

Deeply Virtual Compton Scattering with CLAS12^{and beyond}

Adam HOBART

APCTP Focus Program in Nuclear Physics 2022, 18 – 23 July



GPDs

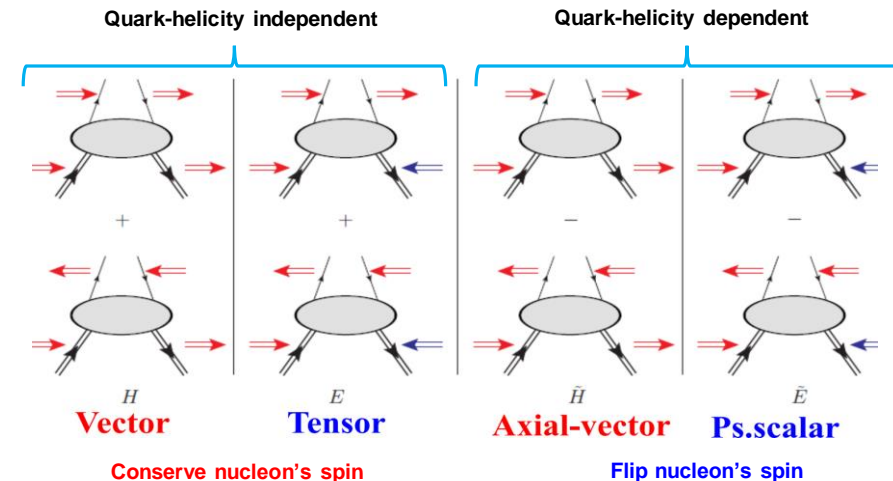
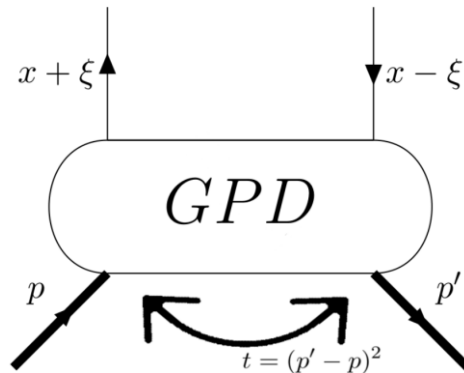
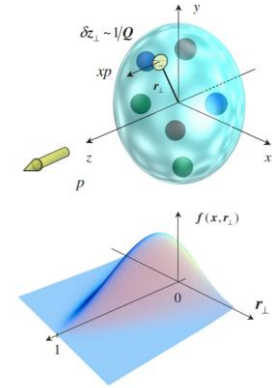
Belitsky, Radyushkin, Physics Reports, 2005

- QCD at low energies: non perturbative regime
 - Need **structure functions** to describe nucleon structure

GPDs

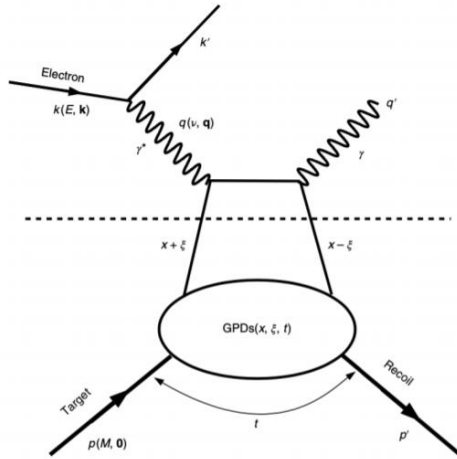
Correlation of transverse position and longitudinal momentum of partons in the nucleon & the spin structure - through Ji's sum rule X. Ji, Phy.Rev.Lett.78,610(1997)

- GPDs can be accessed through **exclusive leptonproduction reactions**
- At leading order QCD, chiral-even (quark helicity is conserved), quark sector: 4 **GPDs** for each quark flavor H, \tilde{H}, E and \tilde{E}
- GPDs depend on x, ξ and $t = (p' - p)^2$





Deeply Virtual Compton Scattering of leptons off nucleons



Invariants

$$Q^2 = -(k - k')^2$$

$$x_0 = Q^2/(2qp)$$

$$y = (qp)/(kp)$$

$$t = (q - q')^2$$

$$\sigma(eN \rightarrow eN\gamma) = \left[\text{DVCS} + \text{Bethe-Heitler (BH)} \right]^2$$

BH is purely electromagnetic and parametrised by FFs

- DVCS allows access to 4 complex GPDs-related quantities:
 - Compton Form Factors (x, ξ, t) (CFFs)

$$\mathcal{H} = \sum_q e_q^2 \left\{ i \pi [H^q(\xi, \xi, t) - H^q(-\xi, \xi, t)] + \mathcal{P} \int_{-1}^1 dx H^q(x, \xi, t) \left[\frac{1}{\xi - x} - \frac{1}{\xi + x} \right] \right\}$$

- x can not be accessed experimentally by DVCS: Models needed to map the x dependence



- Experimentally measured observables:
 - Contribution from both the DVCS and BH amplitudes
 - In order to be sensitive to the DVCS-BH interference part (linear in CFFs) we should have:
 - Beam polarized and/or target polarized
 - Each measured asymmetry gives access to combinations of CFFs
 - The separation of CFFs requires the measurement of several observables
 - Depending on the target (proton or neutron): different sensitivity to the CFFs (GPDs)
 - The flavor separation of GPDs requires measurements on both nucleons

$$(H, E)_u(\xi, \xi, t) = \frac{9}{15} \left[4(H, E)_p(\xi, \xi, t) - (H, E)_n(\xi, \xi, t) \right]$$

$$(H, E)_d(\xi, \xi, t) = \frac{9}{15} \left[4(H, E)_n(\xi, \xi, t) - (H, E)_p(\xi, \xi, t) \right]$$



Polarized beam, unpolarized target

$$\Delta\sigma_{LU} \sim \sin(\phi) \Im\{F_1 \mathbf{H} + \xi(F_1 + F_2) \tilde{\mathbf{H}} - k F_2 \mathbf{E} + \dots\}$$

Unpolarized beam, polarized target

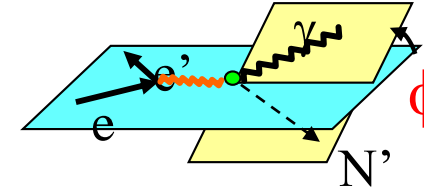
$$\Delta\sigma_{UL} \sim \sin(\phi) \Im\left\{F_1 \tilde{\mathbf{H}} + \xi(F_1 + F_2) \left(\mathbf{H} + \frac{x_b}{2} \mathbf{E}\right) - \xi k F_2 \tilde{\mathbf{E}}\right\}$$

polarized beam, longitudinal polarized target

$$\Delta\sigma_{LL} \sim (A + B \cos(\phi)) \Re\{F_1 \tilde{\mathbf{H}} + \xi(F_1 + F_2) \left(\mathbf{H} + \frac{x_b}{2} \mathbf{E}\right) + \dots\}$$

unpolarized beam, transverse polarized target

$$\Delta\sigma_{UT} \sim \cos(\phi) \sin(\phi_s - \phi) \Im\{k(F_2 \mathbf{H} - F_1 \mathbf{E}) + \dots\}$$



Observable	Proton	Neutron
$\Delta\sigma_{LU}$	$\Im\{\mathbf{H}_p, \tilde{\mathbf{H}}_p, E_p\}$	$\Im\{H_n, \tilde{H}_n, \mathbf{E}_n\}$
$\Delta\sigma_{UL}$	$\Im\{\mathbf{H}_p, \tilde{\mathbf{H}}_p\}$	$\Im\{\mathbf{H}_n, E_n\}$
$\Delta\sigma_{LL}$	$\Re\{\mathbf{H}_p, \tilde{\mathbf{H}}_p\}$	$\Re\{\mathbf{H}_n, E_n\}$
$\Delta\sigma_{UT}$	$\Im\{\mathbf{H}_p, E_p\}$	$\Im\{\mathbf{H}_n\}$

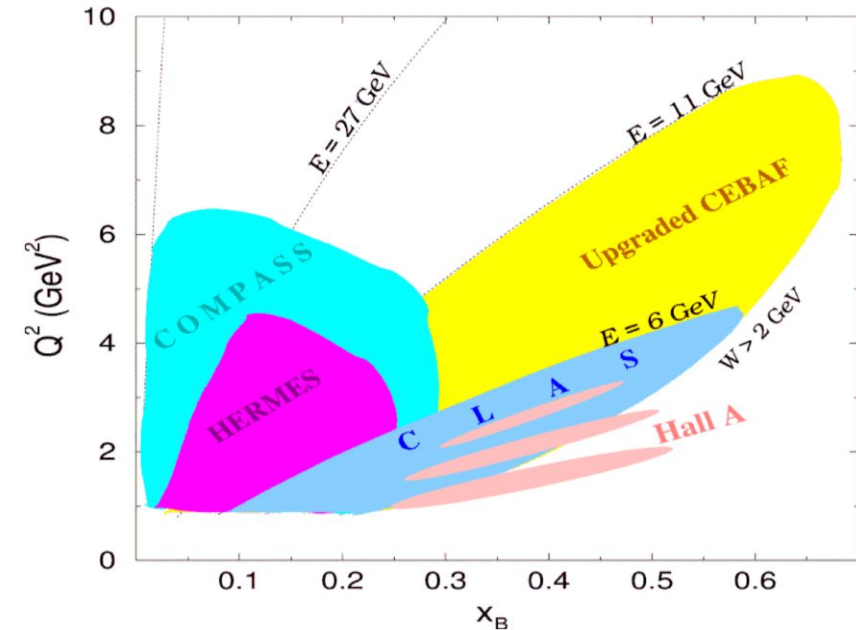
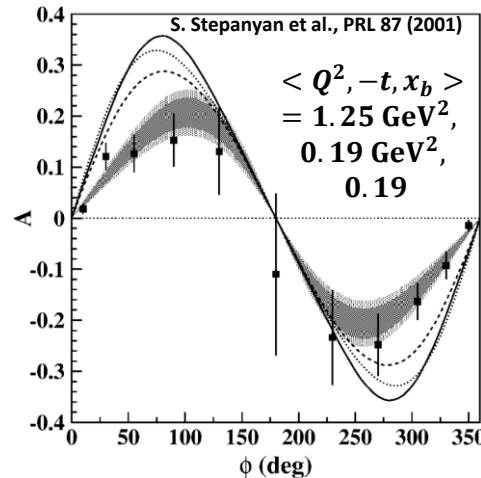


GPD-aimed experiments

- JLAB

- Hall A:
 - Cross sections
 - Beam-polarised cross section differences
- Hall B (CLAS/CLAS12):
 - Beam and target spin asymmetries
 - Cross section measurements over large phase-space acceptance

$ep \rightarrow epX$, from CLAS data:
First observation of DVCS-BH
interference

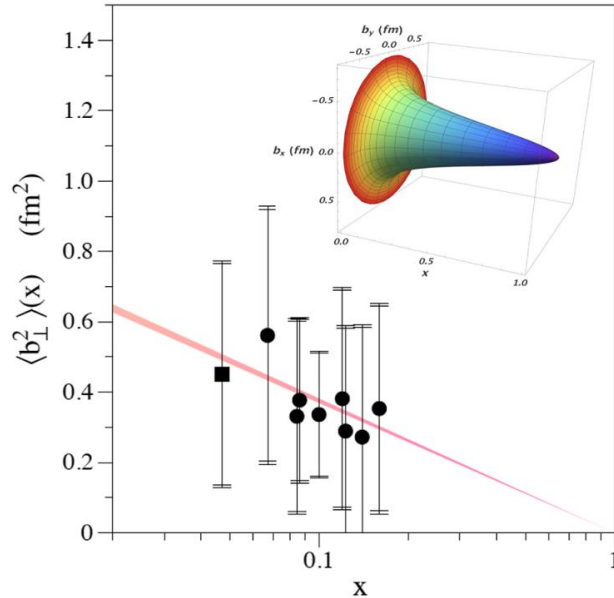




What did we learn from JLAB data at ~6 GeV?

- Tomography of the proton

- Obtained from local fits to HERMES, CLAS, and Hall-A data ($\mathcal{S}H$ + model dependent assumptions for x dependence)

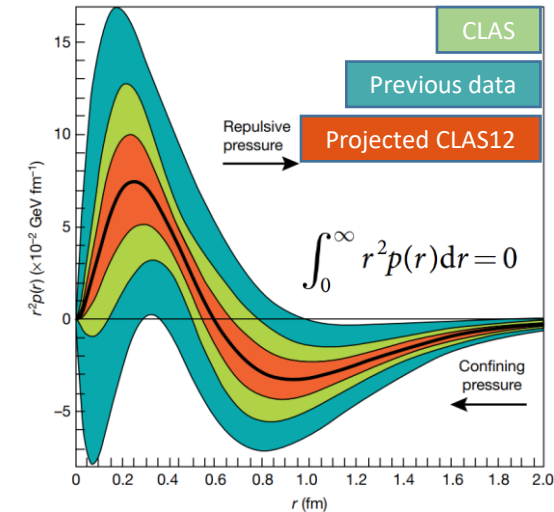
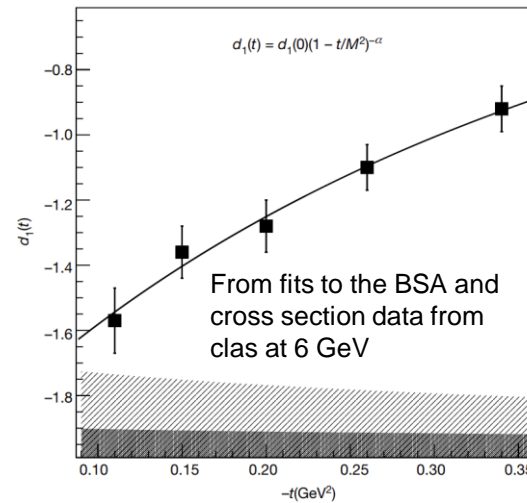


R. Dupré, M. Guidal, M.Vanderhaeghen, PRD95, 011501 (2017)

- The pressure distribution inside the proton

- Gravitational form factor

- $\int x H(x, \xi, t) dx = M_2(t) + \frac{4}{5} \xi^2 d_1(t)$
- $d_1(t) \propto \int \frac{j_0(r\sqrt{-t})}{2t} \rho(r) d^3r$



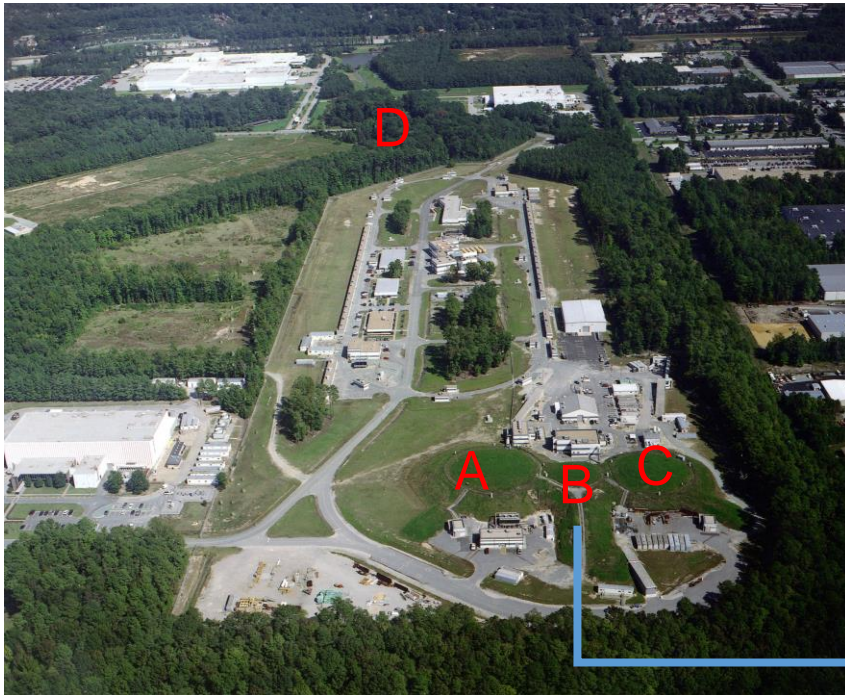
V. Burkert, L. Elouadrhiri, F.X. Girod, Nature 557, 396-399 (2018)



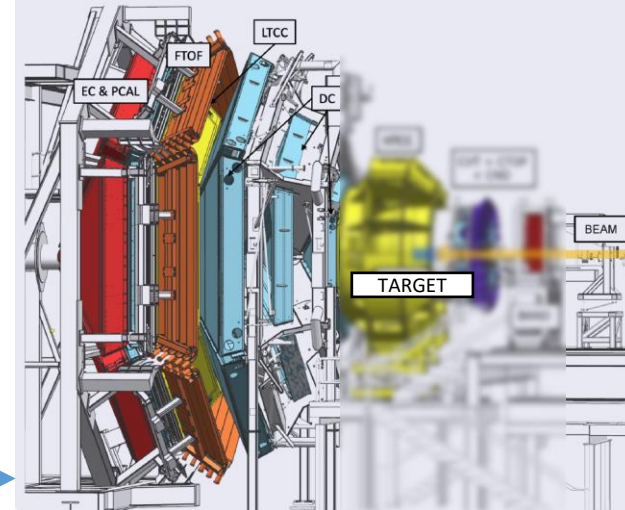
The CEBAF at Jefferson Laboratory

Continuos Electron Beam Accelerator Facility

- Up to 12 GeV electrons



At ~6 GeV



CLAS

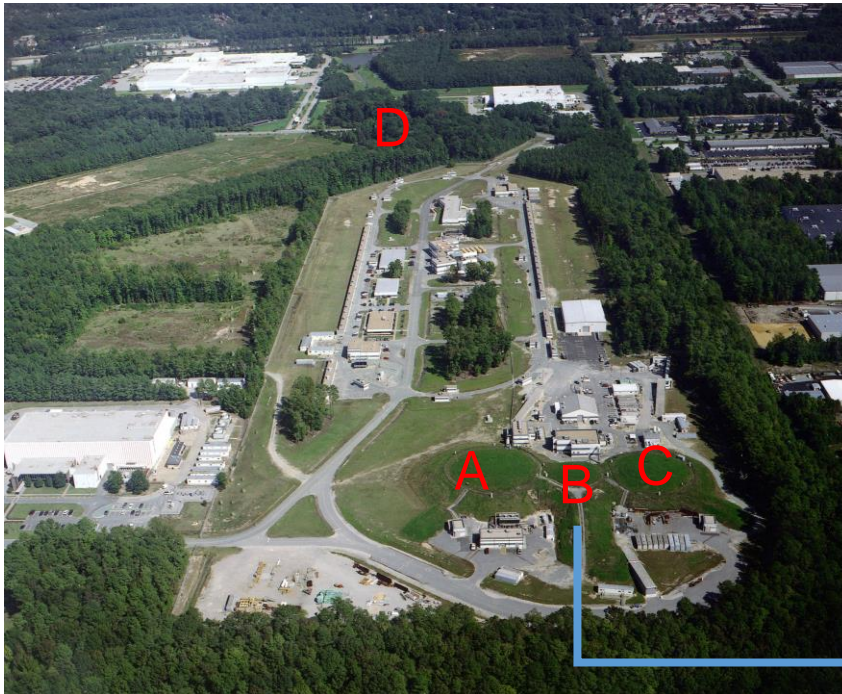


The CEBAF at Jefferson Laboratory

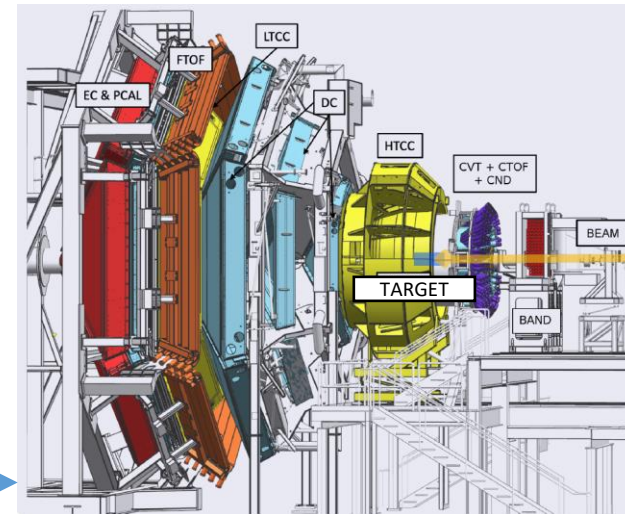
Continuos Electron Beam Accelerator Facility

- Up to 12 GeV electrons

The 12 GeV energy upgrade was associated with detector upgrades



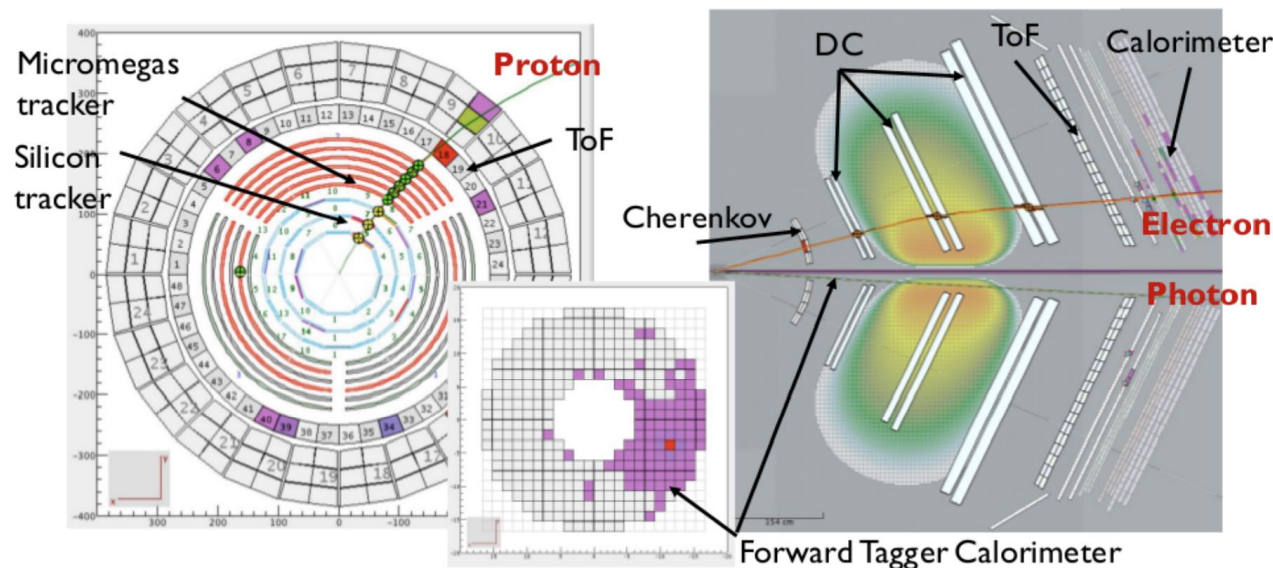
At ~11 GeV



CLAS12

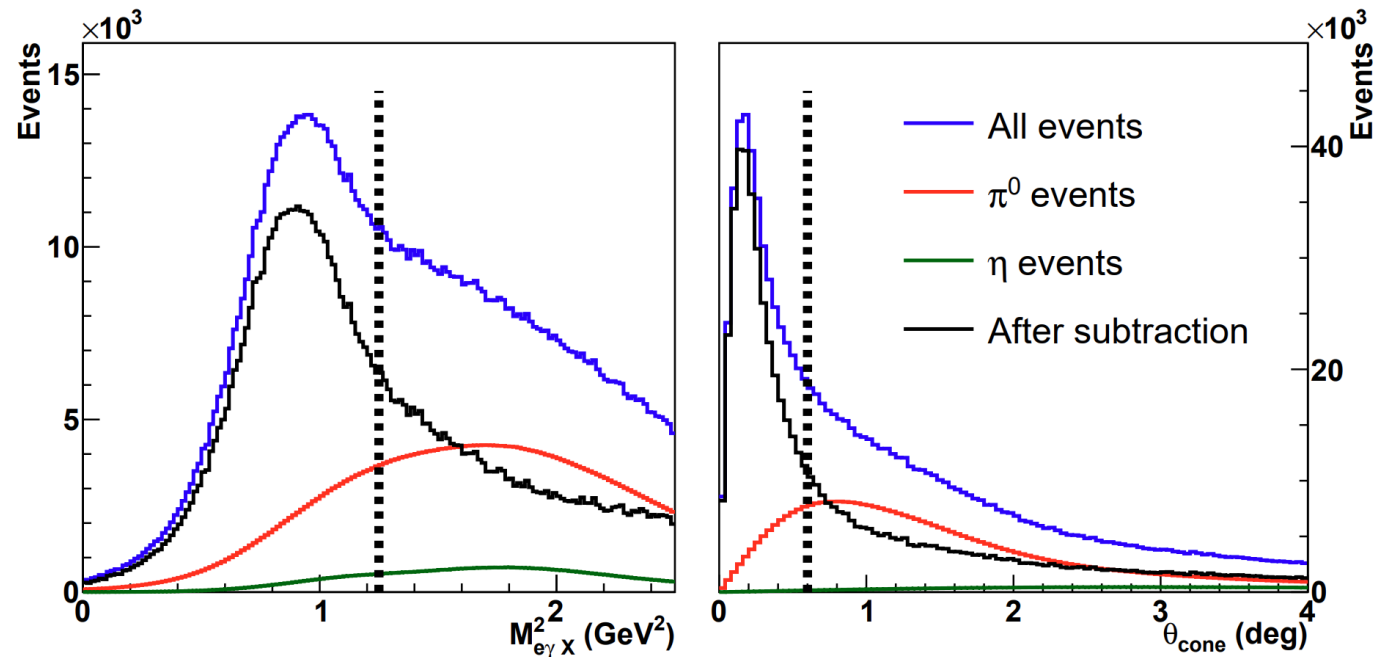


- A **10.6 GeV** electron beam
 - With an average **polarization** of **86%**
 - Scattering off an **unpolarized LH2 target** of **5 cm** length
- The **full exclusivity** of the event is insured by:
 - **Electron detection**: Cerenkov detector, drift chambers and electromagnetic calorimeter
 - **Photon detection**: sampling calorimeter or a small PbWO₄-calorimeter close to the beamline
 - **Proton detection**: Silicon and Micromegas detector





- Exclusivity is enforced by cutting on 5 variables:
 - The missing mass $ep \rightarrow e\gamma pX$
 - The missing mass $ep \rightarrow e\gamma X$
 - The missing energy
 - The missing transverse momentum
 - The cone angle (angle between detected photon and expected photon assuming exclusivity)



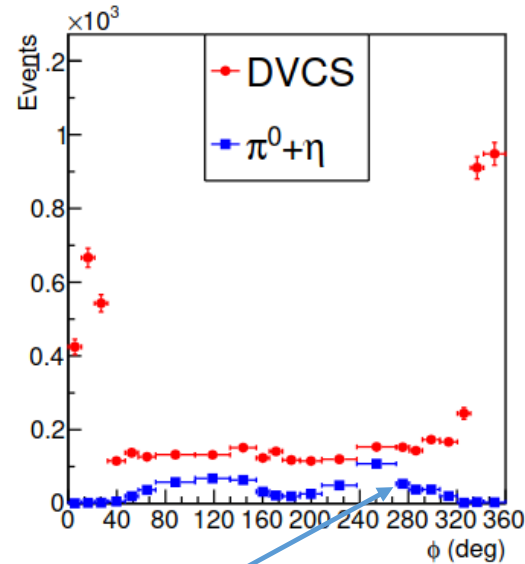
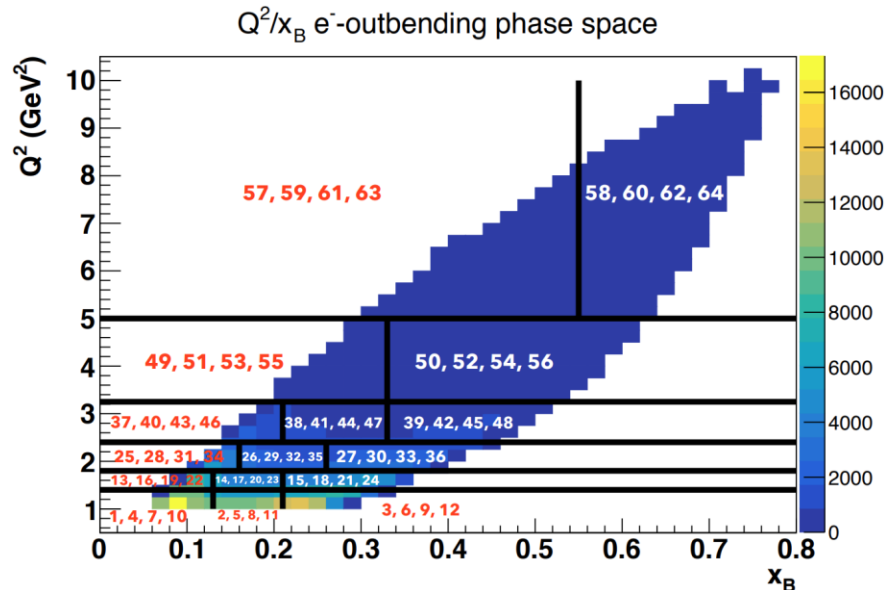


- For each Q^2/x_b , 4 bins in t are defined
 - 64 (Q^2, x_b, t) kinematical bins
- Φ : adaptative binning to accommodate for the steep dependence of the cross section.

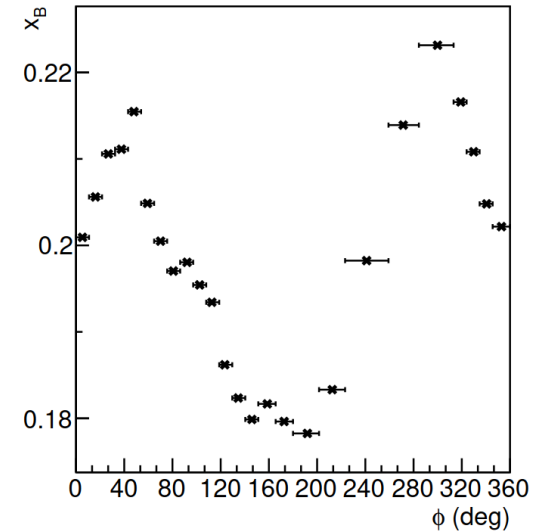
(Q^2, x_b, t) kinematics are Φ dependent

- Binning chosen to accommodate for this variation

Wide kinematical coverage



π^0/η background subtraction is insured using toy MC





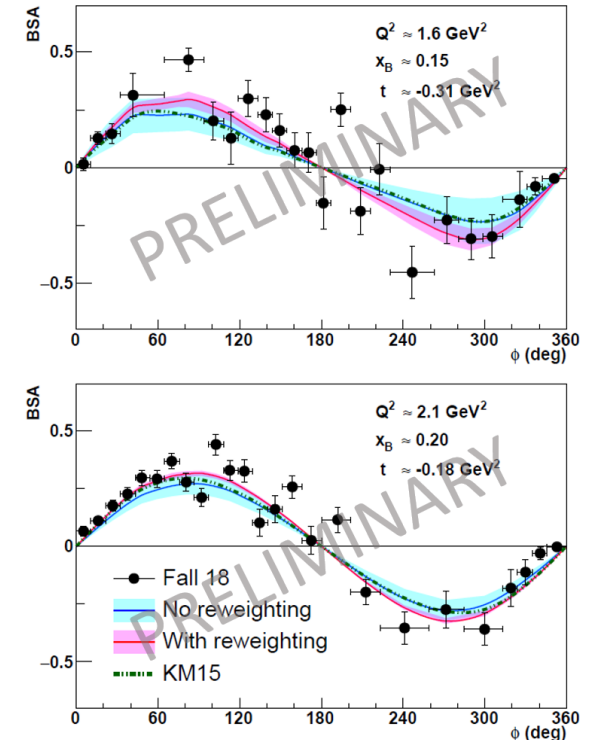
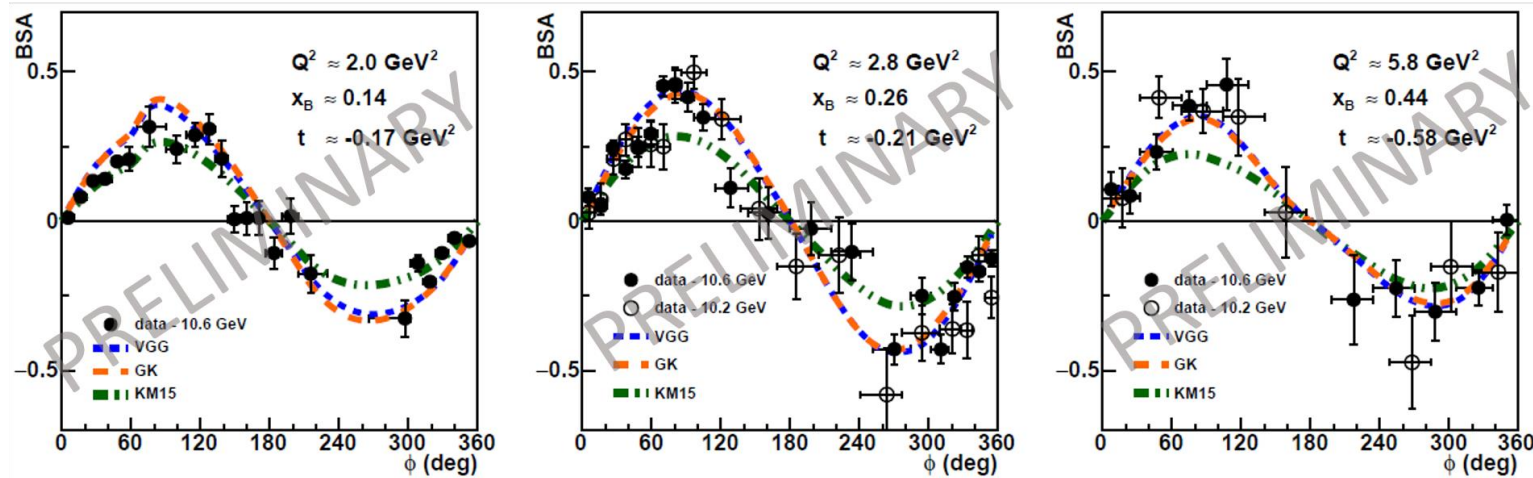
H. Moutarde, P. Sznajder, and J. Wagner, EPJC 79, 614 (2019)

- Deriving the mean and standard deviation of a 100 ANN-predictions produced by a global fit (PARTONS)
 - The new data are shown to be in good agreement
- Comparisons with KM15 and VGG/GK models
 - Data favors the KM15 model for most of the bins
 - At large x_b data favors the VGG/GK models

Kumericki, Kresimir and Müller, Dieter, EPJ Web of Conferences 112, 01012 (2016).

S. V. Goloskokov and P. Kroll, EPJC 65, 10.1140 (2009)

M. Vanderhaeghen, P. A. Guichon, and M. Guidal, Phys.Rev. D60, 094017 (1999)



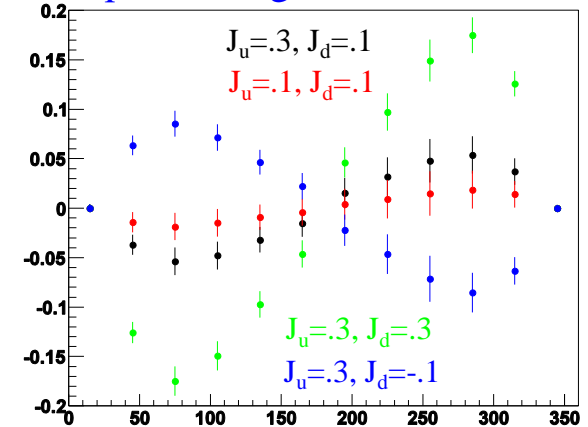


- Physics observable: Beam Spin Asymmetry BSA
 - Scattering off neutron (nDVCS): GPD E
 - Determination of Ji sum rule
 - Contribution of orbital angular momentum of quarks to the nucleon spin

$$\frac{1}{2} \int_{-1}^1 x dx (H(x, \xi, t=0) + E(x, \xi, t=0)) = J = \frac{1}{2} \Delta \Sigma + \Delta L$$

- Scattering off proton (pDVCS): GPD H
 - Quantify medium effects
 - Essential for the extraction of BSA of a “free” neutron (deconvoluting medium effect via comparison with DVCS on hydrogen target)

Model predictions (VGG)
for different values of
quarks' angular momentum



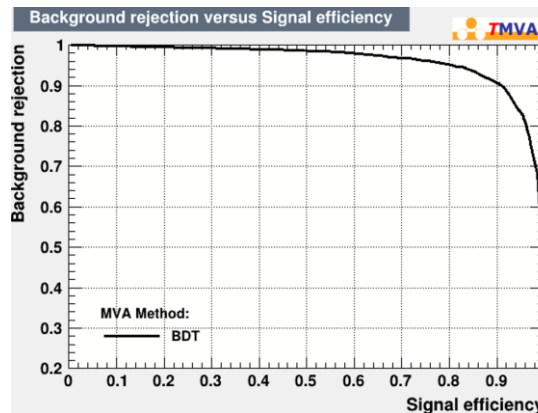
Polarized beam, unpolarized target

$$\Delta \sigma_{LU} \sim \sin(\phi) \Im \{ F_1 \mathbf{H} + \xi (F_1 + F_2) \tilde{\mathbf{H}} - k F_2 \mathbf{E} + \dots \}$$

Observable	Proton	Neutron
$\Delta \sigma_{LU}$	$\Im \{ \mathbf{H}_p, \tilde{\mathbf{H}}_p, E_p \}$	$\Im \{ H_n, \tilde{H}_n, \mathbf{E}_n \}$



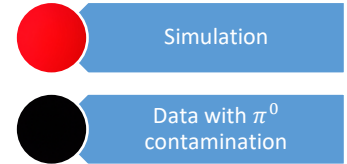
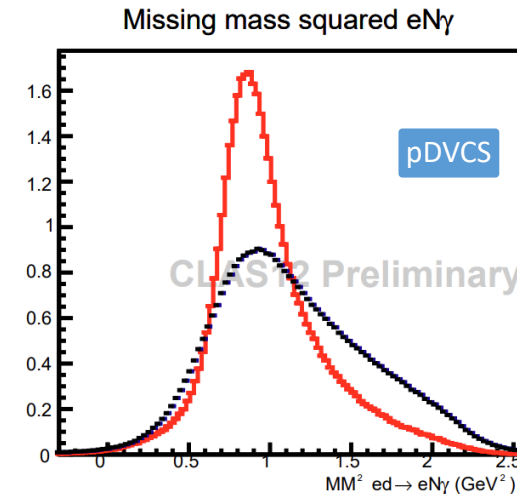
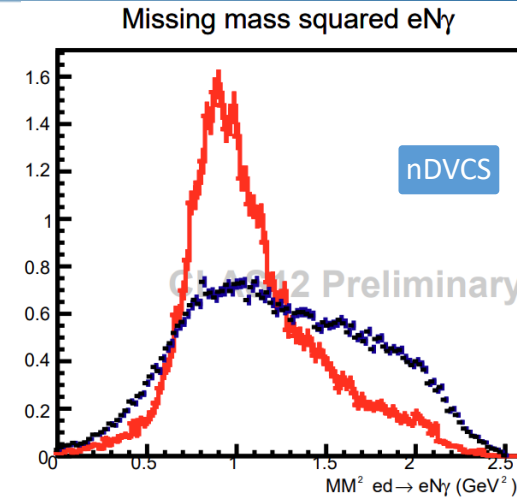
- A 10.6/10.4/10.2 GeV electron beam
 - With an average polarization of 86%
 - Scattering off an unpolarized Liquid Deuterium target of 5 cm length
- The exclusivity of the event is insured by:
 - Electron detection: Cerenkov detector, drift chambers and electromagnetic calorimeter
 - Photon detection: sampling calorimeter or a small PbWO₄-calorimeter close to the beamline
 - Proton detection: Silicon and Micromegas detector OR Neutron detection: Central Neutron Detector
- For Neutron Detection:
 - Machine Learning techniques are applied to improve the Identification and reduce charged particle contamination



The ROC curve of a Boosted Decision Tree classifier indicating the high performance when identifying and removing charged particles contamination from neutron sample



- The nDVCS (pDVCS) final state is selected with the following exclusivity criteria: (N:nucleon)
 - Missing mass
 - $e d \rightarrow e N \gamma X$
 - $e N \rightarrow e N \gamma X$
 - $e N \rightarrow e N X$
 - Missing momentum
 - $e d \rightarrow e N \gamma X$
 - $\Delta\Phi, \Delta t, \theta(\gamma, X)$
 - Difference between two ways of calculating Φ and t
 - Cone angle between measured and reconstructed photon
- Exclusivity selection is optimized with a 4-D χ^2 -like distribution including $\Delta\Phi, \Delta t, \theta(\gamma, X)$ and missing mass $e N \rightarrow e N X$

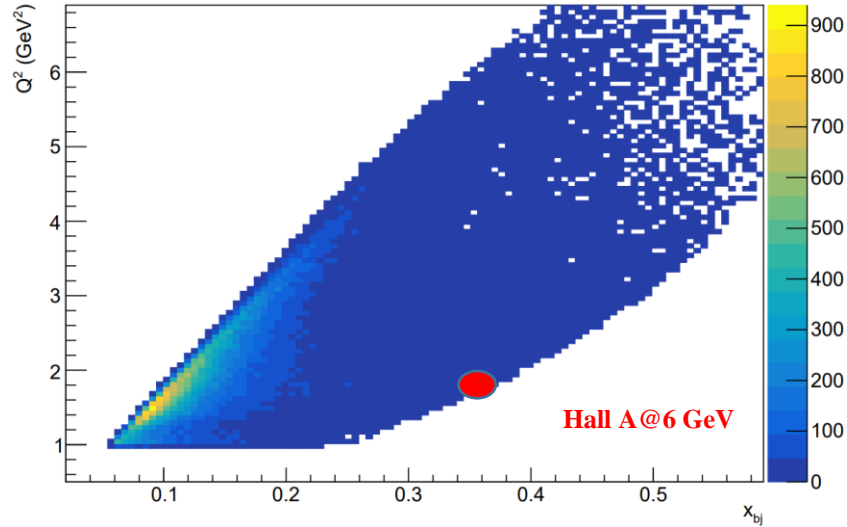


π^0 background contamination is estimated using simulations



First-time measurement of nDVCS with detection of the active neutron

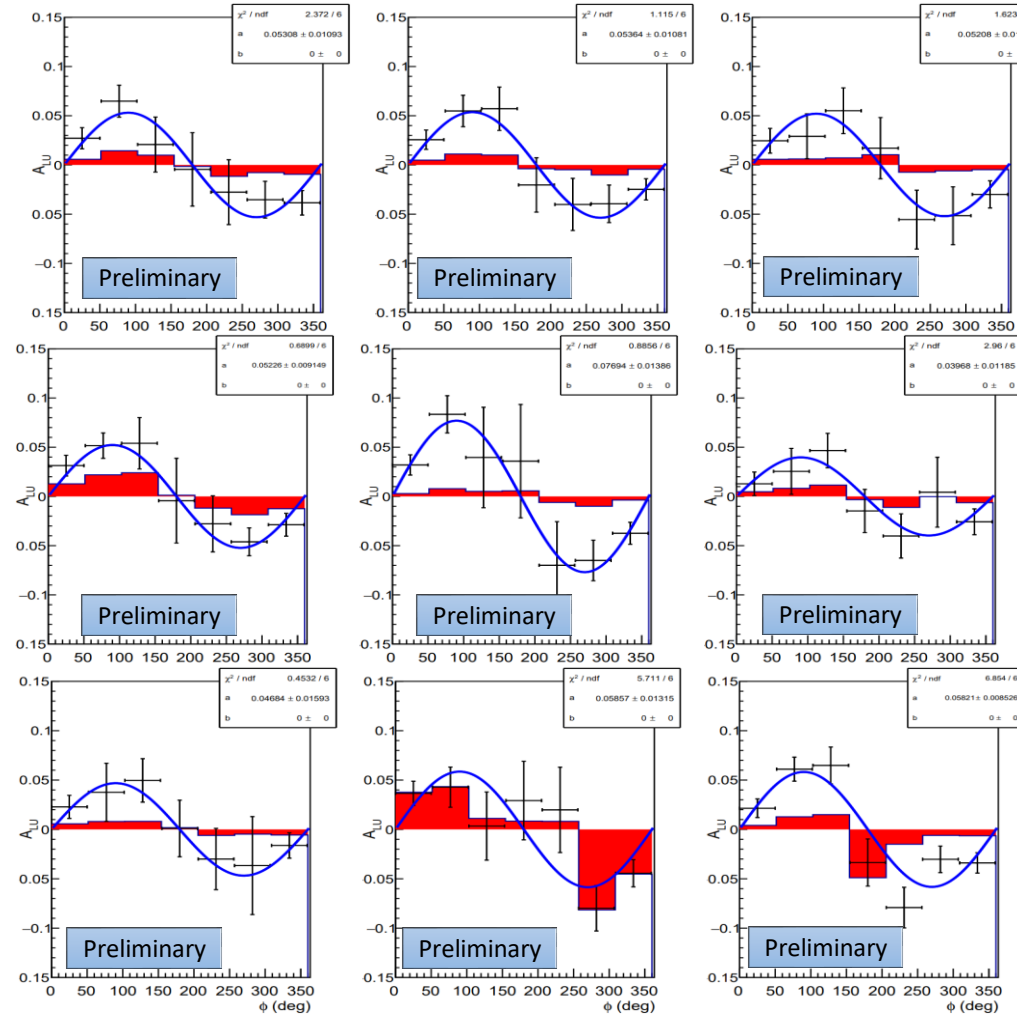
Q^2 vs x_b



Q^2 bins

x_b bins

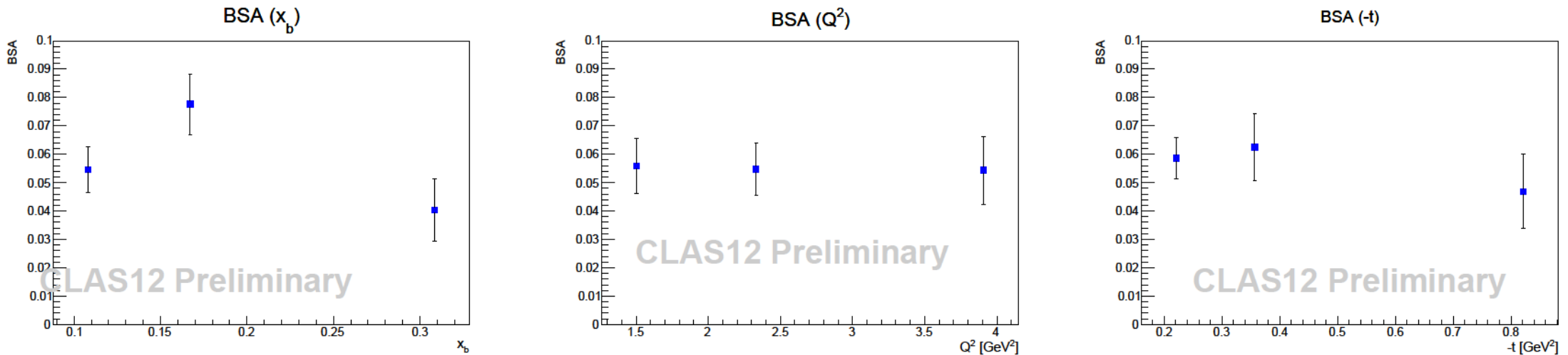
t bins



- Compared to the previous experiment, CLAS12 provides :
 - The possibility to scan the BSA of nDVCS on a wide phase space
 - The possibility to reach the high Q^2 high x_b region of the phase space
 - Exclusive measurement with the detection of the active neutron



$$\Delta\sigma_{LU} \sim \sin(\phi) \Im\{F_1 \mathbf{H} + \xi(F_1 + F_2) \tilde{\mathbf{H}} - k F_2 \mathbf{E} + \dots\}$$



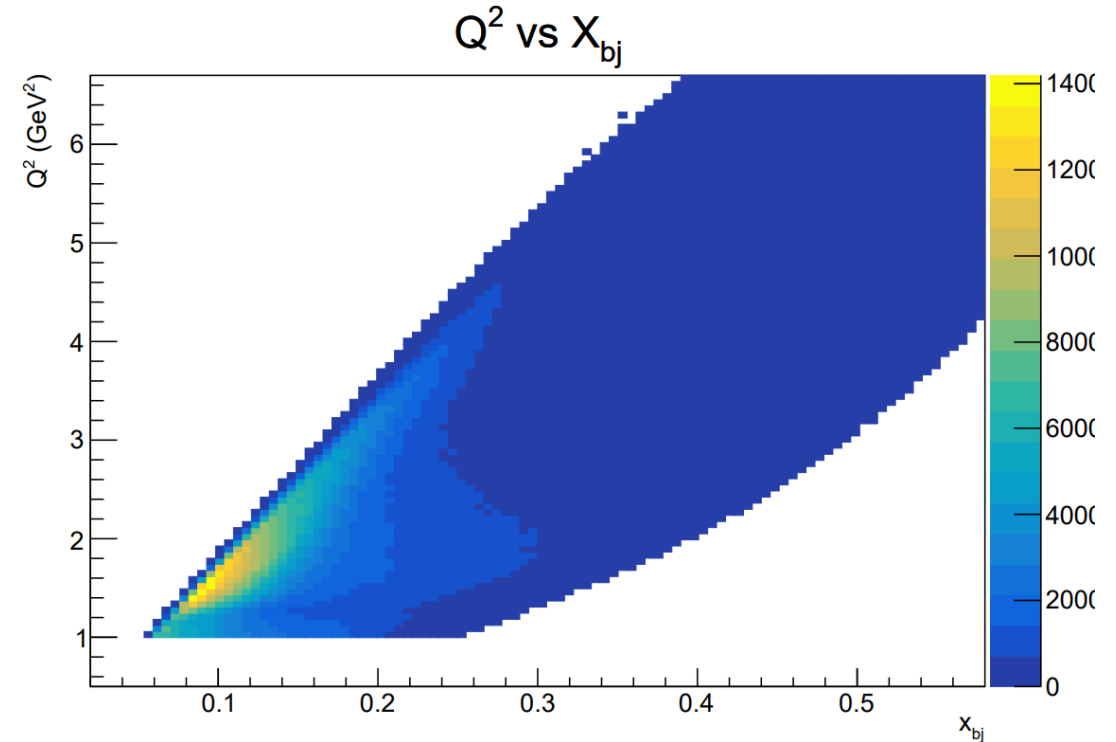
Variation of $\sin(\Phi)$ amplitude

Errors are statistical only



First-time measurement of incoherent
pDVCS on deuteron

Scattering off proton (pDVCS): GPD H
Quantify medium effects:
Essential for the extraction of BSA
of a “free” neutron (de-
convoluting medium effect via
comparison with DVCS on
hydrogen target)

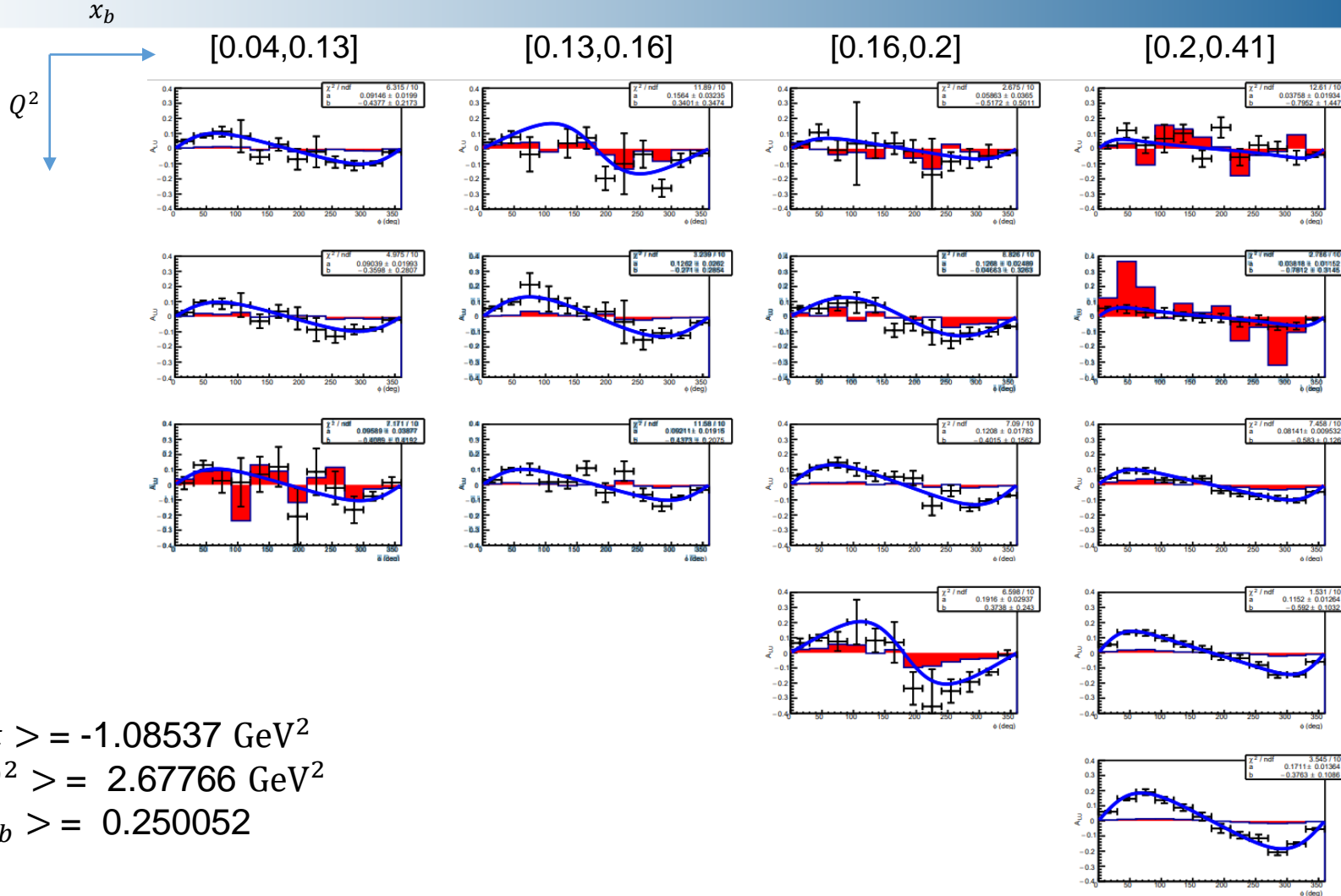




CLAS12: pDVCS with an unpolarized deuterium target

Under review

Preliminary



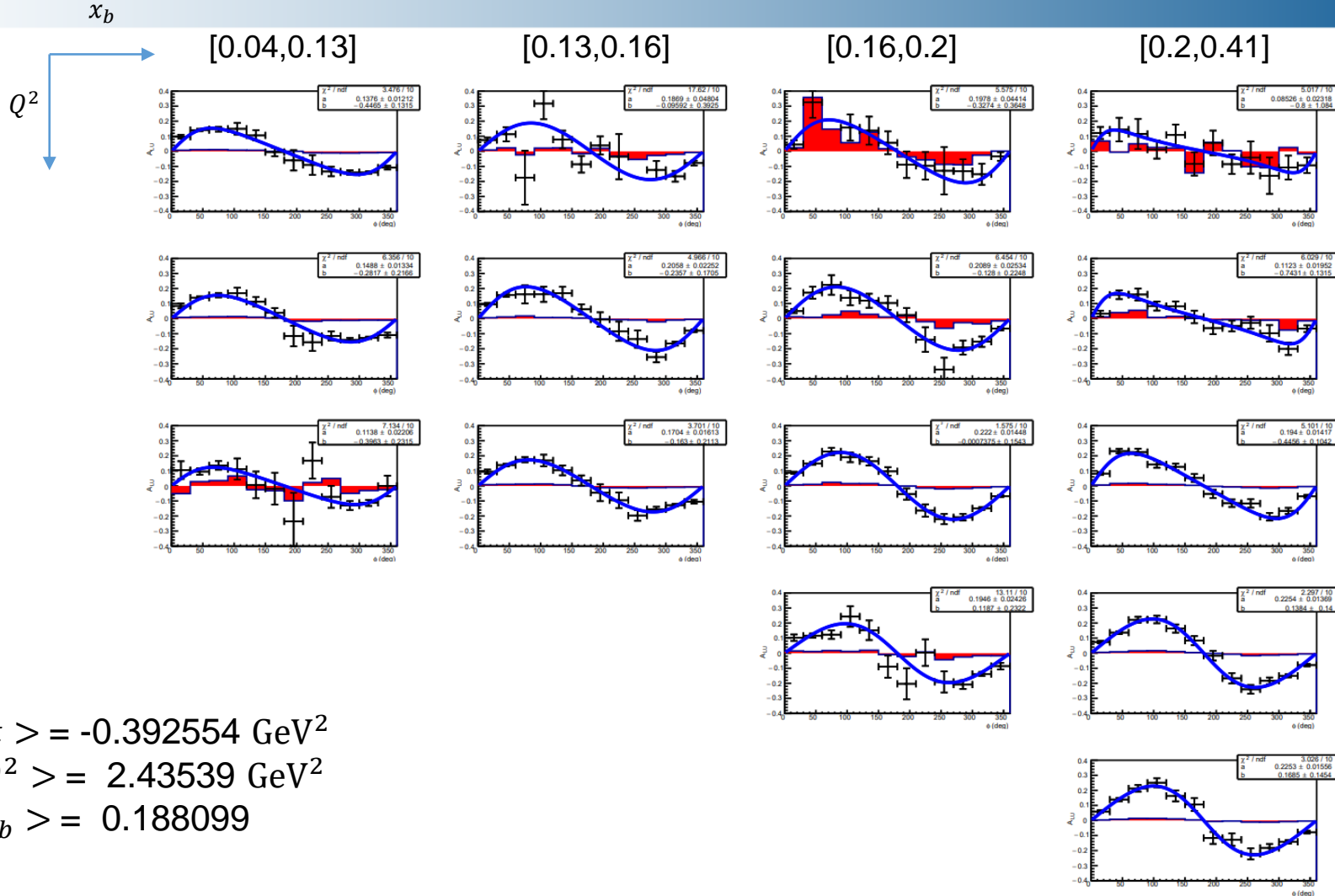
$$\begin{aligned} \langle t \rangle &= -1.08537 \text{ GeV}^2 \\ \langle Q^2 \rangle &= 2.67766 \text{ GeV}^2 \\ \langle x_b \rangle &= 0.250052 \end{aligned}$$



CLAS12: pDVCS with an unpolarized deuterium target

Under review

Preliminary

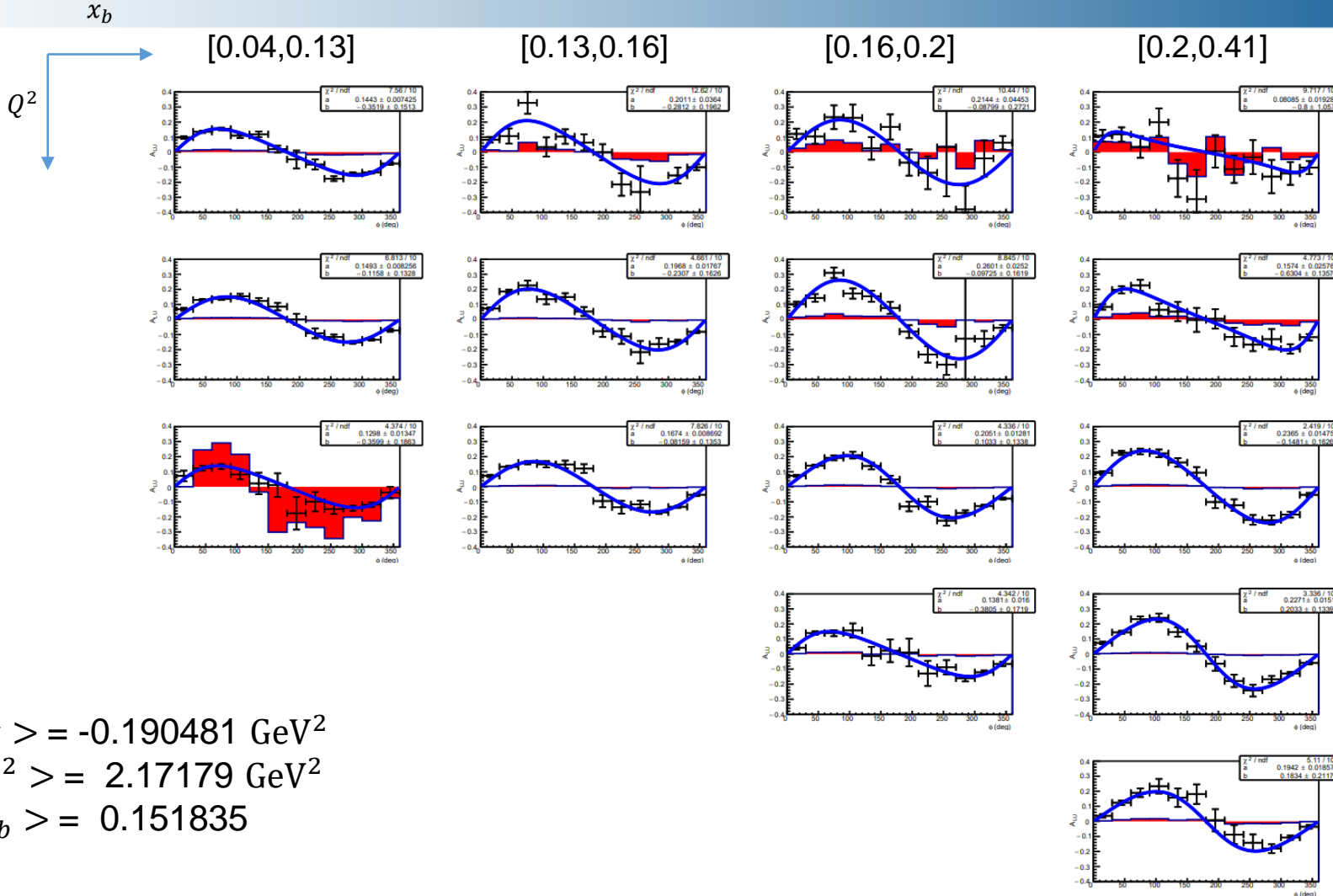


$$\begin{aligned} \langle t \rangle &= -0.392554 \text{ GeV}^2 \\ \langle Q^2 \rangle &= 2.43539 \text{ GeV}^2 \\ \langle x_b \rangle &= 0.188099 \end{aligned}$$

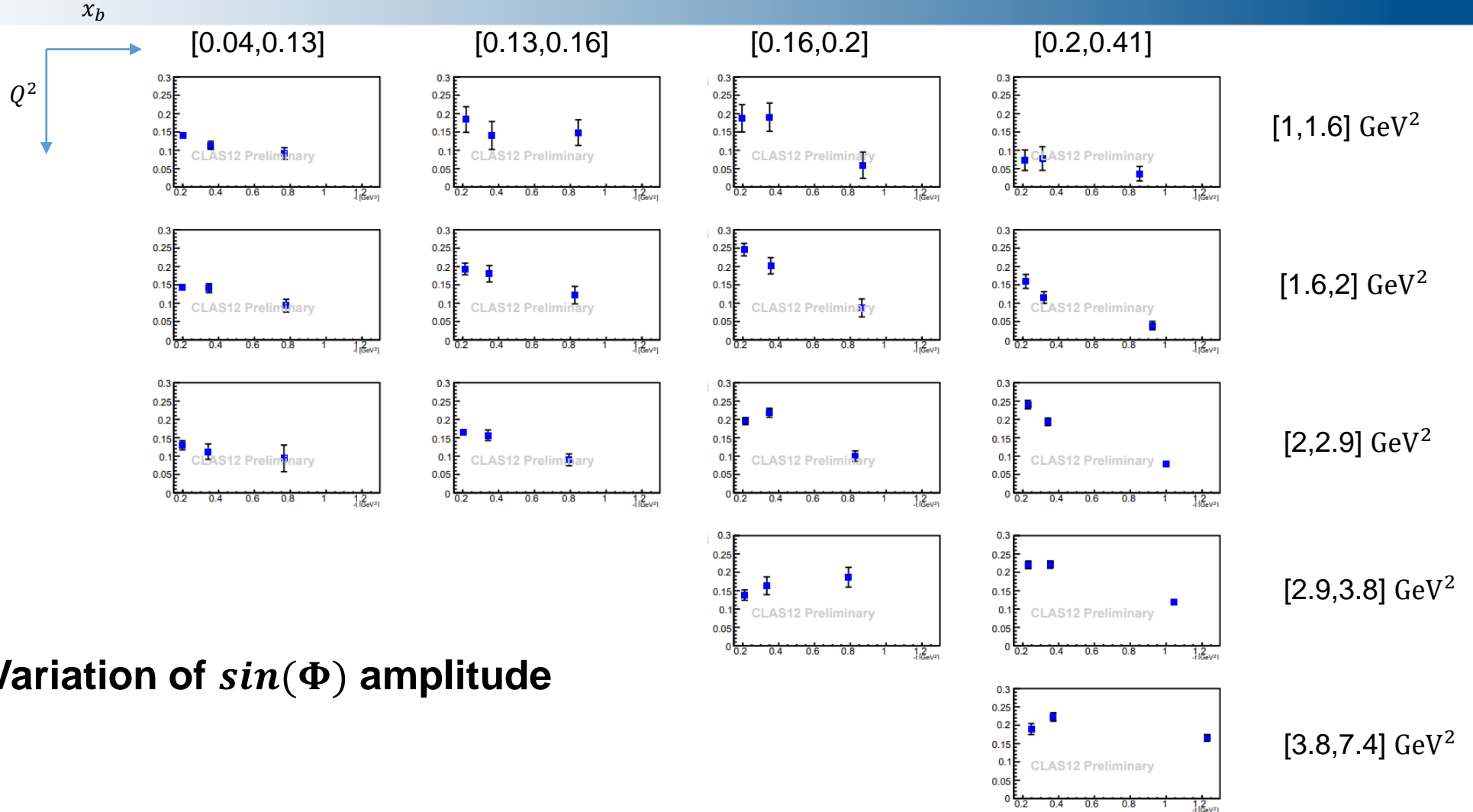


CLAS12: pDVCS with an unpolarized deuterium target

Under review



$$\begin{aligned} \langle t \rangle &= -0.190481 \text{ GeV}^2 \\ \langle Q^2 \rangle &= 2.17179 \text{ GeV}^2 \\ \langle x_b \rangle &= 0.151835 \end{aligned}$$





Observable (target)	CFF sensitivity	Status
ITSA(p), IDSA(p)	$\Im\{\mathbf{H}_p, \tilde{\mathbf{H}}_p\}, \Re\{\mathbf{H}_p, \tilde{\mathbf{H}}_p\}$	Data taking in progress
ITSA(n), IDSA(n)	$\Im\{\mathbf{H}_n\}, \Re\{\mathbf{H}_n\}$	Data taking in progress
tTSA(p)	$\Im\{\mathbf{H}_p\}, \Im\{\mathbf{E}_p\},$	Experiment foreseen for ~2025

- Analysis ongoing: Deuteron DVCS

- Physics observable to extract is Beam Spin Asymmetry (BSA)

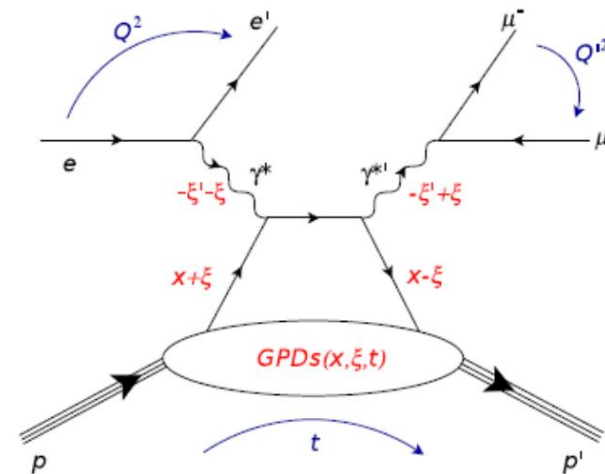
$$\Delta\sigma_{LU} \sim \sin(\phi) \Im \left\{ \frac{2G_1 \mathbf{H}_1 + (G_1 - \tau G_3)(\mathbf{H}_1 - 2\tau \mathbf{H}_3) + \frac{2}{3} \tau G_3 \mathbf{H}_5}{2G_1^2 + (G_1 - 2\tau G_3)^2} \right\}$$

- Spin 1: 9 GPDs for each quark flavor
 - BSA of DVCS off deuterium is sensitive to 3 of them



Upgrades?

- pDVCS measurements will benefit from energy upgrade
 - Wider kinematical coverage for a complementary scan of GPDs
 - CLAS12 acceptance is OK(?) Associate energy upgrade with detector upgrades(?)
- nDVCS will surely benefit from both luminosity and energy upgrade
 - nDVCS BSA is of the order of ~ 2 to ~ 4 %: high luminosity beneficial for precise measurements
 - Wider kinematical coverage: comparable to pDVCS, beneficial for simultaneous fits of GPDs (flavor separation)
- But also: scan x dependence of GPDs; luminosity upgrade crucial for DDVCS!
 - Full kinematics mapping of GPDs: virtuality of final photon allows a unique direct access to GPDs at $x \neq \pm \xi$
 - Improved detection of muons



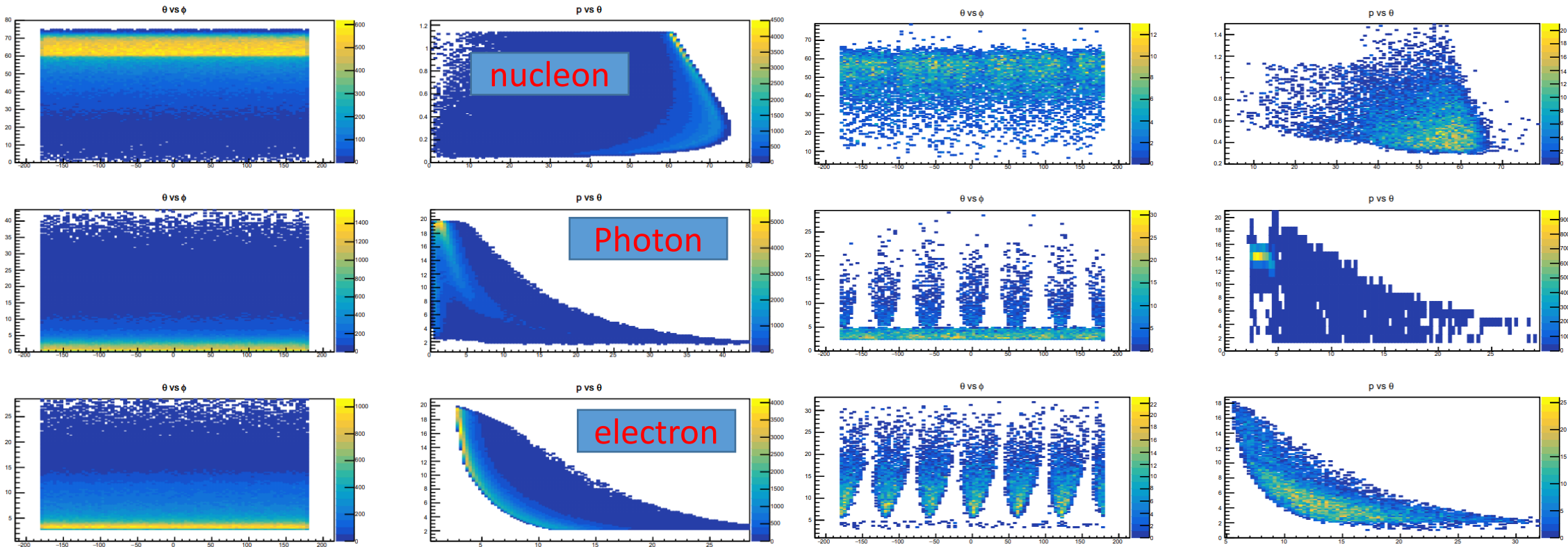


- Simulated pDVCS events at 22 GeV (H. Avakian)
- GEMC of current CLAS12
- Standard CLAS12 reconstruction

1% acceptance

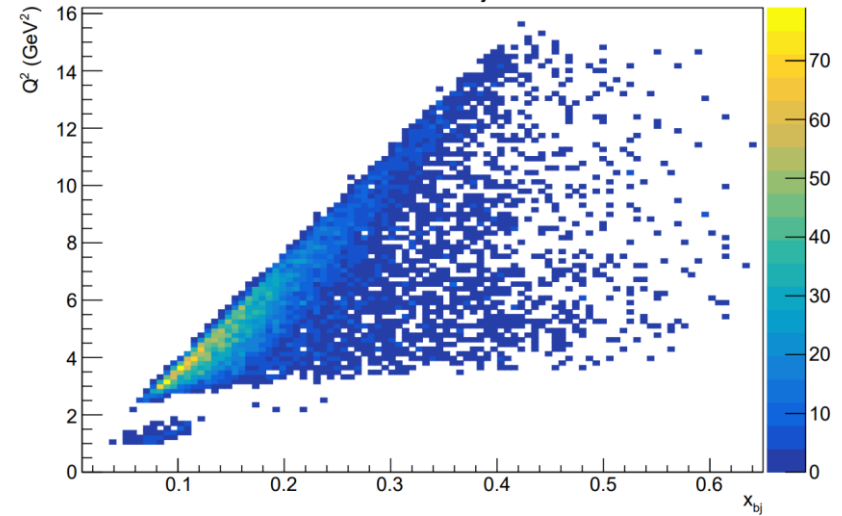
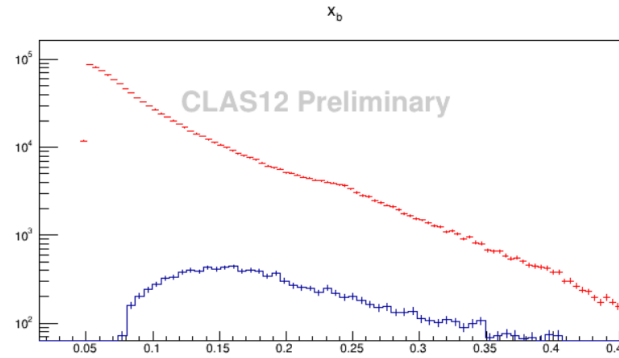
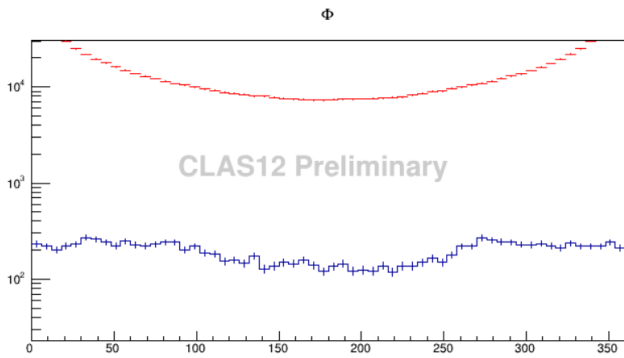
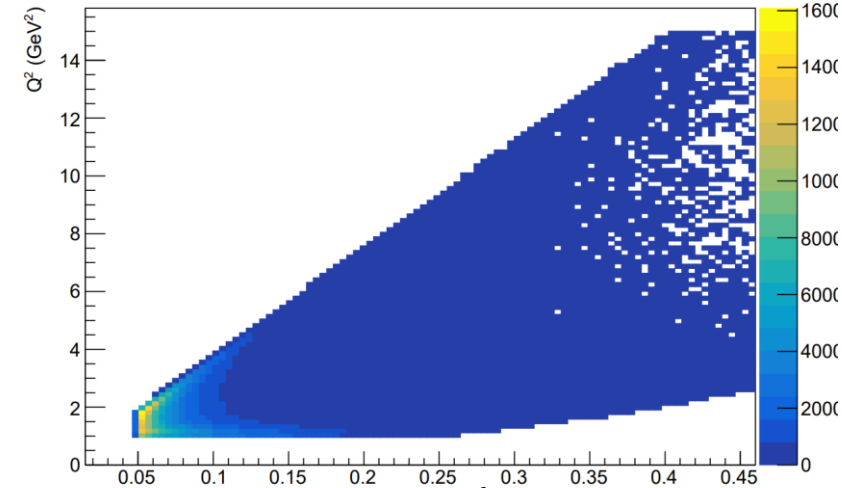
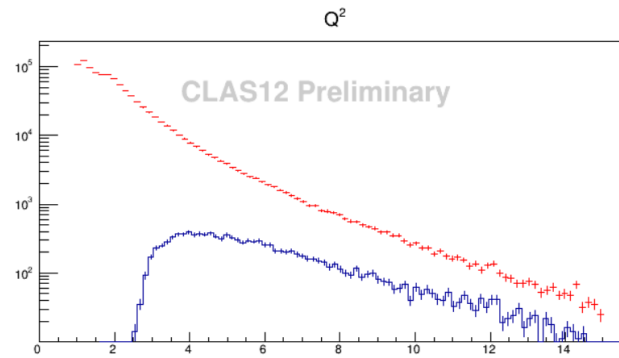
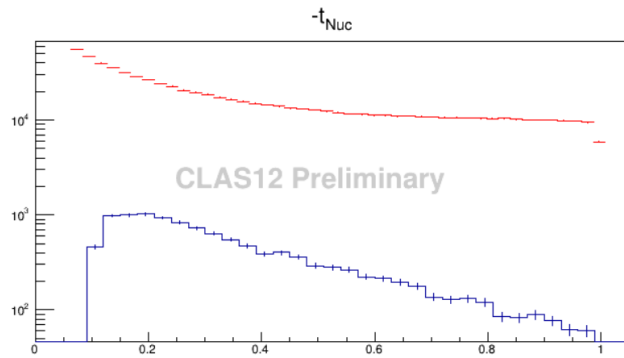
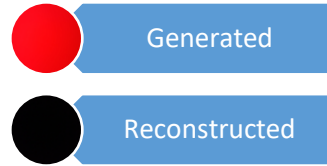
Generated

Reconstructed



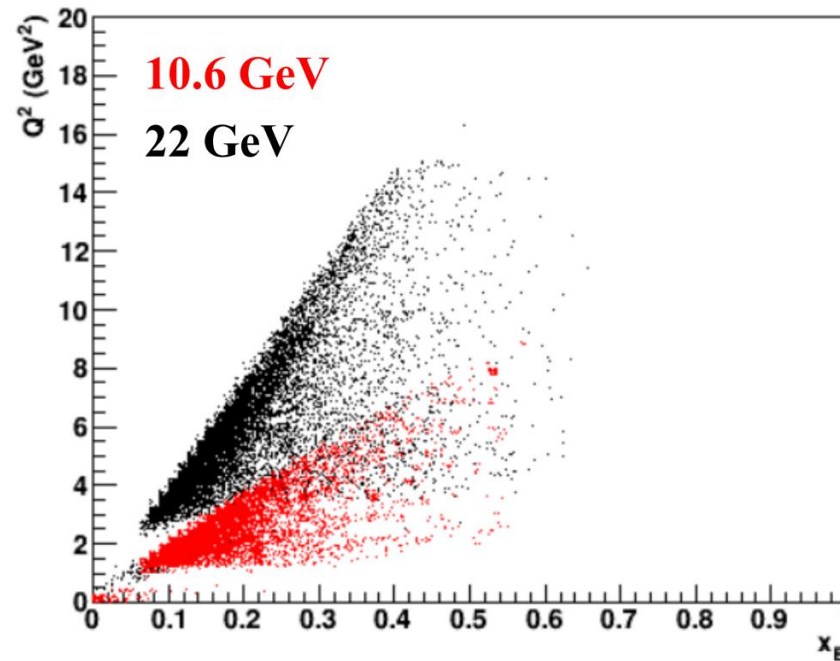


- Simulated pDVCS events at 22 GeV (H. Avakian)
- GEMC of current CLAS12
- Standard CLAS12 reconstruction



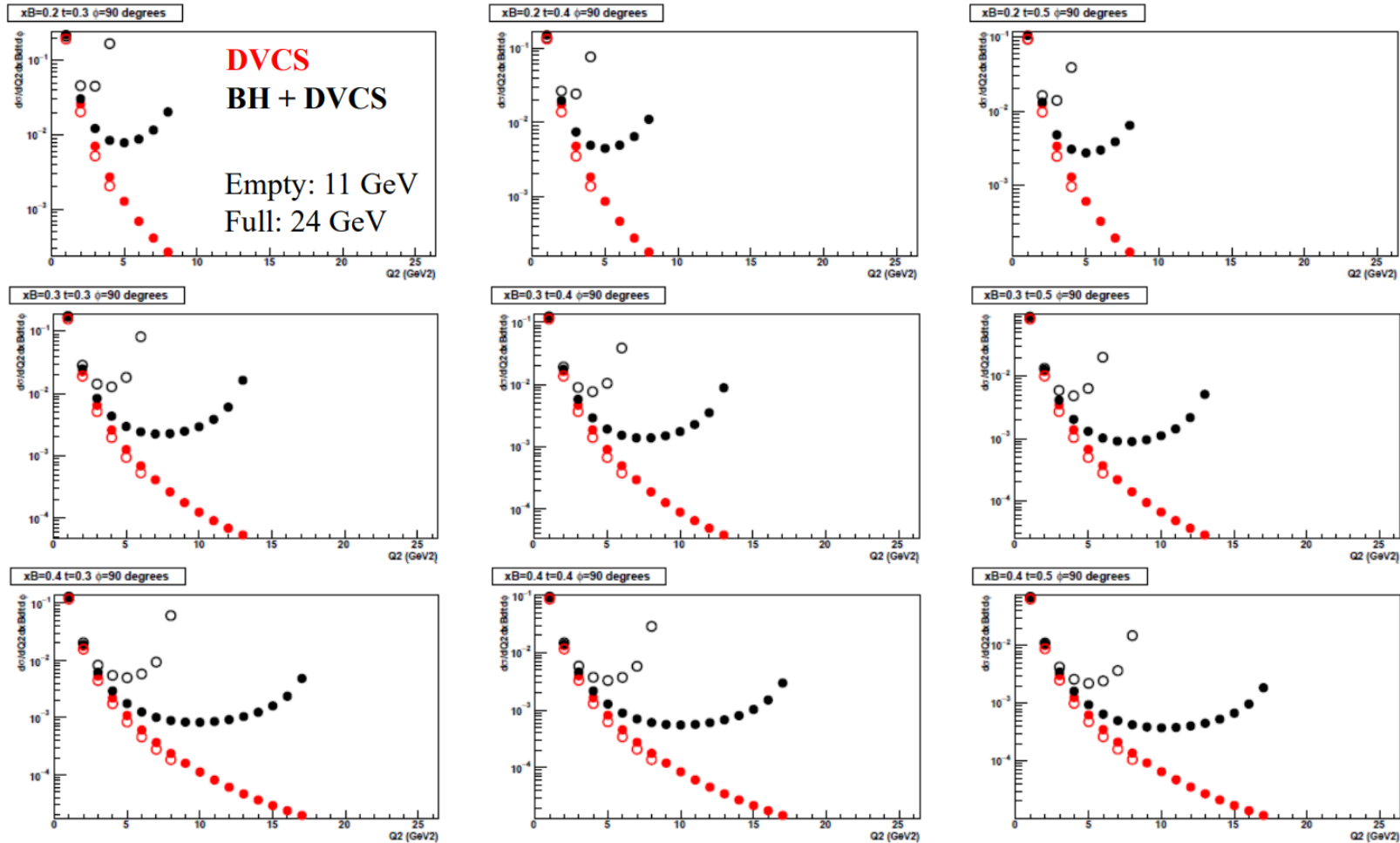


- Simulated pDVCS events at 22 GeV (H. Avakian)
- GEMC of current CLAS12
- Standard CLAS12 reconstruction





Process cross section: predictions from VGG



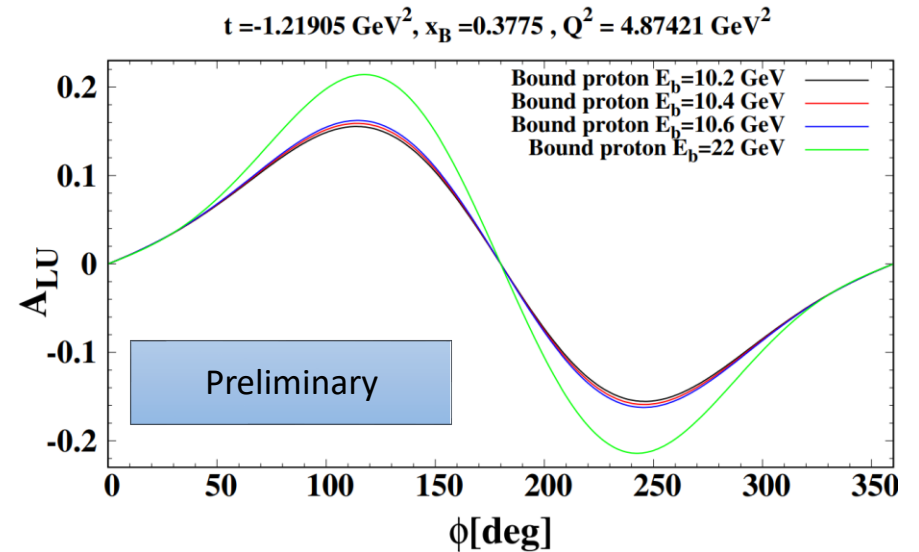
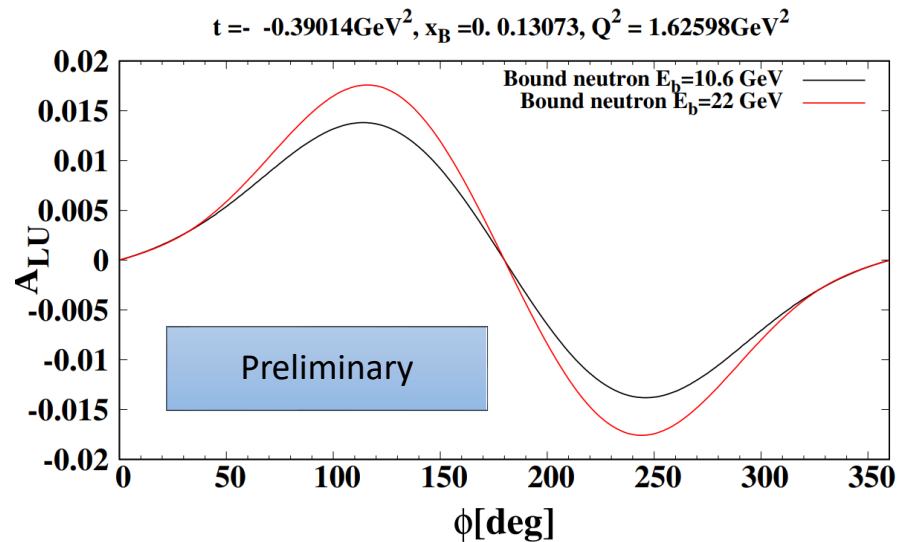


BSA at high energies

- BH dominance at higher energies
- However, BSA seems to be larger
 - pDVCS: Achieve good measurements with reasonable statistics
 - nDVCS: higher luminosity favorable given the low BSA

S. Fucini, R. Dupré, S. Scopetta to be published soon

Model predicting small nuclear effects as expected for deuteron



Larger effects are predicted for ^3He (S. Fucini, S. Scopetta, M. Viviani PRC 102 (2021) and PRD 101 (2020))



Conclusions

- GPDs are powerful tool to explore the structure of the nucleons and nuclei
 - Nucleon tomography, quark angular momentum, distribution of forces in the nucleon
- Exclusive reactions can provide important information on nucleon structure via the extraction of GPDs
- CLAS12 offers a wide kinematical reach over which the GPDs dependence on different kinematical variables can be scanned
 - Data to add constraints on GPDs in unexplored regions of the phase space
 - Possibilities to measure new observables using different experimental configurations
 - Flavor separation of GPDs
- The 20+ GeV upgrade offers a unique opportunity for a wide phase-space scan of GPDs
 - Combined with luminosity upgrade would transform CLAS into a high precision experiment
 - Will allow for x-dependence GPD scan using DDVCS
- And maybe Positron beam?