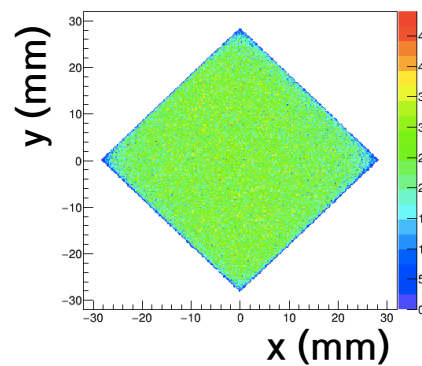
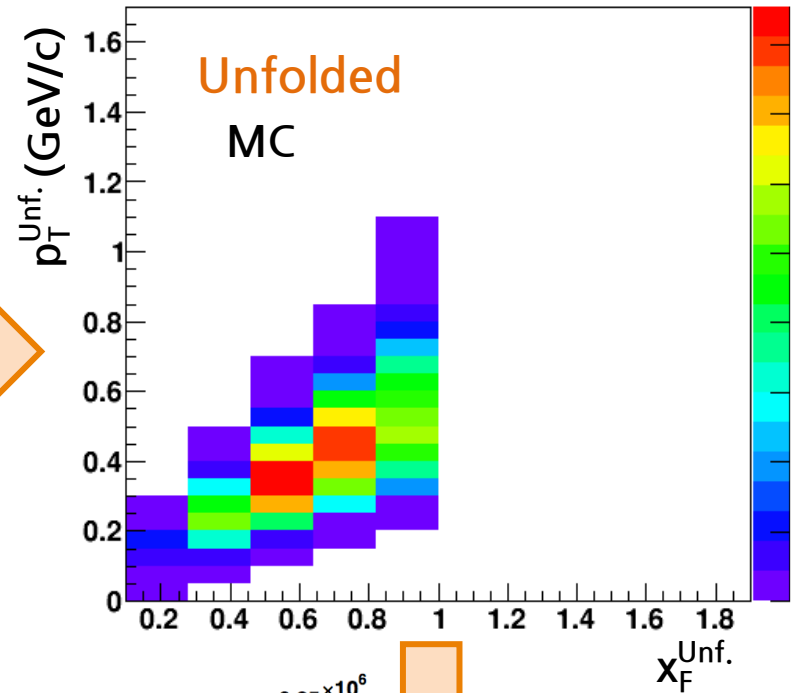
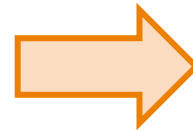
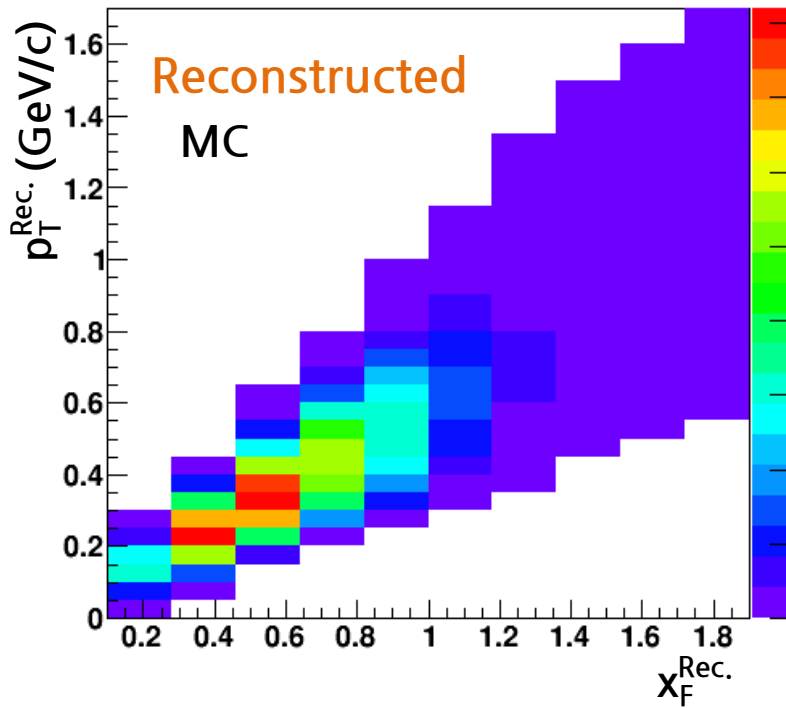


# Charged background



- We applied a threshold to the front counter ADC distribution to suppress the charged hadron events.
- We performed a template of the front counter ADC distribution using those of the neutron and charged hadron events obtained from QGSJET II-04 MC sample.
- Charged hadron contamination below the threshold was estimated and subtracted.

# Unfolding



- Kinematic variables of neutrons,  $x_F$ ,  $p_T$ , and  $\phi$ , were unfolded using Bayesian unfolding.
- For prior distribution, neutrons were uniformly generated to the detector.