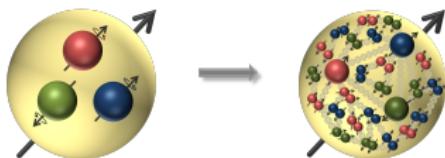


# NUCLEON WITH ONE DYNAMICAL GLUON ON THE LIGHT FRONT



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in collaboration with

**Siqi Xu, Zhi Hu, Sreeraj Nair,  
Xingbo Zhao, Yang Li and James P. Vary**

Light Cone 2022, September 20, 2022



Introduction  
○○○

BLFQ  
○○○○

$|qqq\rangle$   
○○○○○○

$|qqq\rangle + |qqqg\rangle$   
○○○○○○○○○○○○○○○○○○

Conclusions  
○

## Overview



### Introduction

Basis Light-Front Quantization (BLFQ) Approach to

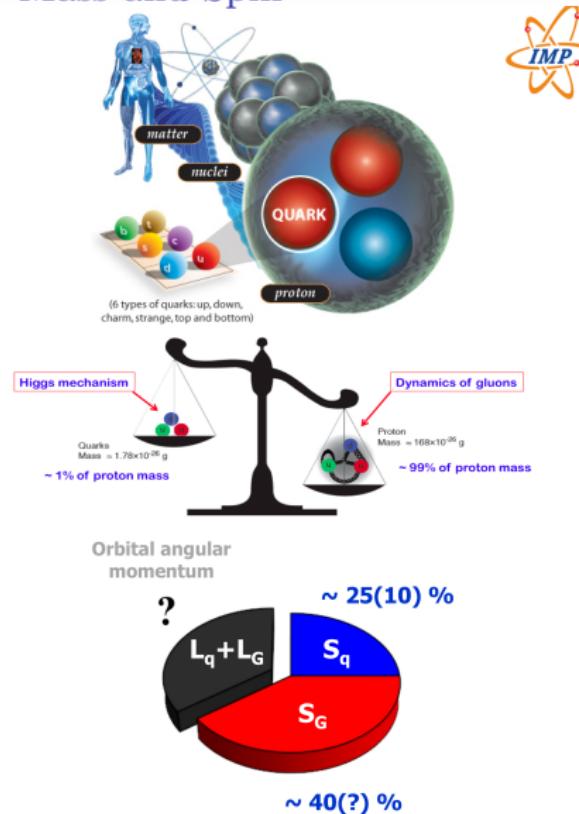
Nucleon:  $|qqq\rangle$  (Based on PRD 104,094036 (2021))

Nucleon:  $|qqq\rangle + |qqqg\rangle$  (Based on arXiv:2209.08584)

### Conclusions

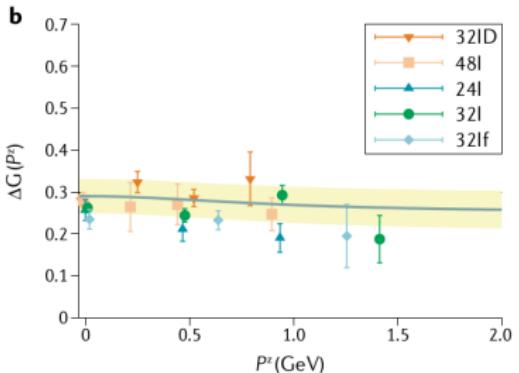
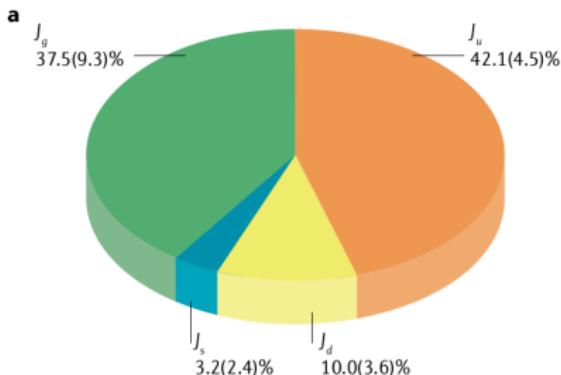
## Fundamental Properties: Mass and Spin

- About 99% of the visible mass is contained within nuclei
- Nucleon: composite particles, built from nearly massless quarks ( $\sim 1\%$  of the nucleon mass) and gluons
- *How does 99% of the nucleon mass emerge?*
- Quantitative decomposition of nucleon spin in terms of quark and gluon degrees of freedom is not yet fully understood.
- *To address these fundamental issues → nature of the subatomic force between quarks and gluons, and the internal landscape of nucleons.*



<sup>1</sup> Pictures (top to bottom) taken from A. Signori's talk, J. Qui talk, C. Lorce's talk

| Spin sum rule   | Formula  | Terms   | Characteristics  |
|---|--|---|--|
| Frame independent ( $J_i$ ) <sup>30</sup>             | $\frac{1}{2}\Delta\Sigma + L_q^z + J_g = \frac{\hbar}{2}$                | $\Delta\Sigma/2$ is the quark helicity<br>$L_q^z$ is the quark OAM<br>$J_g$ is the gluon contribution         | The quark and gluon contributions, $J_q$ and $J_g$ , can be obtained from the GPD moments.<br>A similar sum rule also works for the transverse angular momentum and has a simple parton interpretation                                   |
| Infinite-momentum frame (Jaffe–Manohar) <sup>31</sup> | $\frac{1}{2}\Delta\Sigma + \Delta G + \ell_q + \ell_g = \frac{\hbar}{2}$ | $\Delta G$ is the gluon helicity<br>$\ell_q$ and $\ell_g$ are the quark and gluon canonical OAM, respectively | All terms have partonic interpretations; $\ell_q$ and $\ell_g$ are twist-three quantities.<br>$\Delta G$ is measurable in experiments, including the RHIC spin and the EIC; $\ell_q$ and $\ell_g$ can be extracted from twist-three GPDs |



<sup>1</sup> X. Ji, F. Yuan and Y. Zhao, Nature Reviews Physics 3, 65 (2021)

<sup>2</sup> Y.-B. Yang, R.S. Sufian, A. Alexandru et al., Phys. Rev. Lett. 118, 102001 (2017)

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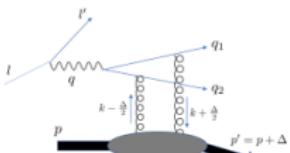
PHYSICAL REVIEW LETTERS 128, 182002 (2022)

## Signature of the Gluon Orbital Angular Momentum

Shohini Bhattacharya<sup>1,\*</sup>, Renaud Boussarie,<sup>2,†</sup> and Yoshitaka Hatta<sup>1,3,‡</sup><sup>1</sup>Physics Department, Brookhaven National Laboratory, Upton, New York 11973, USA<sup>2</sup>CPHT, CNRS, Ecole Polytechnique, Institut Polytechnique de Paris, 91128 Palaiseau, France<sup>3</sup>RIKEN BNL Research Center, Brookhaven National Laboratory, Upton, New York 11973, USA

(Received 30 January 2022; revised 15 March 2022; accepted 4 April 2022; published 2 May 2022)

We propose a novel observable for the experimental detection of the gluon orbital angular momentum (OAM) that constitutes the proton spin sum rule. We consider longitudinal double spin asymmetry in exclusive dijet production in electron-proton scattering and demonstrate that the  $\cos\phi$  azimuthal angle correlation between the scattered electron and proton is a sensitive probe of the gluon OAM at small  $x$  and its interplay with the gluon helicity. We also present a numerical estimate of the cross section for the kinematics of the Electron-Ion Collider.

<sup>1</sup> X. Ji, F. Yuan and Y. Zhao, Nature Reviews Physics 3, 65 (2021)<sup>2</sup> Shohini Bhattacharya's talk: 9/21/22, 1:00 PM

## Basis Light-Front Quantization (BLFQ)

A computational framework for solving relativistic many-body bound state problems in quantum field theories <sup>1</sup>

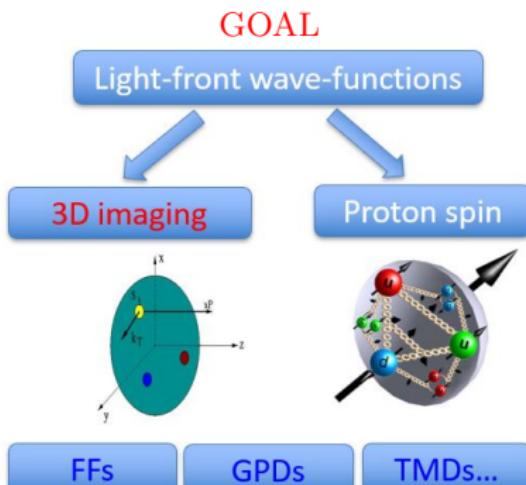


$$P^- P^+ |\Psi\rangle = M^2 |\Psi\rangle$$

- $P^- \equiv P^0 - P^3$  : light-front Hamiltonian
- $P^+ \equiv P^0 + P^3$  : longitudinal momentum
- $|\Psi\rangle$  mass eigenstate
- $M^2$  : mass squared eigenvalue for eigenstate  $|\Psi\rangle$
- First-principle / effective Hamiltonian as input
- Evaluate observables

$$O \sim \langle \Psi | \hat{O} | \Psi \rangle$$

- direct access to light-front wavefunction of bound states



<sup>1</sup>Vary, Honkanen, Li, Maris, Brodsky, Harindranath, *et. al.*, Phys. Rev. C 81, 035205 (2010).

## General Procedure of BLFQ



- ✓ Construct the basis state:  $|\alpha\rangle \equiv |qqq\rangle, |qqqg\rangle, \dots$ 
  - Single-particle coordinates
  - Longitudinal direction: plane-wave basis  $\Psi_k(x^-) = \frac{1}{2L} e^{i\frac{\pi}{L} kx^-}$ .
  - Transverse direction: 2D harmonic oscillator basis,  $\phi_{nm}(\vec{p}_\perp; b)$ .
- ✓ Derive/modelling the Light-Front Hamiltonian:  $H_{\text{eff}}$ 
  - Kinetic Energy + Confinement + Spin dependent interactions
  - Incorporates transverse CM motion; mixed up with intrinsic motion
  - Introduce a constraint term  $H' = \lambda_L(H_{\text{c.m.}} - 2b^2 I)$  into  $H_{\text{eff}}$  to factorize transverse CM motion from intrinsic motion
  - The effective Hamiltonian we diagonalize:  
$$H'_{\text{eff}} = H_{\text{eff}} - \left(\sum_i \vec{p}_{i\perp}\right)^2 + \lambda_L(H_{\text{c.m.}} - 2b^2 I)$$
  - $\lambda_L$ : sufficiently large and positive
- ✓ Evaluate observables using resulting LFWFs
  - Structural information of bound state

# Solution proposed by BLFQ



## Discrete basis and their direct product

2D HO  $\phi_{nm}(p^\perp)$  in the transverse plane

Plane-wave in the longitudinal direction

Light-front helicity state for spin d.o.f.

## Truncation

$$\sum_i (2n_i + |m_i| + 1) \leq N_{\max}$$

$$\sum_i k_i = K, \quad x_i = \frac{k_i}{K}$$

$$\sum_i (m_i + \lambda_i) = M_J$$

$$\alpha_i = (k_i, n_i, m_i, \lambda_i)$$

$$|\alpha\rangle = \otimes_i |\alpha_i\rangle$$

Fock sector truncation

- Fock expansion of hadronic bound states:

$$|\text{Meson}\rangle = \psi_{(q\bar{q})} |q\bar{q}\rangle + \psi_{(q\bar{q}+1g)} |q\bar{q}g\rangle + \psi_{(q\bar{q}+q\bar{q})} |q\bar{q}q\bar{q}\rangle + \dots,$$

$$|\text{Baryon}\rangle = \psi_{(3q)} |qqq\rangle + \psi_{(3q+1g)} |qqqg\rangle + \psi_{(3q+q\bar{q})} |qqqq\bar{q}\rangle + \dots,$$

<sup>1</sup>Vary, Honkanen, Li, Maris, Brodsky, Harindranath, *et. al.*, Phys. Rev. C 81, 035205 (2010).

## Applications of BLFQ

### QCD systems



- **Heavy mesons:** spectrum, decay constant, elastic form factor, radii, radiative transitions, distribution amplitude, PDFs, GPDs

—Li, Chen, Zhao, Maris, Vary, Adhikari, M Li, Tang, A El-Hady, Lan, Wu, CM (2016 - 2022)

- **Light mesons:** spectrum, decay constant, elastic form factor, radii, distribution amplitude, PDFs, GPDs, TMDs

—Jia, Vary, Lan, Zhao, Qian, Li, Fu, J. Chen, Wu, CM (2018 - 2022)

- **Baryons:** EMFFs, axial form factor, radii, PDFs, GPDs, TMDs, OAM

—Xu, Hu, Peng, Zhu, Zhao, Li, Chakrabarti, Vary, Lan, Liu, CM (2019-2022)

- **Tetraquarks:** Masses of all-charm tetraquarks

—Kuang, Serafin, Zhao, Vary (2022)

### QED systems

- **Electron:** anomalous magnetic moments, GPDs
- **positronium:** wave function, spectroscopy, FFs, GPDs
- **Photon:** wave function, structure functions, GPDs, TMDs

—Zhao, Wiecki, Li, Honkanen, Maris, Vary, Brodsky, Fu, Hu, Nair, CM (2013 - 2022)

## Nucleon within BLFQ

- The LF eigenvalue equation:  $H_{\text{eff}}|\Psi\rangle = M^2|\Psi\rangle$



$$\begin{aligned} H_{\text{eff}} = & \sum_a \frac{\vec{p}_{\perp a}^2 + m_a^2}{x_a} + \frac{1}{2} \sum_{a \neq b} \kappa^4 \left[ x_a x_b (\vec{r}_{\perp a} - \vec{r}_{\perp b})^2 - \frac{\partial_{x_a} (x_a x_b \partial_{x_b})}{(m_a + m_b)^2} \right] \\ & + \frac{1}{2} \sum_{a \neq b} \frac{C_F 4\pi \alpha_s}{Q_{ab}^2} \bar{u}_{s'_a}(k'_a) \gamma^\mu u_{s_a}(k_a) \bar{u}_{s'_b}(k'_b) \gamma^\nu u_{s_b}(k_b) g_{\mu\nu} \end{aligned}$$

- For the first Fock sector:
$$|qqq\rangle = |n_{q_1}, m_{q_1}, k_{q_1}, \lambda_{q_1}\rangle \otimes |n_{q_2}, m_{q_2}, k_{q_2}, \lambda_{q_2}\rangle \otimes |n_{q_3}, m_{q_3}, k_{q_3}, \lambda_{q_3}\rangle$$
- Transverse : 2D harmonic oscillator basis  $\phi_{nm}(\vec{p}_\perp)$ ;  
Plane wave basis in longitudinal direction.
- The valence wavefunction in momentum space:

$$\Psi_{\{x_i, \vec{p}_\perp i, \lambda_i\}}^{M_J} = \sum_{n_i, m_i} \left[ \psi(\alpha_i) \prod_{i=1}^3 \phi_{n_i m_i}(\vec{p}_\perp i) \right]$$

## Nucleon Form Factors



$$G_E(q^2) = F_1(q^2) - \frac{q^2}{4M^2} F_2(q^2),$$

$$\langle P'; \uparrow | \frac{J^+(0)}{2P^+} | P; \uparrow \rangle = F_1(q^2)$$

$$\langle P'; \uparrow | \frac{J^+(0)}{2P^+} | P; \downarrow \rangle = -(q^1 - iq^2) \frac{F_2(q^2)}{2M}$$

Drell & Yan (PRL, 70); West (PRL, 70)

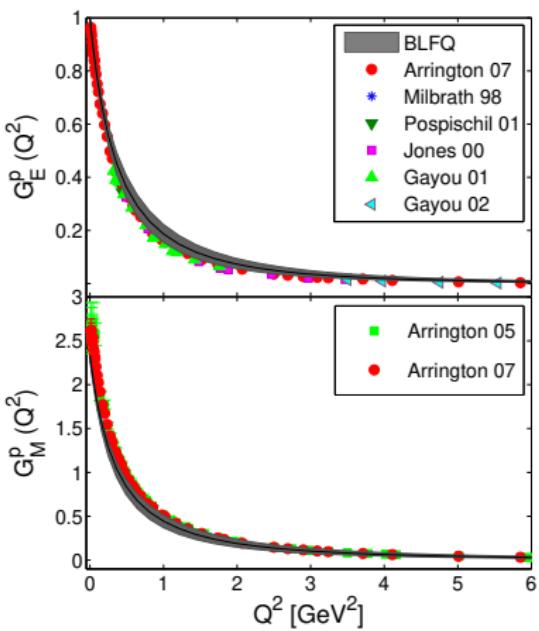
TABLE I. Model parameters for the basis truncations  $N_{\max} = 10$  and  $K = 16.5$ .

| $m_{q/k}$ | $m_{q/g}$ | $\kappa$ | $\alpha_s$    |
|-----------|-----------|----------|---------------|
| 0.3 GeV   | 0.2 GeV   | 0.34 GeV | $1.1 \pm 0.1$ |

$$N_{\max} = 10, K = 16.5$$

Basis truncation :

$$\sum_i (2n_i + |m_i| + 1) \leq N_{\max}; \quad K = \sum_i k_i$$



<sup>1</sup> CM, Siqi Xu, et. al., Phys. Rev. D **102**, 016008 (2020)

<sup>2</sup> Xu, CM, Lan, Zhao, Li, and Vary, Phys. Rev. D **104**, 094036 (2021)

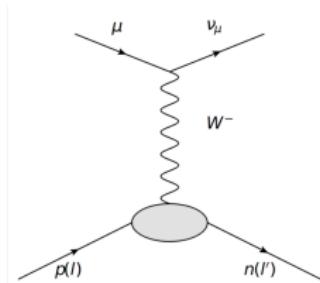
## Axial Form Factor

$$\langle N(p)|A^\mu|N(p')\rangle = \bar{u}(p') \left[ \gamma^\mu G_A(t) + \frac{(p' - p)^\mu}{2m} G_p(t) \right] \gamma_5 u(p)$$

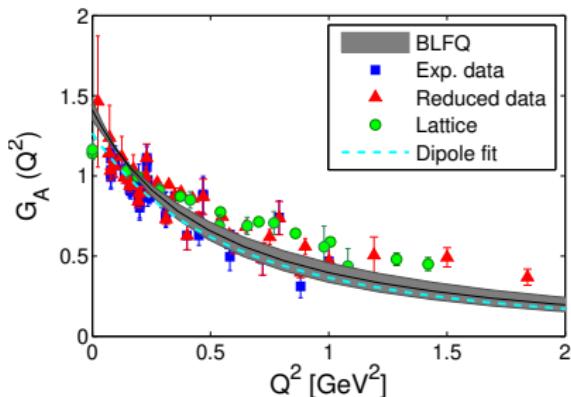


- Axial vector current:  
 $A^\mu = \bar{q}\gamma^\mu\gamma_5 q$
- Measured by ordinary muon capture (OMC)

$$\mu^-(l) + p(r) \rightarrow \nu_\mu(l') + n(r')$$



- Provide information on spin distributions

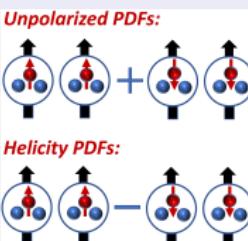
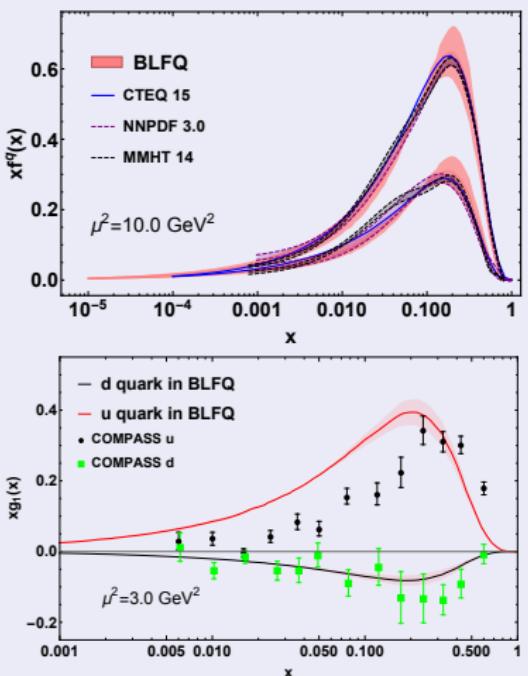


$$G_A(Q^2) = G_u(Q^2) - G_d(Q^2)$$

<sup>1</sup> CM, Siqi Xu *et. al.*, Phys. Rev. D **102**, 016008 (2020)

# Parton Distribution Functions

Xu, CM, Lan, Zhao, Li, and Vary, PRD 104, 094036 (2021)



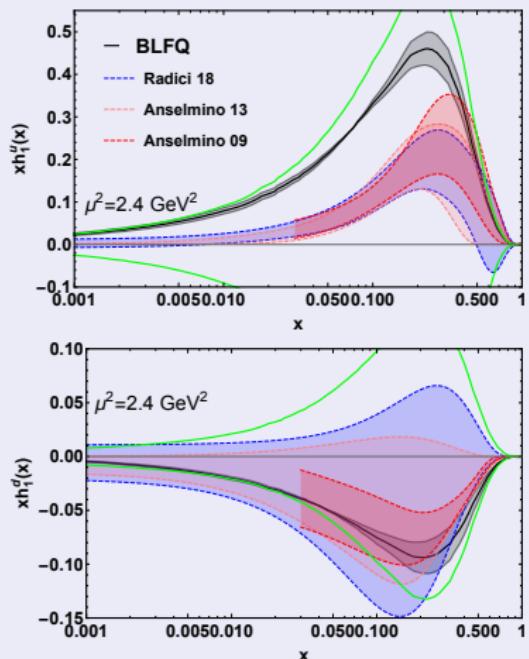
- **Unpolarized PDFs  $f_1(x)$**  : longitudinal momentum distribution of unpol. quark in unpol. proton.
- **Helicity PDFs  $g_1(x)$**  : longitudinal momentum distribution of the polarized quark
- Results correspond to leading Fock sector only.

<sup>1</sup> NNPDF, EPJC 77, 663 (2017); HMMT, EPJC 75, 204 (2015); CTEQ, PRD 93, 033006 (2016).

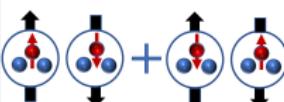
<sup>2</sup> COMPASS Collaboration, Phys. Lett. B 693, 227 (2010).

# Transversity Distribution

Xu, CM, Lan, Zhao, Li, and Vary, PRD 104, 094036 (2021)



## Transversity PDFs :



- **Transversity PDFs** describe correlation between the transverse polarization of the nucleon and the transverse polarization of the parton.
- Satisfy Soffer Bound:
$$|h_1(x)| \leq \frac{1}{2} |f_1(x) + g_1(x)|$$
- Results correspond to leading Fock sector only, **missing higher Fock sectors**.

<sup>1</sup> M. Radici and A. Bacchetta, Phys. Rev. Lett. 120, 192001 (2018).

<sup>2</sup> M. Anselmino, et. al., Phys. Rev. D 87, 094019 (2013).

## x-Dependent Squared Radius



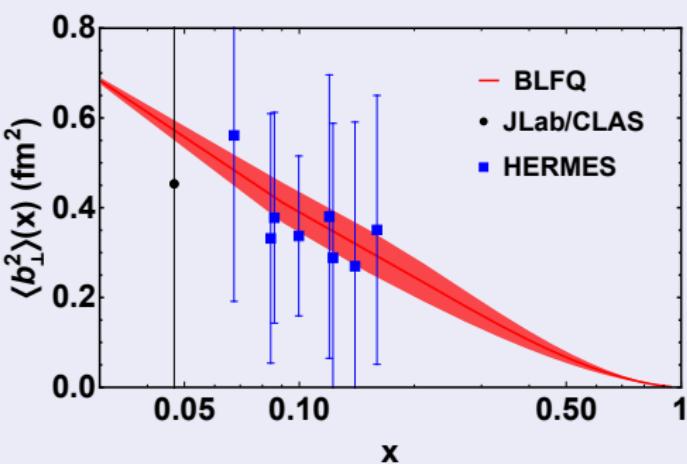
$$\langle b_{\perp}^2 \rangle^q(x) = \frac{\int d^2 \vec{b}_{\perp} b_{\perp}^2 H^q(x, \vec{b}_{\perp})}{\int d^2 \vec{b}_{\perp} H^q(x, \vec{b}_{\perp})},$$

- Transverse squared radius:

$$\langle b_{\perp}^2 \rangle = \sum_q e_q \int_0^1 dx f^q(x) \langle b_{\perp}^2 \rangle^q(x)$$

- BLFQ result:  $\langle b_{\perp}^2 \rangle = 0.40 \pm 0.04$  fm<sup>2</sup>

- Experimental data <sup>2</sup>:  
 $\langle b_{\perp}^2 \rangle_{\text{exp}} = 0.43 \pm 0.01$  fm<sup>2</sup>



<sup>1</sup>Xu, CM, Lan, Zhao, Li, and Vary, PRD 104, 094036 (2021)

<sup>2</sup>R. Dupre, M. Guidal and M. Vanderhaeghen, PRD 95, 011501 (2017).

# Effective Hamiltonian with One Dynamical Gluon



$$| \text{Baryon} \rangle = a | qqq \rangle + b | qqqg \rangle \boxed{|} + c | qqq\bar{q} \rangle + \dots$$

kinetic energy

$$H_{\text{eff}} = \sum_a \frac{\vec{p}_{\perp a}^2 + m_a^2}{x_a} + \frac{1}{2} \sum_{a \neq b} \kappa^4 [x_a x_b (\vec{r}_{\perp a} - \vec{r}_{\perp b})^2]$$

transverse confining potential [2]

$$- \frac{1}{2} \sum_{a \neq b} \kappa^4 \left[ \frac{\partial x_a (x_a x_b \partial x_b)}{(m_a + m_b)^2} \right] +$$

 $H_{\text{vertex}} + H_{\text{inst}}$ 

longitudinal confining potential [3]

QCD interactions [4]

<sup>1</sup> S. Xu, CM, X. Zhao, Y. Li, J. P. Vary, 2209.08584 [hep-ph].

<sup>2</sup> Brodsky, Teramond, Dosch and Erlich, Phys. Rep. 584, 1 (2015).

<sup>3</sup> Li, Maris, Zhao and Vary, Phys. Lett. B (2016).

<sup>4</sup> Brodsky, Pauli, and Pinsky, Phys. Rep. 301, 299 (1998).

## QCD Interactions

$$|P_{baryon}\rangle = \psi_1|qqq\rangle + \psi_2|qqqg\rangle$$

$$H_{\text{Interact}} = H_{\text{Vertex}} + H_{\text{inst}} = g\bar{\psi} \gamma^\mu T^a \psi A_\mu^a + \frac{g^2 C_F}{2} j^+ \frac{1}{(i\partial^+)^2} j^+$$



$$N_{\max} = 9, K = 16.5$$

Higher Fock Sector effect

Remove Soft Gluon effect

| $m_u$    | $m_d$    | $\kappa$ | $m_g$    | $m_{int}$ | $b_{inst}$ | $b$      | $g$  |
|----------|----------|----------|----------|-----------|------------|----------|------|
| 0.30 GeV | 0.25 GeV | 0.54 GeV | 0.50 GeV | 1.80 GeV  | 3.00 GeV   | 0.70 GeV | 2.40 |

$\downarrow$

Different Mass  
Asymmetry of u and d

$\uparrow$

UV Cutoff  
In Instantaneous term



## Nucleon Form Factors



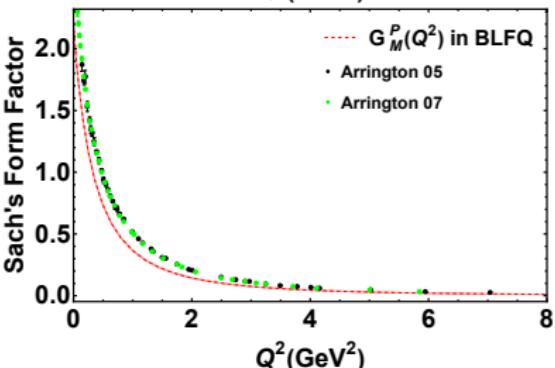
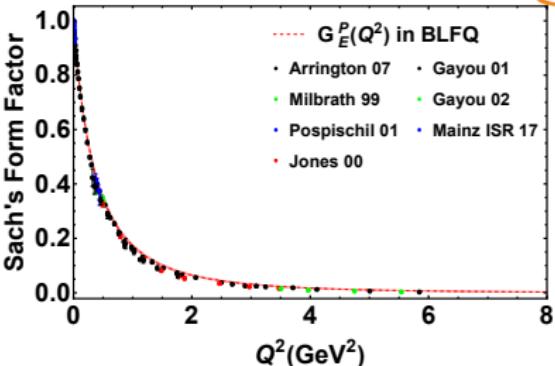
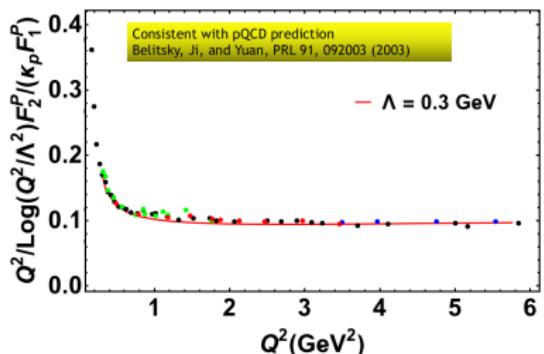
$$\langle P'; \uparrow | \frac{J^+(0)}{2P^+} | P; \uparrow \rangle = F_1(q^2)$$

$$\langle P'; \uparrow | \frac{J^+(0)}{2P^+} | P; \downarrow \rangle = -(q^1 - iq^2) \frac{F_2(q^2)}{2M}$$

Drell & Yan (PRL, 70); West (PRL, 70)

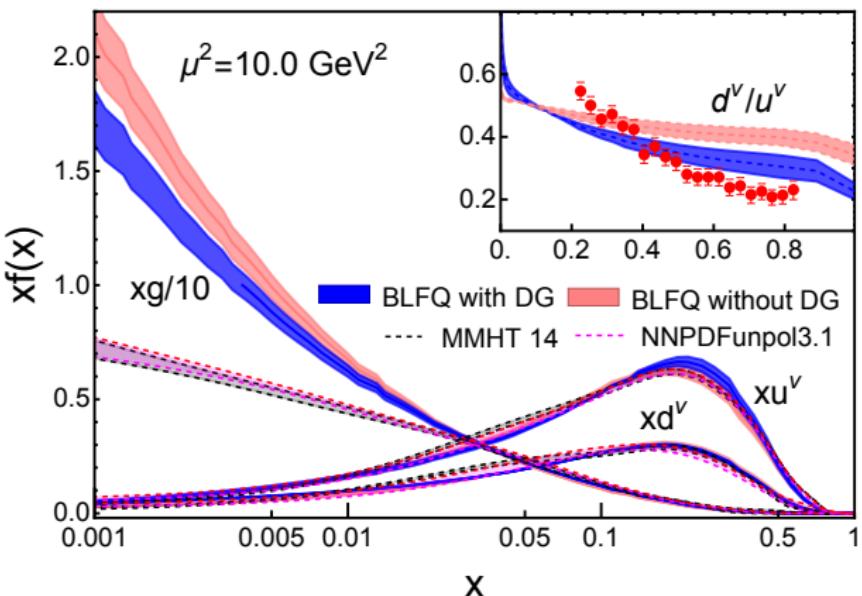
$$\sqrt{\langle r_E^2 \rangle} = 0.85 \pm 0.01 (0.840^{+0.003}_{-0.002}) \text{ fm}$$

$$\sqrt{\langle r_M^2 \rangle} = 0.88 \pm 0.07 (0.849^{+0.003}_{-0.003}) \text{ fm}$$



<sup>2</sup>Xu, CM, Zhao, Li, and Vary, 2209.08584 [hep-ph]

## Unpolarized PDFs



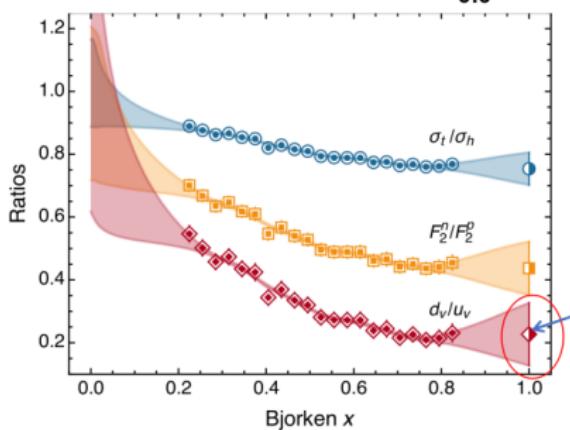
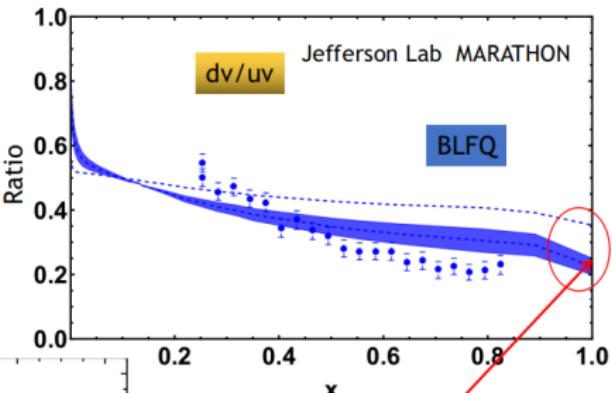
- Within  $|qqq\rangle$ , model scale  $\mu_0^2 = 0.195 \pm 0.020 \text{ GeV}^2$
- Gluon is generated dynamically from the QCD evolution.
- Including dynamical gluon (DG), model scale  $\mu_0^2 = 0.23 - 0.25 \text{ GeV}^2$
- Including DG, gluon distribution is closer to global fits.

<sup>1</sup> S. Xu, CM, X. Zhao, Y. Li, J. P. Vary, 2209.08584 [hep-ph].

# Valence Quark Ratio in Proton

Ratio of structure functions

$$\frac{F_2^n(x)}{F_2^p(x)} \underset{x \gtrsim 0.2}{=} \frac{1 + 4d_v(x)/u_v(x)}{4 + d_v(x)/u_v(x)}$$



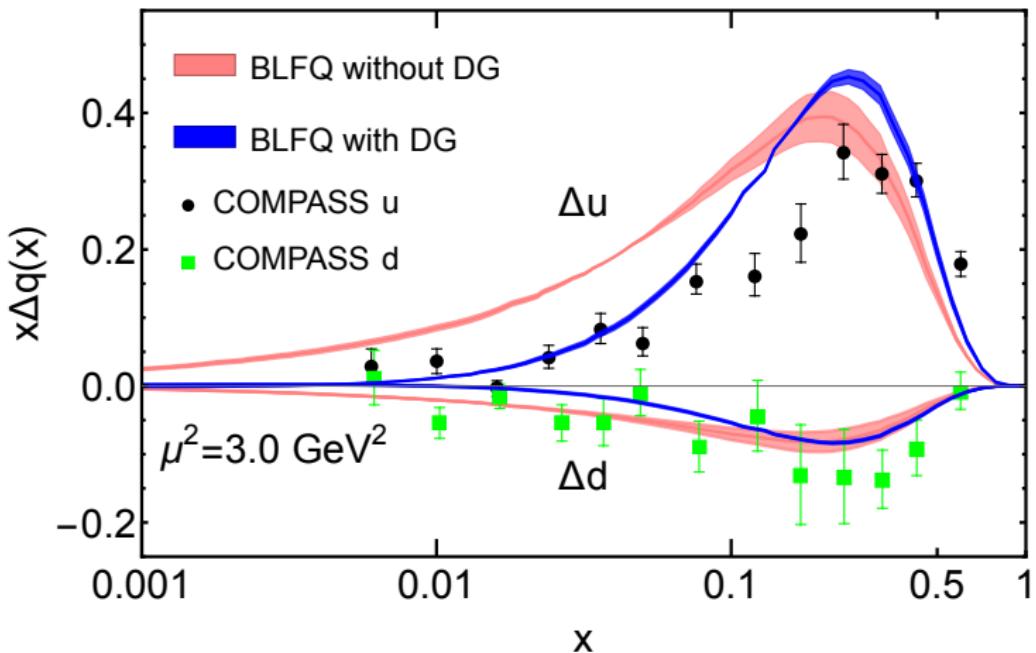
$0.225 \pm 0.025$   
 $0.230 \pm 0.057$



<sup>1</sup>S. Xu, CM, X. Zhao, Y. Li, J. P. Vary, 2209.08584 [hep-ph].

<sup>2</sup>Cui, Gao, Binosi, Chang, Roberts and Schmidt, Chin. Phys. Lett. **39**, 041401 (2022).

## Quark Helicity PDFs

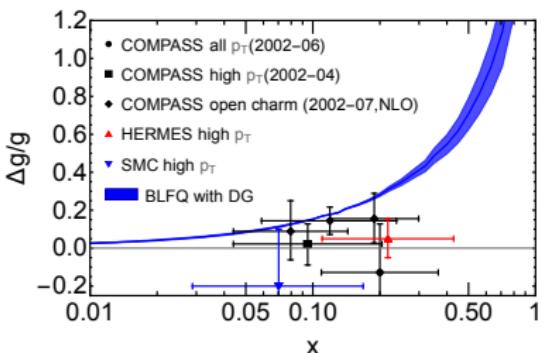
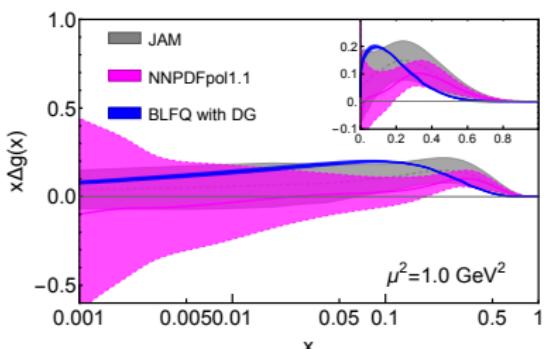


- Distributions improve at small  $x$  and  $x > 0.5$ .
- $\frac{1}{2}\Sigma_u = 0.438 \pm 0.004$ , strongly dominates over  $\frac{1}{2}\Delta\Sigma_d = -0.080 \pm 0.002$ .

<sup>1</sup>S. Xu, CM, X. Zhao, Y. Li, J. P. Vary, 2209.08584 [hep-ph].

<sup>2</sup>M. G. Alekseev et al. (COMPASS Collaboration), Phys. Lett. B 693, 227 (2010).

# Gluon Helicity Distribution



- $\Delta G = \int_0^1 dx \Delta g(x) = 0.131 \pm 0.003$ , is sizeable to the proton spin.
- PHENIX Collaboration:  $\Delta G^{[0.02, 0.3]} = 0.2 \pm 0.1$

PRL 103 (2009) 012003

<sup>1</sup> S. Xu, CM, X. Zhao, Y. Li, J. P. Vary, 2209.08584 [hep-ph].

<sup>2</sup> N. Sato *et al.* [JAM], PRD93 (2016); E. R. Nocera *et al.* [NNPDF], NPB 887 (2014).

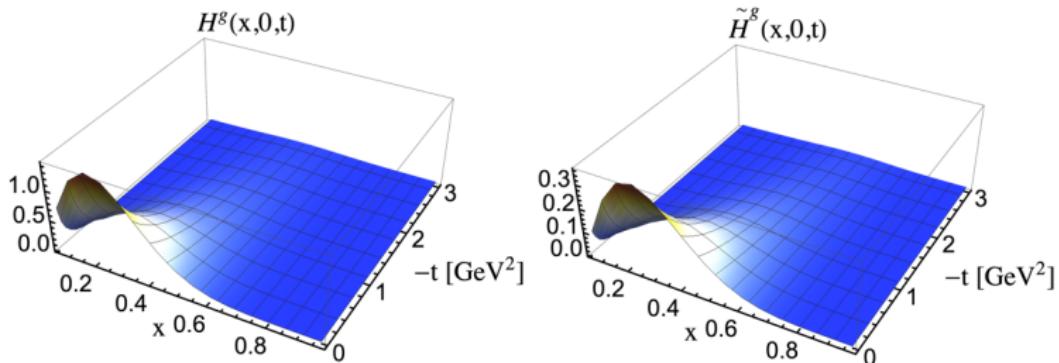
## Gluon GPDs

Meissner, Metz and Goeke, PRD 76, 034002 (2007)



$$F^g(x, \Delta; \lambda, \lambda') = \frac{1}{2P^+} \bar{u}(p', \lambda') \left( \gamma^+ H^g(x, \xi, t) + \frac{i\sigma^{+\mu} \Delta_\mu}{2M} E^g(x, \xi, t) \right) u(p, \lambda),$$
$$\tilde{F}^g(x, \Delta; \lambda, \lambda') = \frac{1}{2P^+} \bar{u}(p', \lambda') \left( \gamma^+ \gamma_5 \tilde{H}^g(x, \xi, t) + \frac{\Delta^+ \gamma_5}{2M} \tilde{E}^g(x, \xi, t) \right) u(p, \lambda).$$

## Preliminary

<sup>1</sup> B. Lin, S. Xu, CM, X. Zhao, et. al., work in progress.

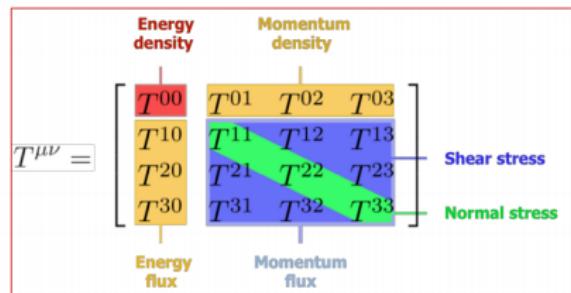
## Nucleon Gravitational Form Factors



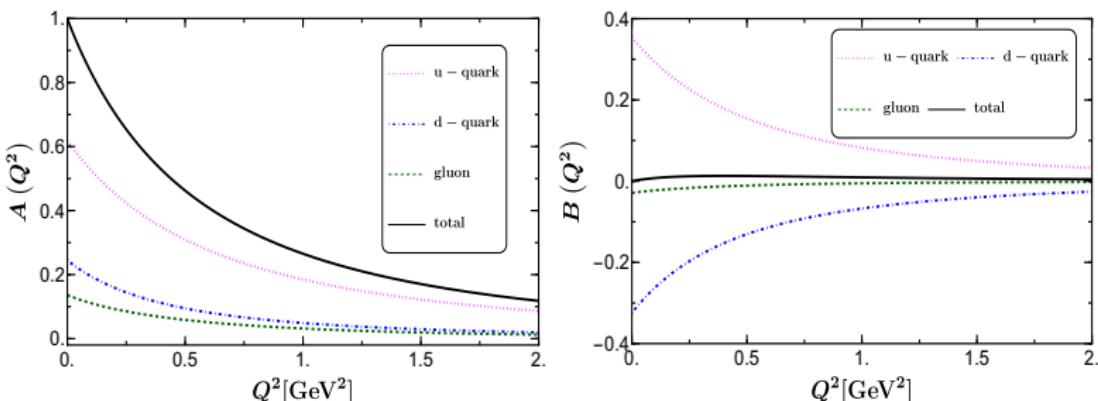
Nucleon scattering by the classical gravitational field is described by the gravitational (energy momentum tensor) form factors (GFFs).

$$\begin{aligned} \langle P' | T_i^{\mu\nu}(0) | P \rangle = \bar{U}' & \left[ -B_i(q^2) \frac{\bar{P}^\mu \bar{P}^\nu}{M} + (A_i(q^2) + B_i(q^2)) \frac{1}{2} (\gamma^\mu \bar{P}^\nu + \gamma^\nu \bar{P}^\mu) \right. \\ & \left. + C_i(q^2) \frac{q^\mu q^\nu - q^2 g^{\mu\nu}}{M} + \bar{C}_i(q^2) M g^{\mu\nu} \right] U \end{aligned}$$

- Matrix elements of the energy momentum tensor (EMT) contain fundamental information about various mechanical properties.



<sup>1</sup> Adam Freese's talk on 19th September, 3:45 PM.

Preliminary Results :  $A_q(Q^2)$  and  $B_q(Q^2)$ 

- $A(Q^2)$  and  $B(Q^2)$  are extracted from the  $T^{++}$  component.
- Spin sum rule:  $J^i = \frac{1}{2} (A^i(0) + B^i(0))$

$$\sum_i A^i(0) = 1 \text{ and } \sum_i B^i(0) = 0$$

<sup>1</sup>S. Nair *et. al.*, in preparation.

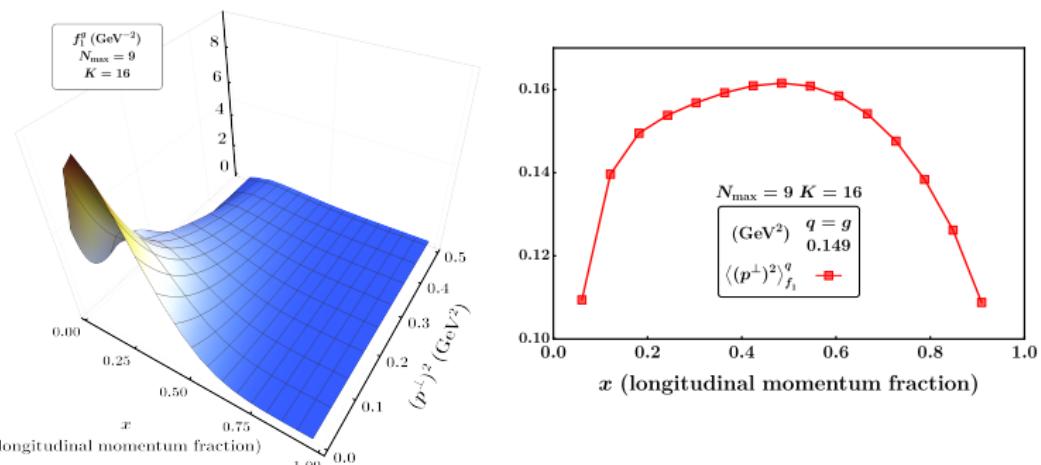
## Gluon TMDs

Meissner, Metz and Goeke, PRD 76, 034002 (2007)



$$\Phi^g(x, \vec{k}_\perp; S) = f_1^g(x, \vec{k}_\perp^2) - \frac{\epsilon_\perp^{ij} k_\perp^i S_\perp^j}{M} f_{1T}^{\perp g}(x, \vec{k}_\perp^2).$$

## Preliminary

<sup>1</sup> Z. Hu, et. al., work in progress.

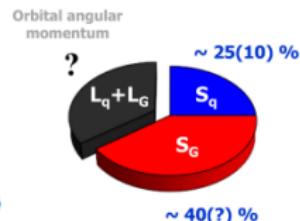
## GTMDs &amp; OAM



Generalized Transverse-Momentum Parton Distribution functions

$$W_{\lambda\lambda'}^{[\gamma^+]}(P, x, \vec{k}_\perp, \Delta) = \frac{1}{2} \int \frac{dz^- d^2\vec{z}_\perp}{(2\pi)^3} e^{ik\cdot z} \left\langle p', \lambda' \left| \bar{\psi}\left(-\frac{1}{2}z\right) \gamma^+ \psi\left(\frac{1}{2}z\right) \right| p, \lambda \right\rangle$$

$$W_{\lambda\lambda'}^{[\delta^{ij}]}(P, x, \vec{k}_\perp, \Delta) = \frac{1}{x P^+} \int \frac{dz^- d^2\vec{z}_\perp}{(2\pi)^3} e^{ik\cdot z} \left\langle p', \lambda' \left| G^{+i}\left(-\frac{1}{2}z\right) G^{+i}\left(\frac{1}{2}z\right) \right| p, \lambda \right\rangle$$

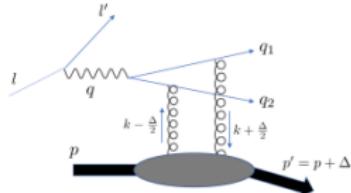


Parameterization:

$$\begin{aligned} W_{\lambda\lambda'}^{[\gamma^+]}(P, x, \vec{k}_\perp, \Delta) &= W_{\lambda\lambda'}^{[\delta^{ij}]}(P, x, \vec{k}_\perp, \Delta) \\ &= \frac{1}{2M} \bar{u}(p', \lambda') \left[ F_{1,1} + \frac{i\sigma^{j+}}{p^+} (k_\perp^j F_{1,2} + \Delta_\perp^j F_{1,3}) + i \frac{\sigma^{ij} k_\perp^i \Delta_\perp^j}{M^2} F_{1,4} \right] u(p, \lambda) \end{aligned}$$

 $F_{1,4}$  is related to the orbital angular momentum

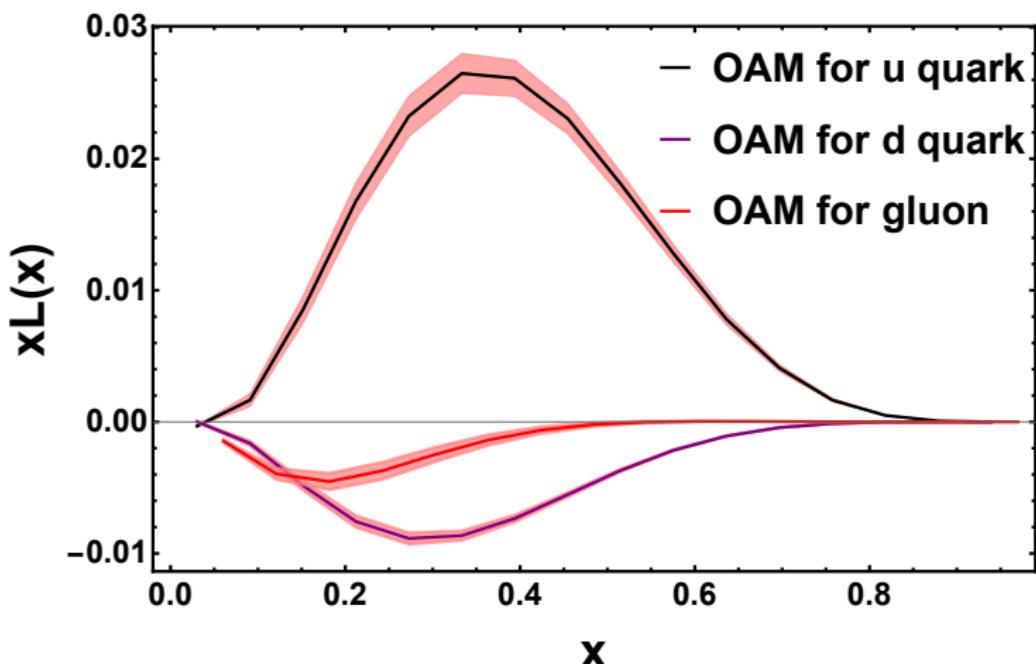
$$L_{q,g}(x) = - \int d^2k_\perp \frac{k_\perp^2}{M^2} F_{1,4}(x, k_\perp, \Delta_\perp = 0)$$




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<sup>1</sup> S. Bhattacharya, R. Boussarie and Y. Hatta, arXiv-hep:2201.08709 (2022)

## Preliminary Results of OAM



- Final remark :  $\frac{1}{2}\Delta\Sigma_q = 0.36$ ;  $\Delta\Sigma_g = 0.13$ ;  $l_q = 0.02$ ;  $l_g = -0.01$

## CONCLUSIONS



- Discussed structure of proton from eigenstates of LF effective Hamiltonians
- Considered  $|qqq\rangle$  and  $|qqqg\rangle$ .
- **LF Hamiltonian**  $\Rightarrow$  **Wavefunctions**  $\Rightarrow$  **Observables**.
- Provides good description of data/global fits for various observables.
- Discussed gluon distributions of the nucleon.
- With one dynamical gluon, the quark spin contributes 70%; the gluon spin plays a substantial role (26%) in understanding the nucleon spin.
- *This is not a complete picture ... long way to go.*

Enormous amount of possibilities with future EICs . . . . Thank You