

Fixed targets at LHC

Pasquale Di Nezza



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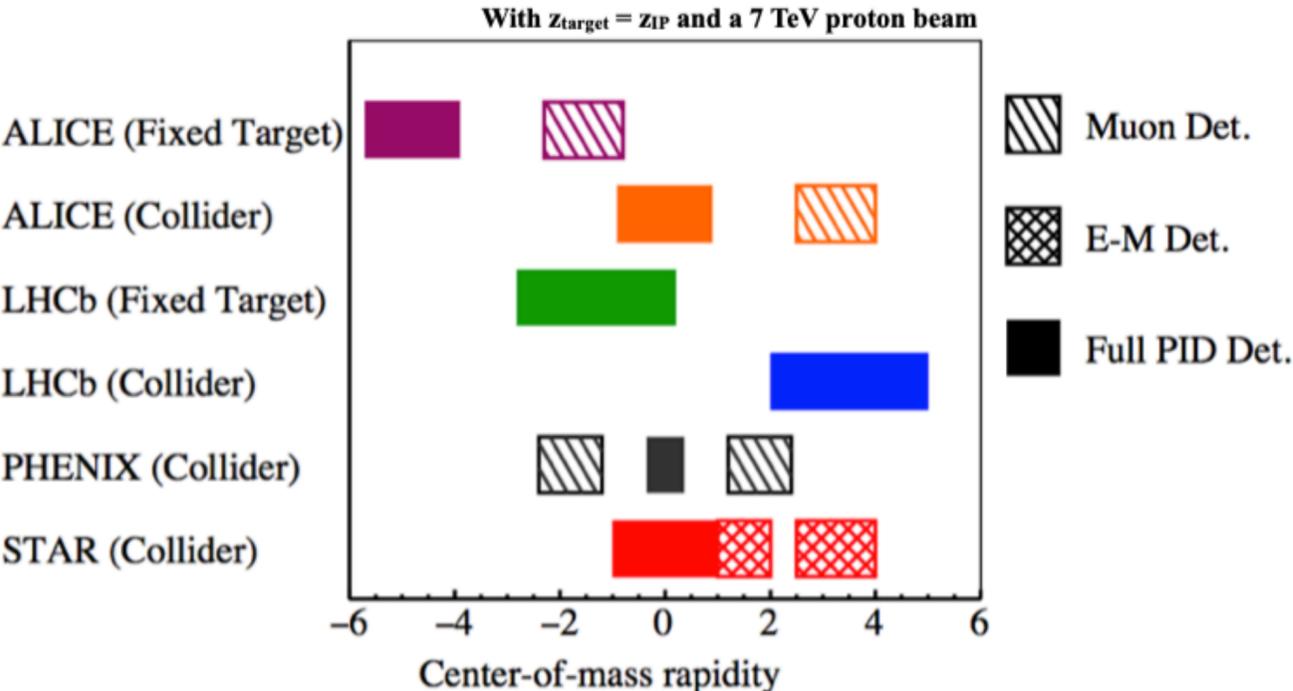


ALICE

Solid (unpolarised) target

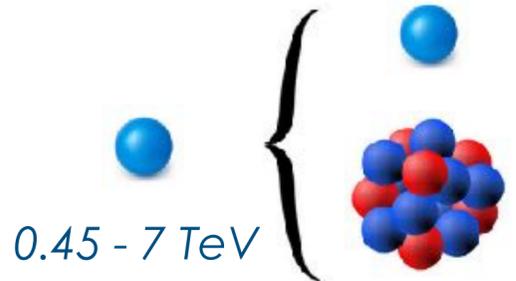


Acceptance in center-of-mass rapidity



gaseous targets @

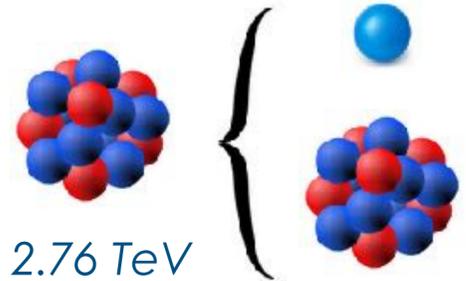




pp or pA collisions: 0.45 - 7 TeV beam on fix target

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 41 - 115 \text{ GeV}$$

$$y_{CMS} = 0 \rightarrow y_{lab} = 4.8$$



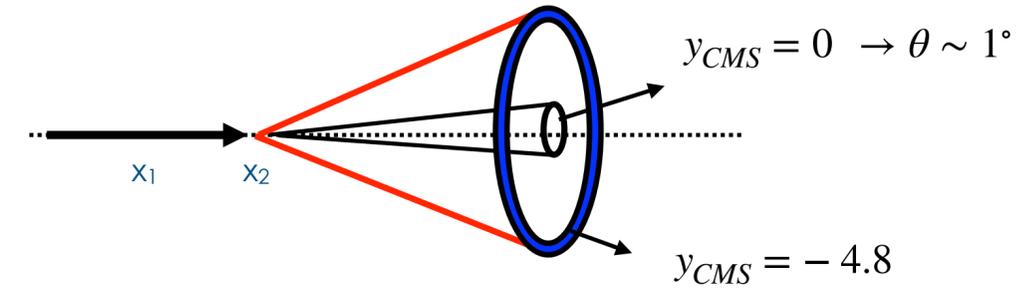
AA collisions: 2.76 TeV beam on fix target

$$\sqrt{s_{NN}} \simeq 72 \text{ GeV}$$

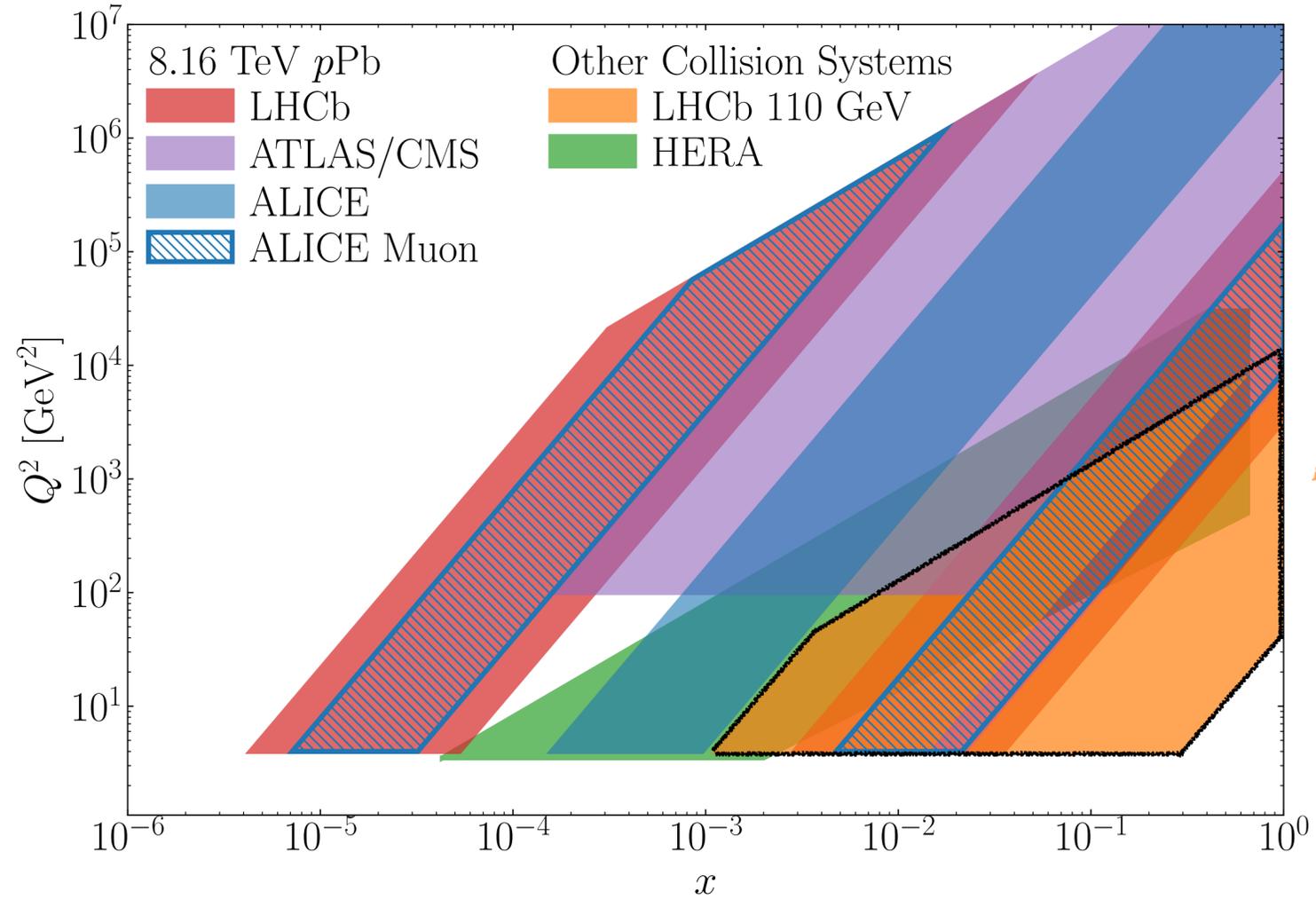
$$y_{CMS} = 0 \rightarrow y_{lab} = 4.3$$

1: beam; 2: target

Large CM boost, large x_2 values ($x_F < 0$) and small x_1



$$\gamma = \frac{\sqrt{s_{NN}}}{2m_p} \simeq 60$$



Broad and poorly explored kinematic range

SMOG2

an unpolarised target at



JINST 3 (2008) S08005
IJMPA 30 (2015) 1530022

Luminosity

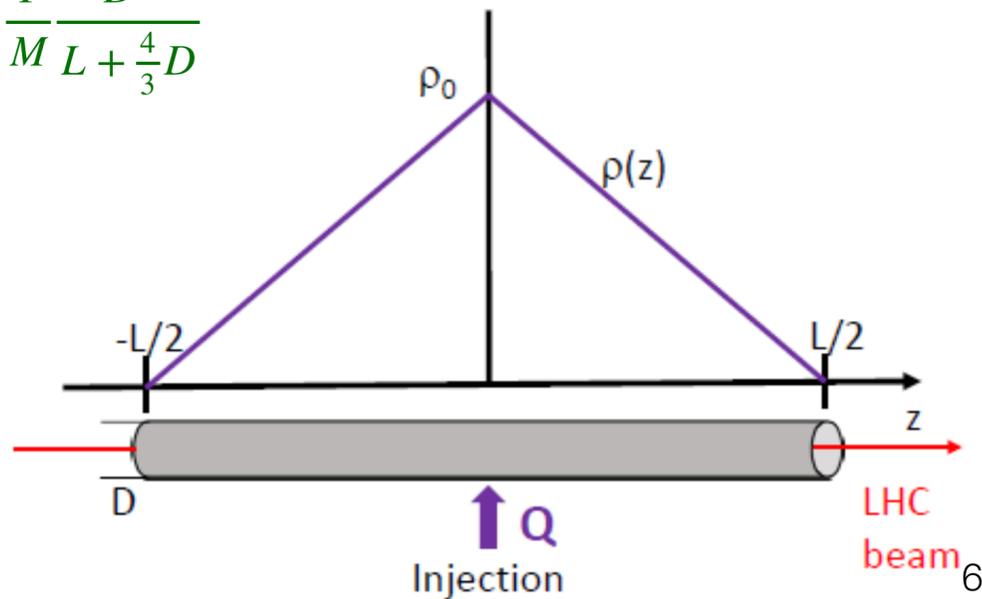
$$L_{ist} = \theta N_p f_{rev}$$

areal density
 $N_{plb} \cdot N_b$ (number of particles)

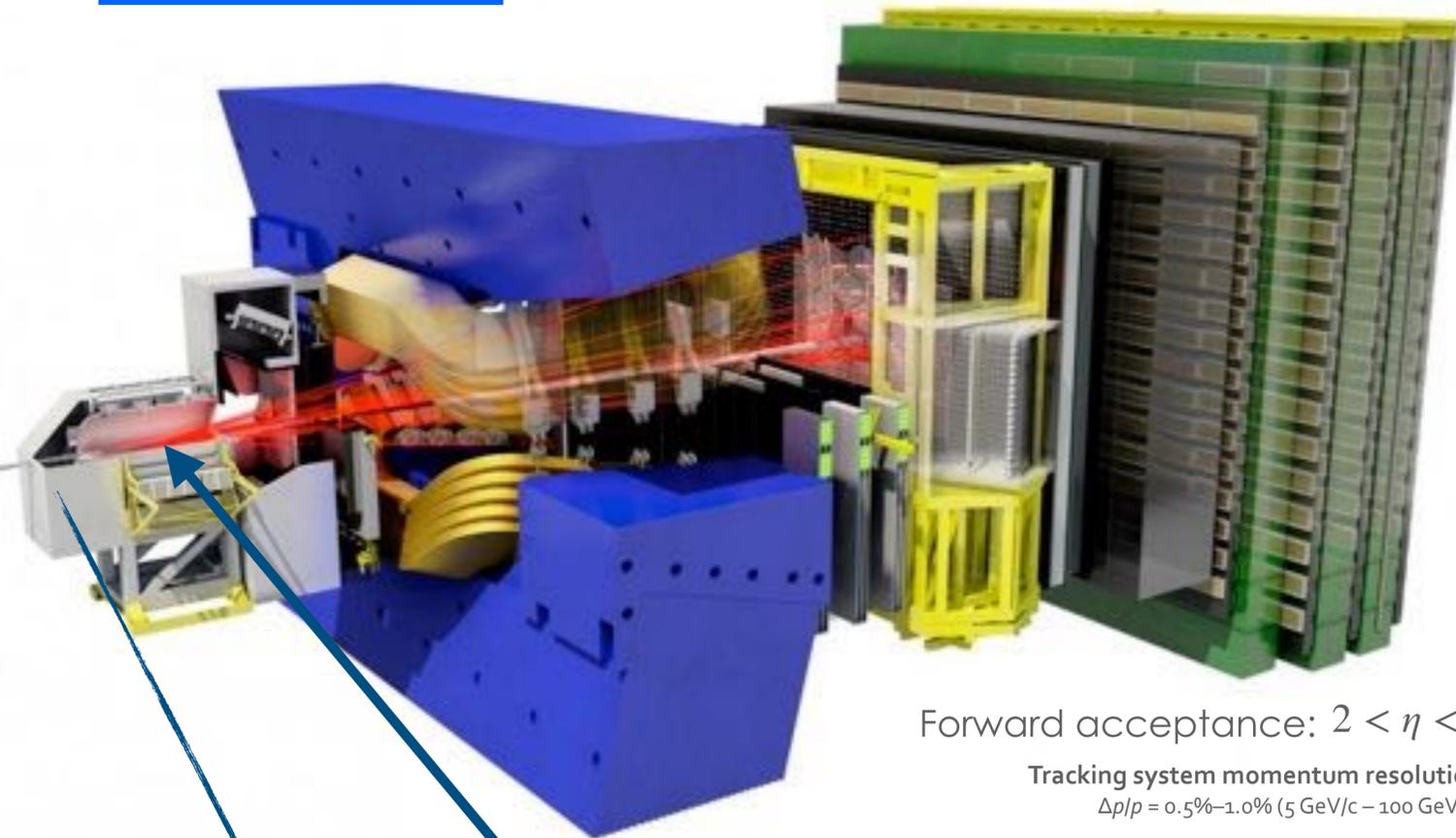
$$\rho_0 \frac{L}{2} = \frac{\Phi L}{C 2}$$

$$C = 3.81 \sqrt{\frac{T}{M}} \frac{D^3}{L + \frac{4}{3}D}$$

Storage cell concept



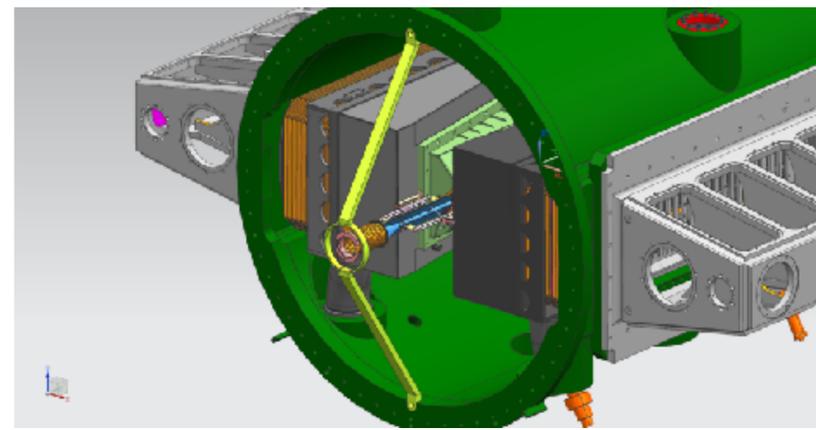
LHC beam



Forward acceptance: $2 < \eta < 5$

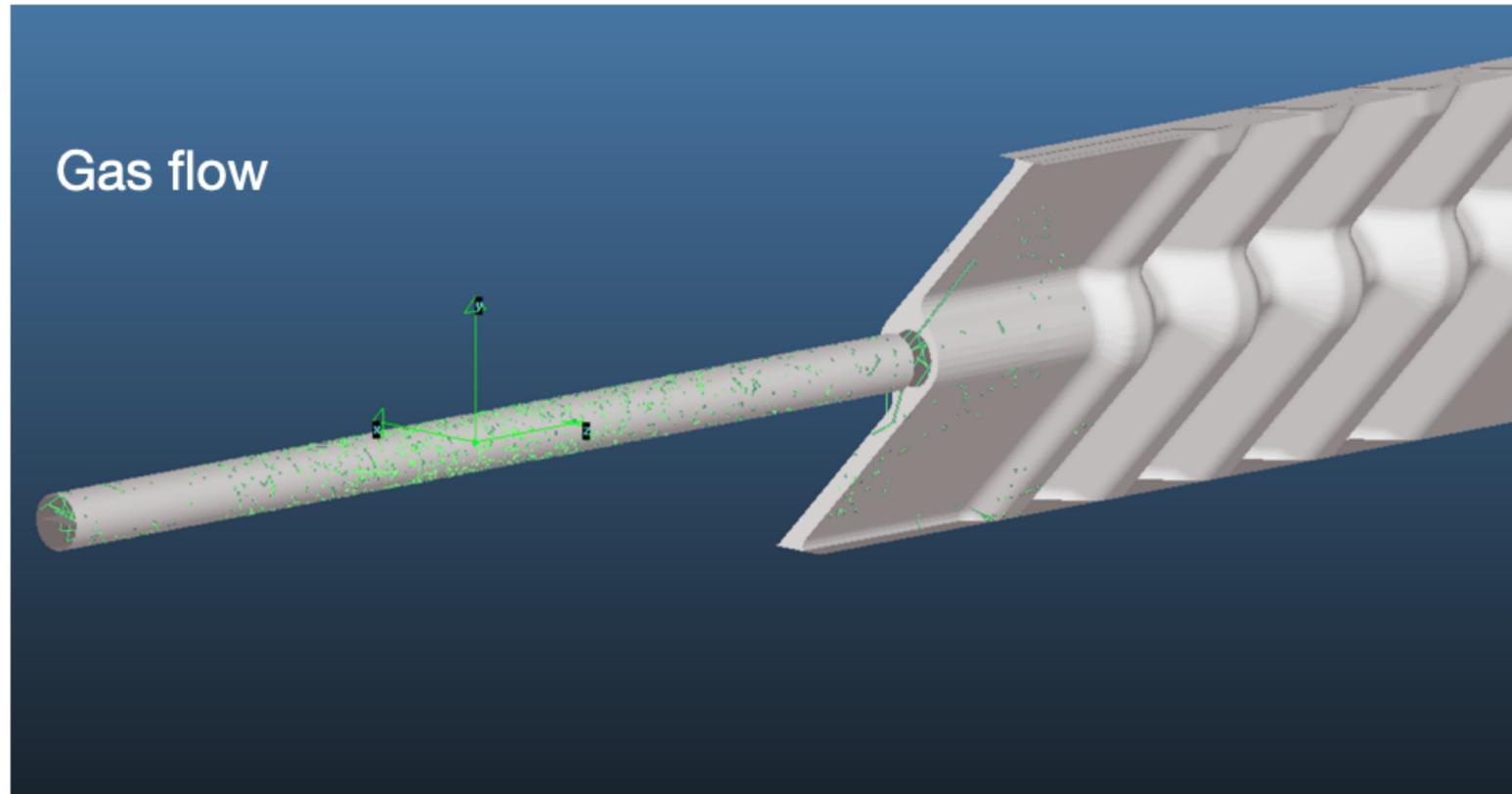
Tracking system momentum resolution
 $\Delta p/p = 0.5\% - 1.0\%$ (5 GeV/c - 100 GeV/c)

beam-beam collisions



UNpolarised target (beam-gas)

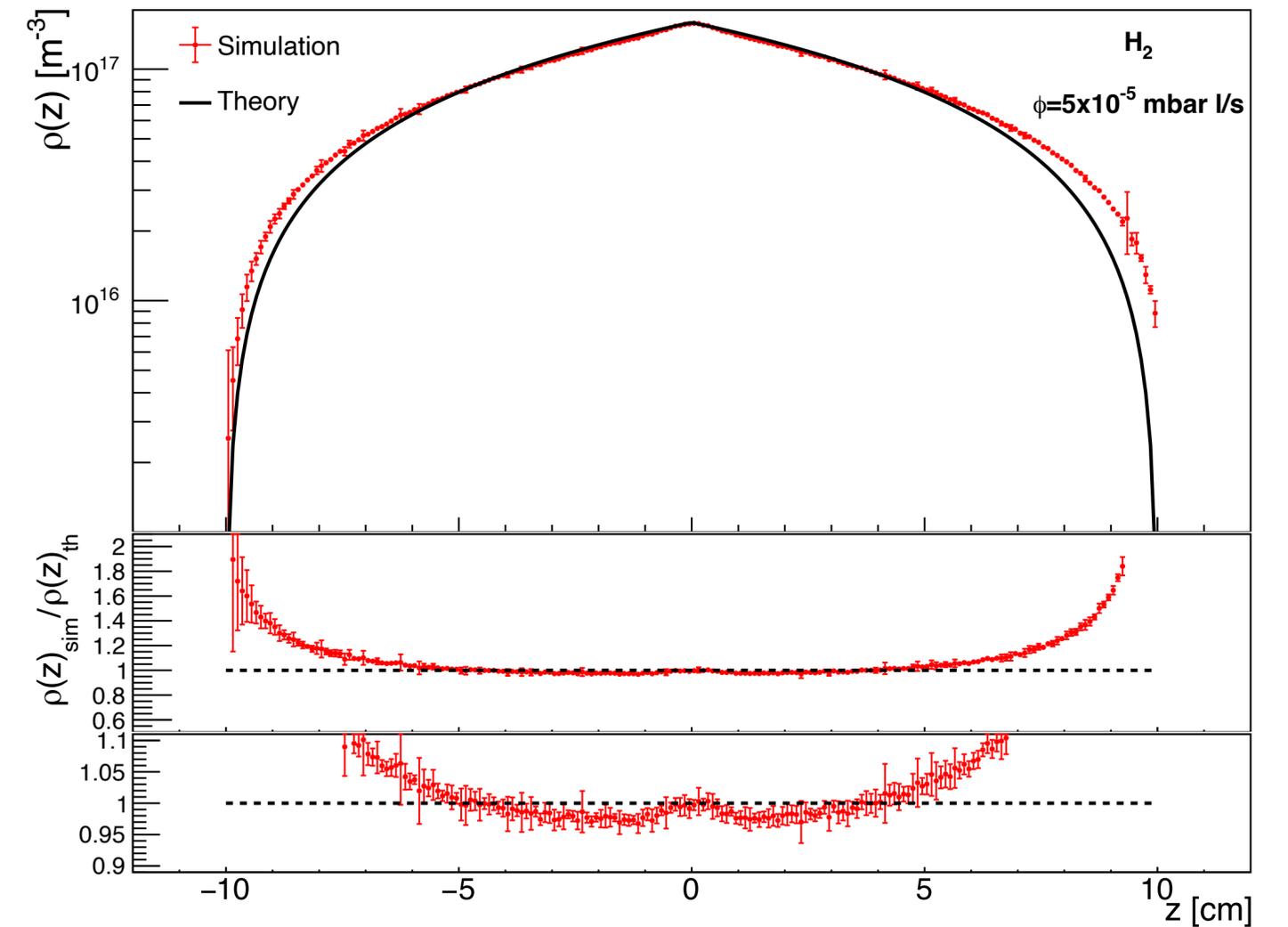
The storage cell advantage



SMOG2 example pAr @115 GeV in 1yr of data taking

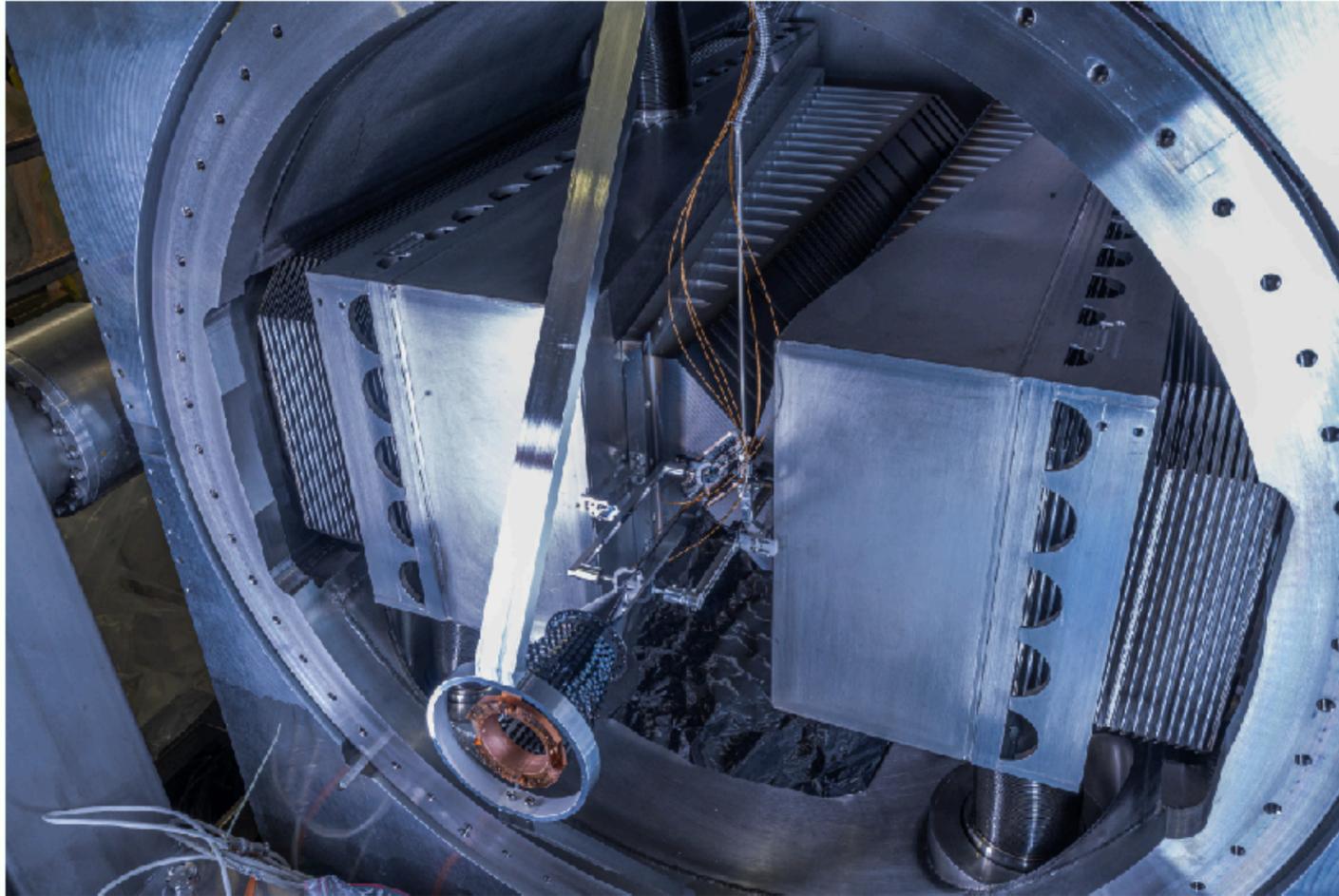
Int. Lumi.		80/pb
Sys.error of J/Ψ xsection		~3%
J/Ψ yield		28 M
D^0 yield		280 M
Λ_c yield		2.8 M
Ψ' yield		280 k
$Y(1S)$ yield		24 k
$DY \mu^+\mu^-$ yield		24 k

Very high statistics with a low gas flow

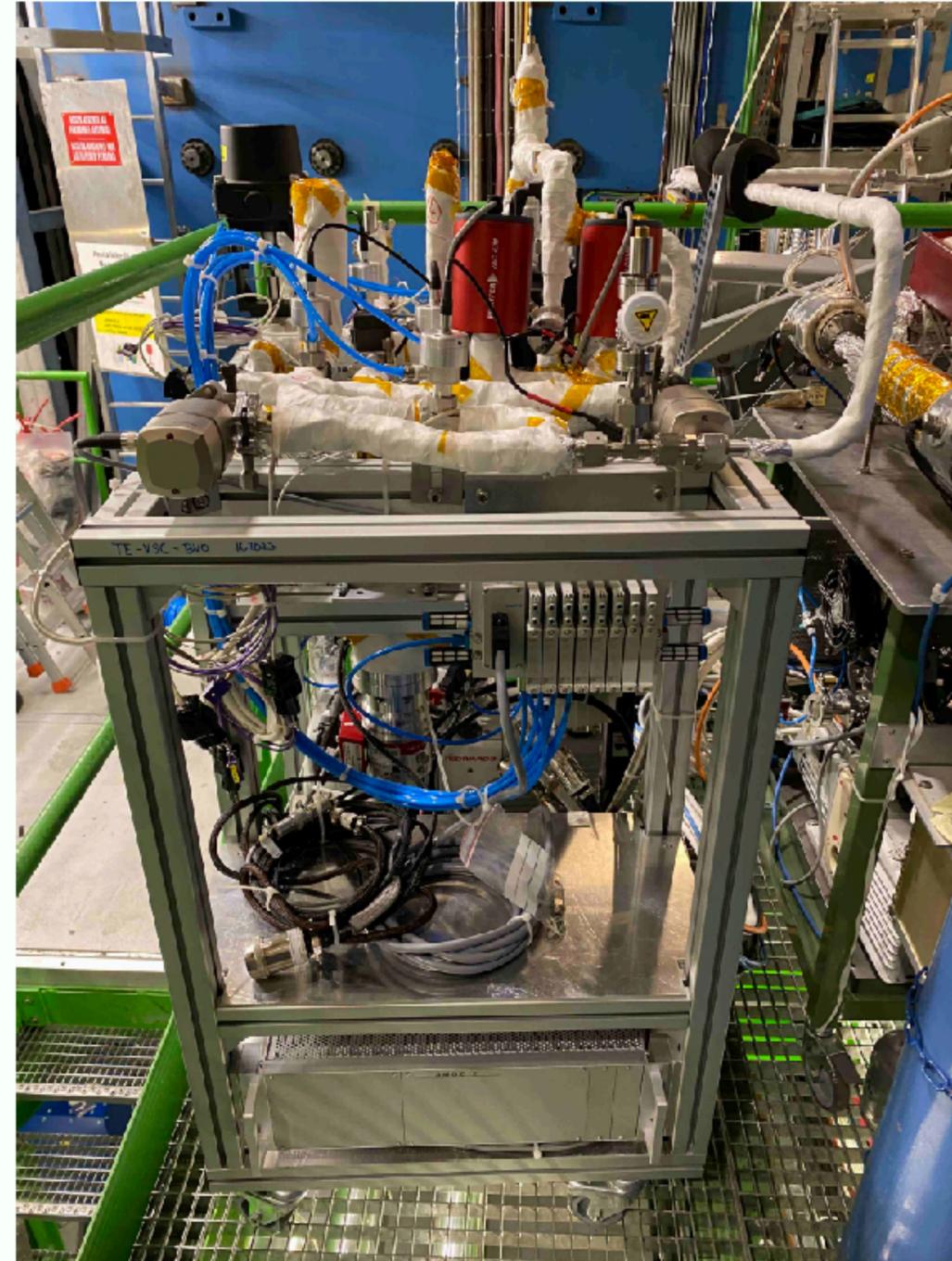


	$\theta_{sim} \times 10^{12} [cm^{-2}]$	$\theta_{th} \times 10^{12} [cm^{-2}]$	C_θ
Hydrogen	1.627	1.592	1.022
Argon	7.274	7.120	1.022

SMOG2

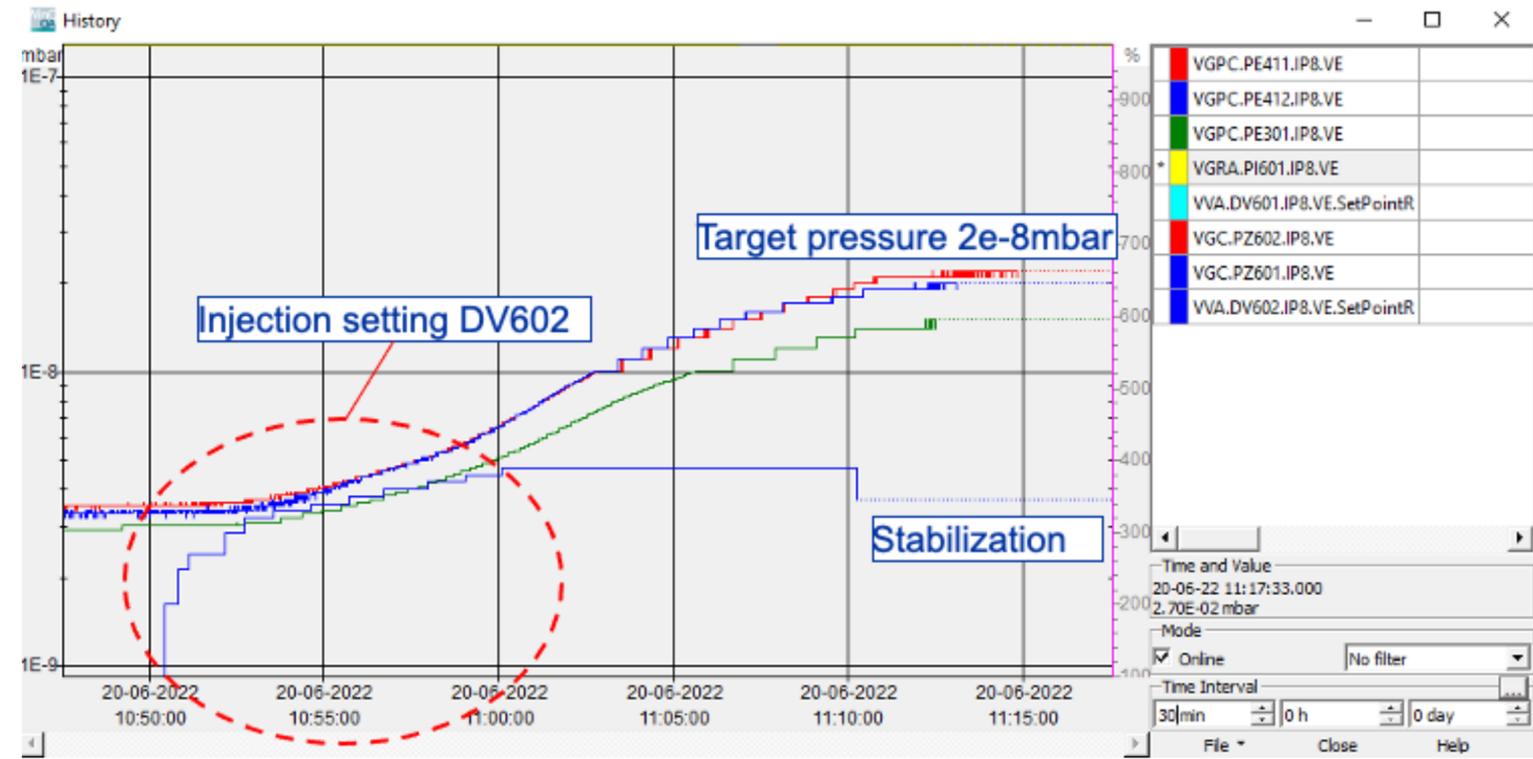


- The system is completely installed (storage cell + GFS + triggers + reconstruction)
- Negligible impact on the beam lifetime ($\tau_{beam-gas}^{p-H_2} \sim 2000$ days , $\tau_{beam-gas}^{Pb-Ar} \sim 500$ h)
- Injectable gases: He, Ne, Ar ... H₂, D₂, N₂, O₂, Kr, Xe

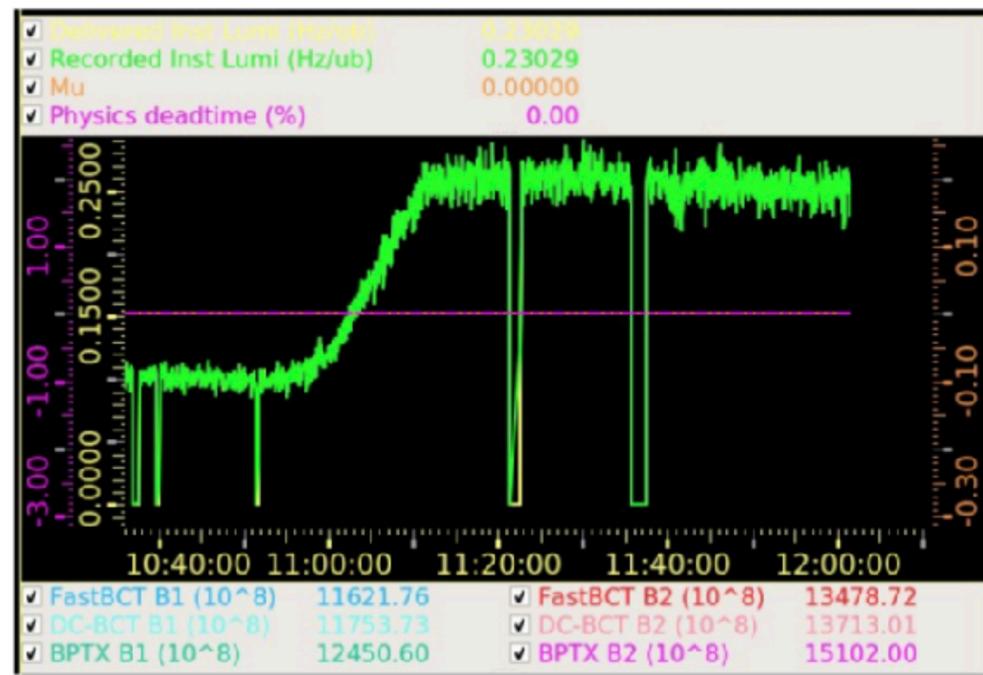
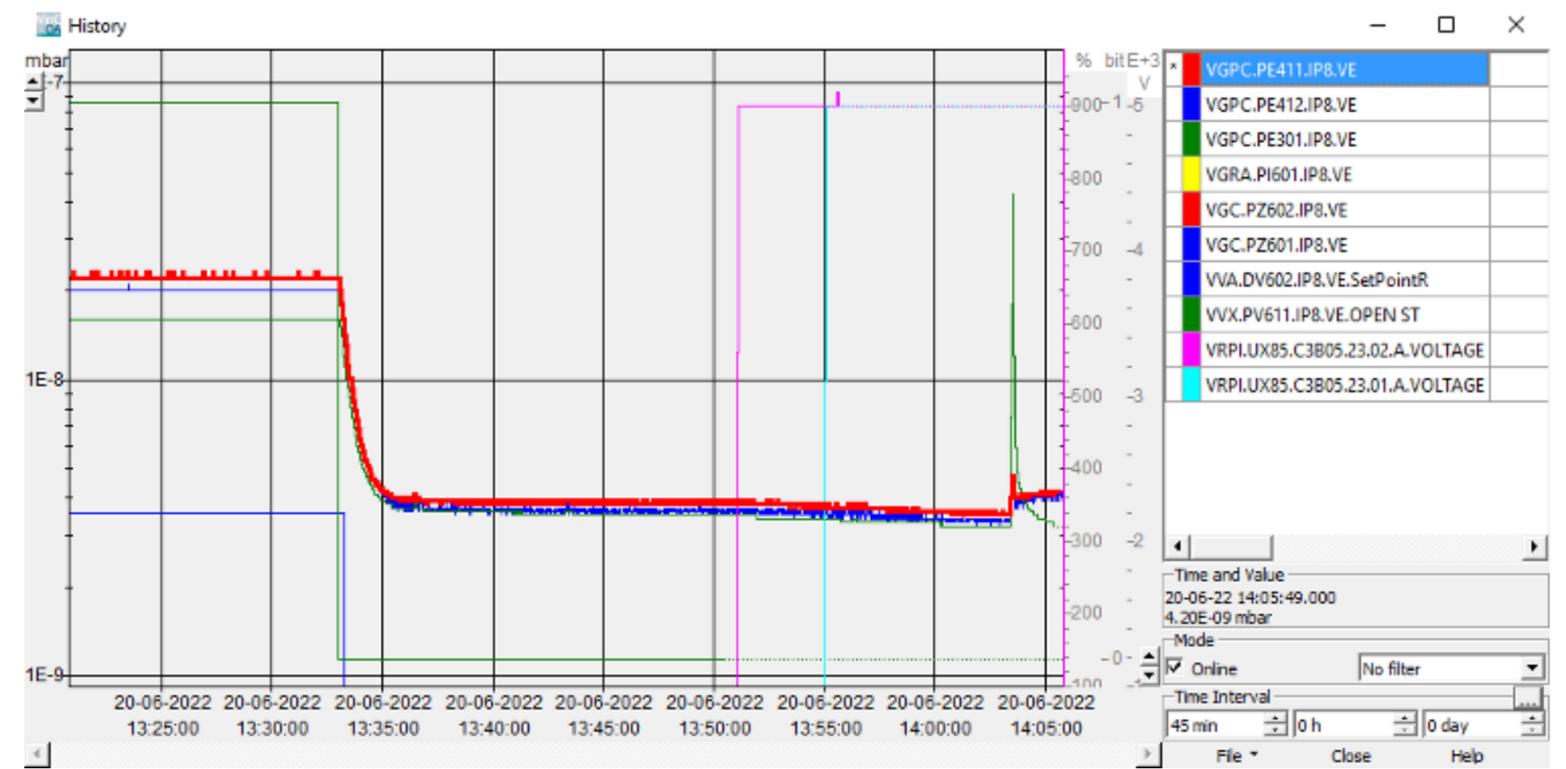


SMOG2 gas injection at LHC Run3 started a couple of weeks ago

Pressure increase into the primary vacuum



Vacuum recovery after the gas injection stop



Luminosity increase seen by the LHCb luminometer (Plume)

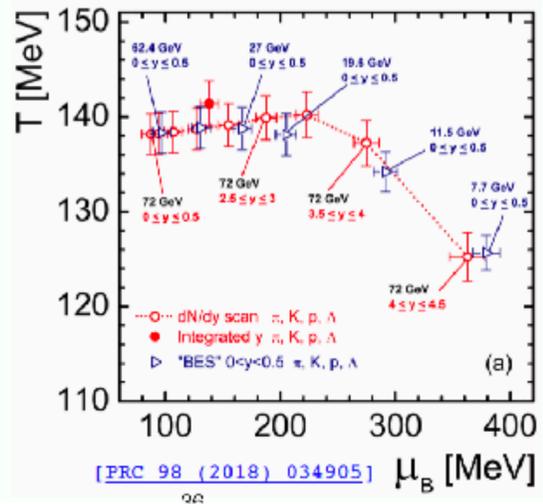
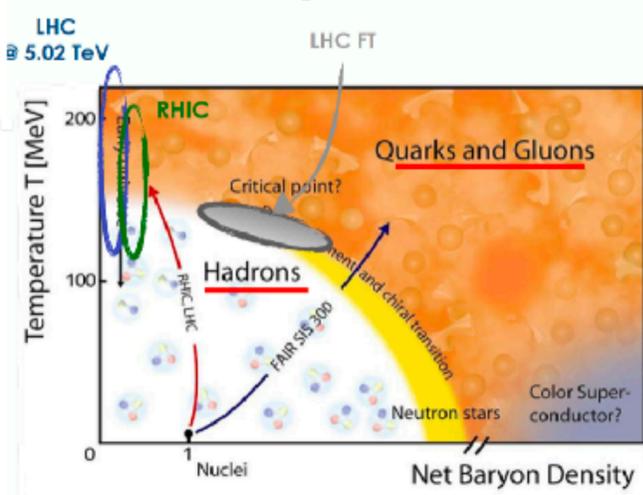
LHC official statement

No negative feedback when there is gas injection. Green light to inject when needed

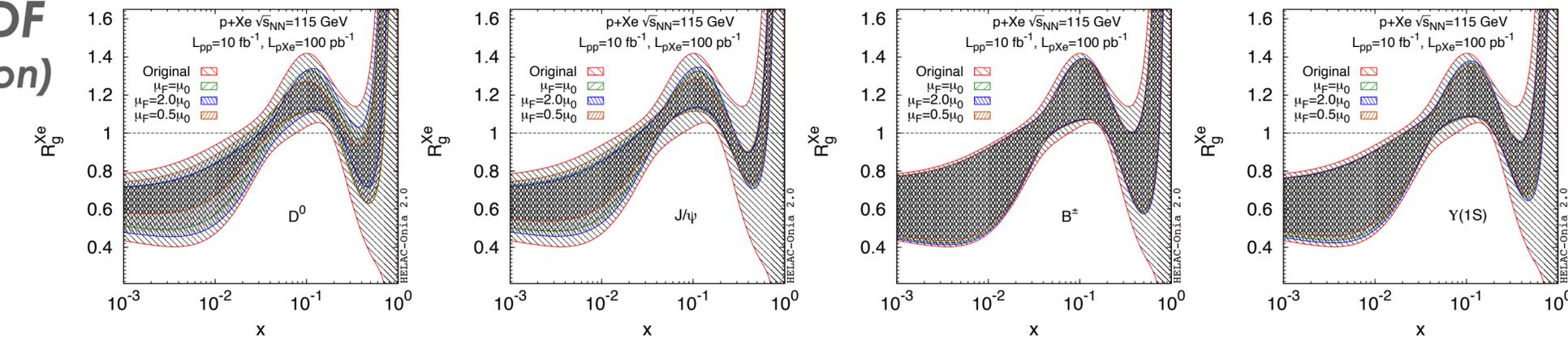
SMDQ2 ... few highlights

<http://cds.cern.ch/record/2649878/files/>

Heavy-Ion and QCD phase space

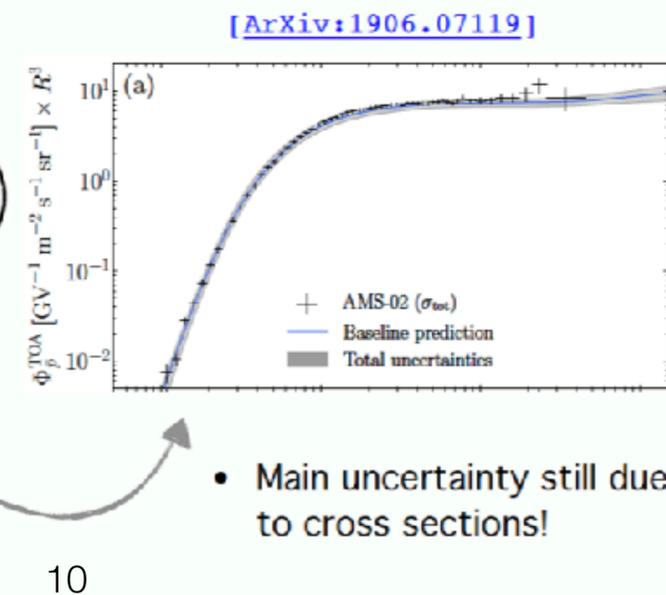
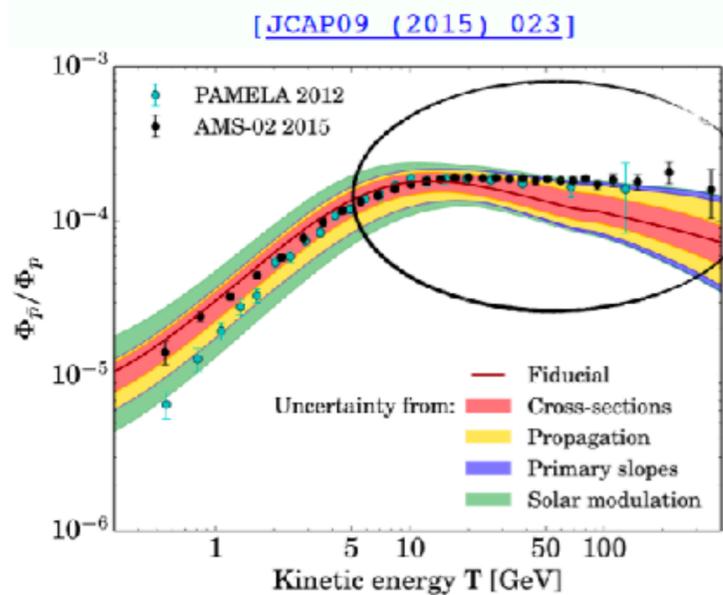
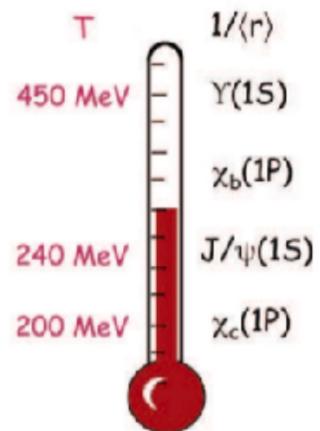


nPDF (gluon)

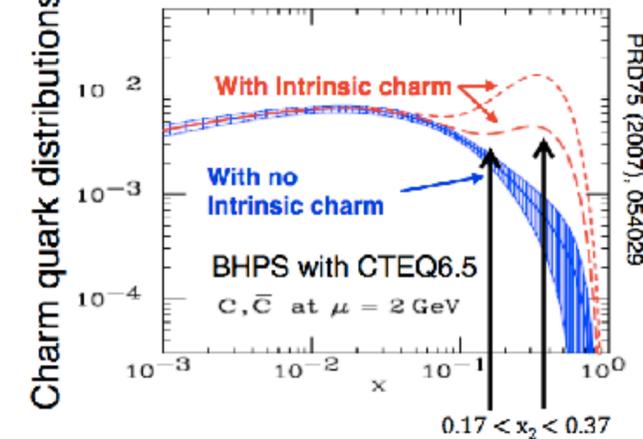


Astroparticle (DM and CR)

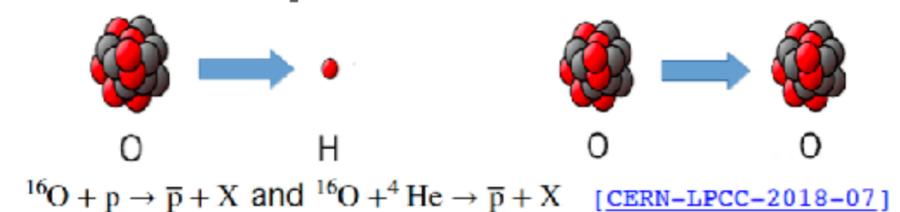
c \bar{c} bound states



Intrinsic charm

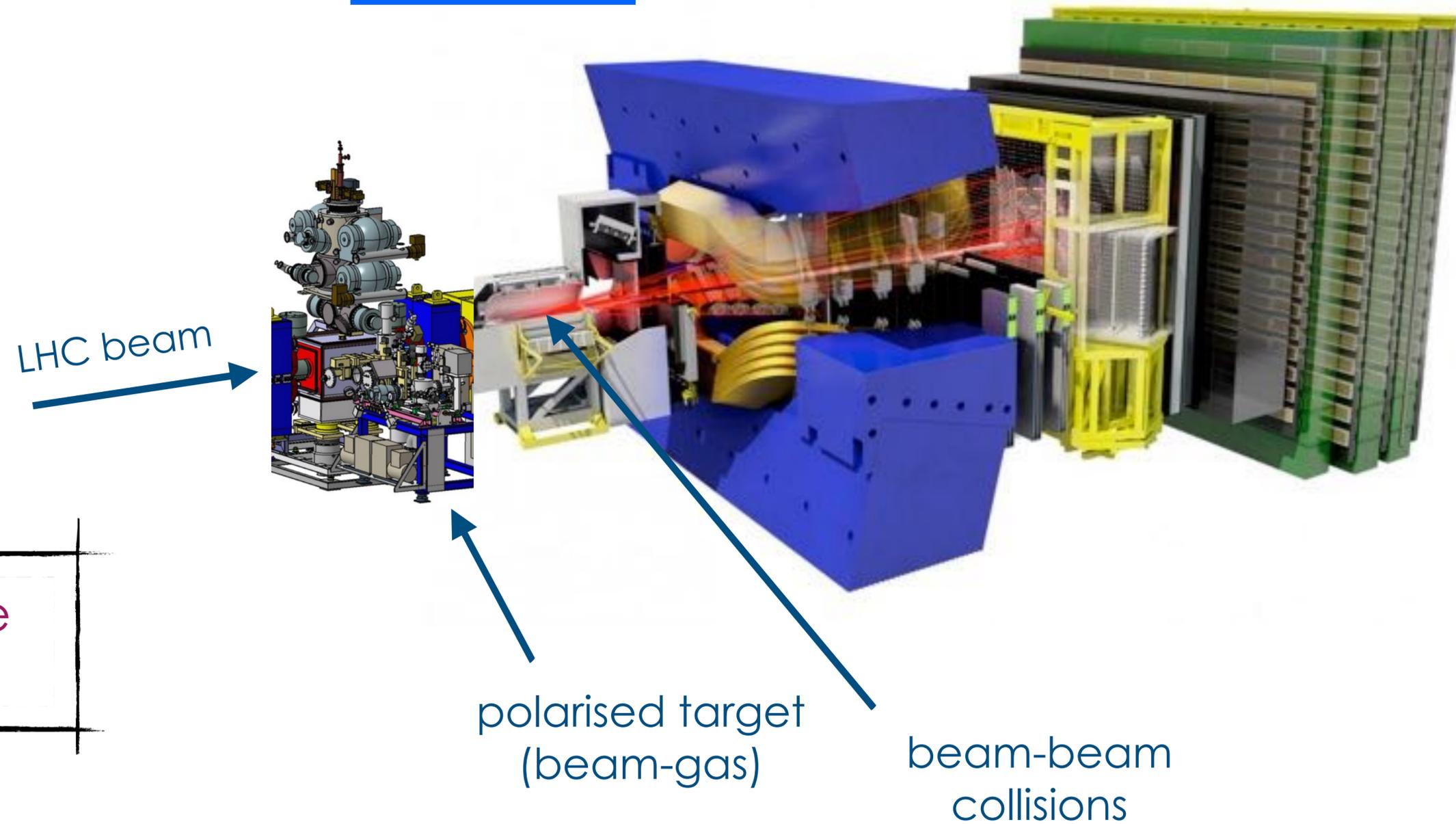


Special Runs





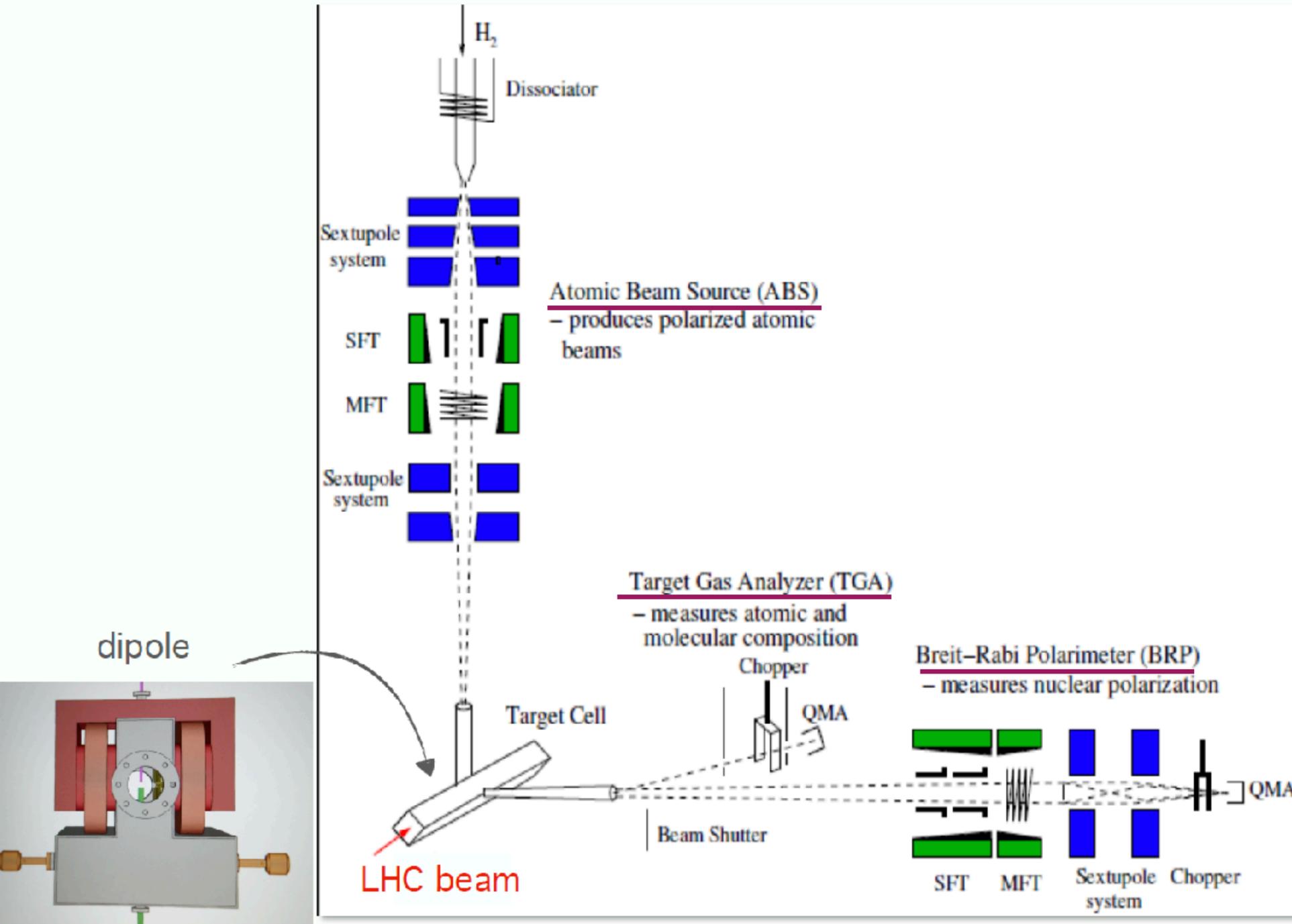
a polarised target at



The LHC beams cannot be polarised

SMOG2 is not only a unique project itself, but also a great playground for LHC spin

LHCspin experimental setup

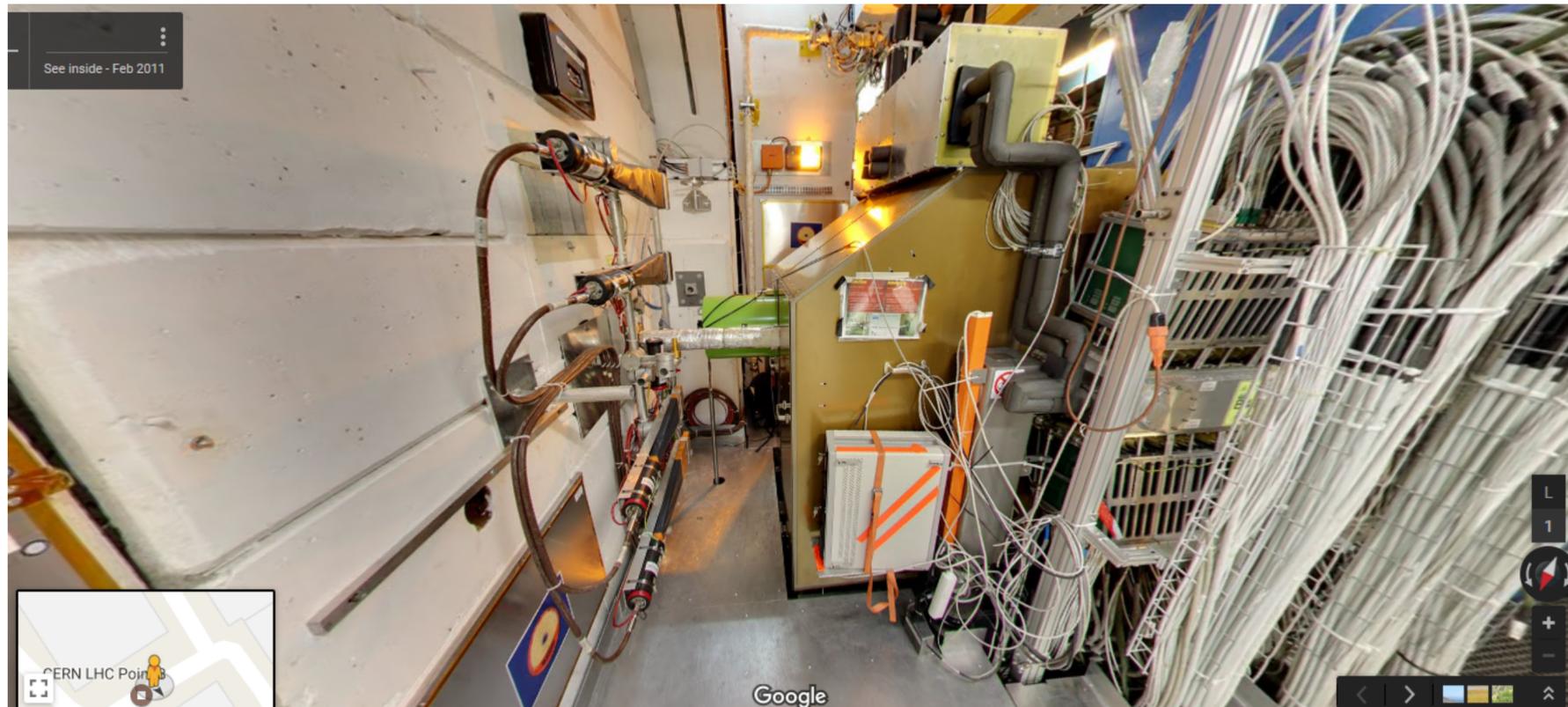
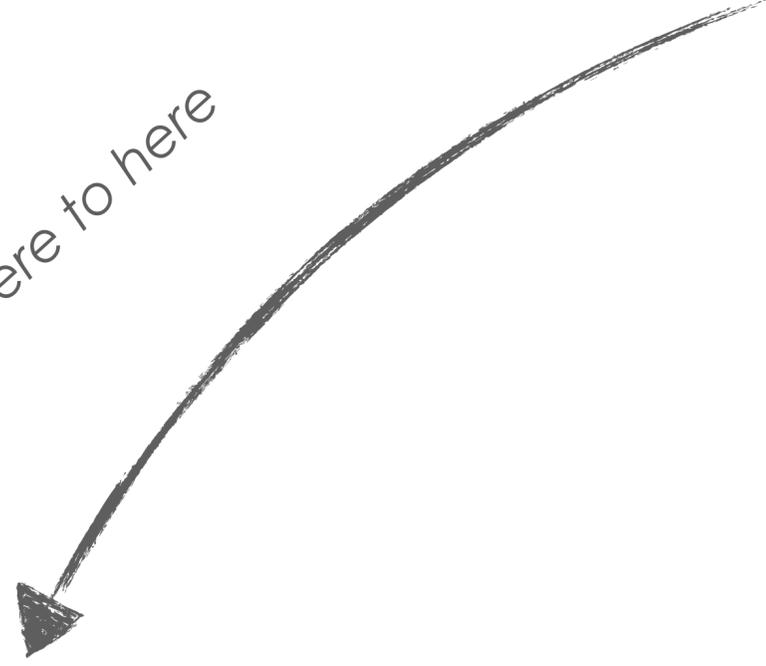


- Start from the well established HERMES setup @ DESY...
- ... to create the next generation of fixed target polarisation techniques!



HERMES PGT

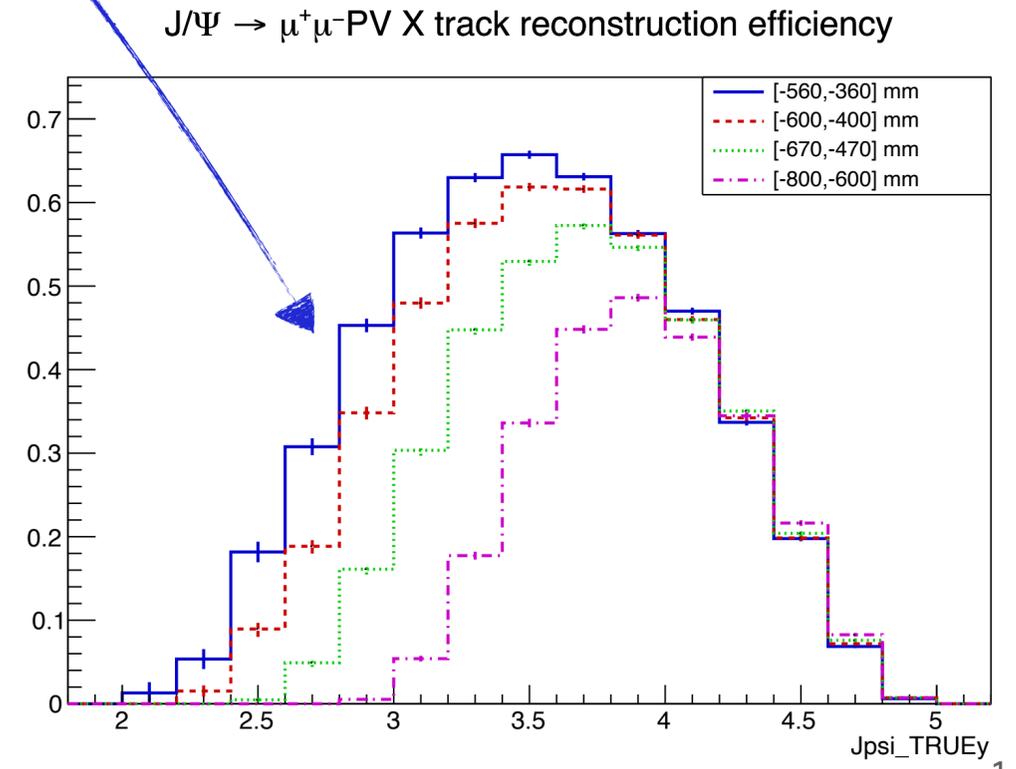
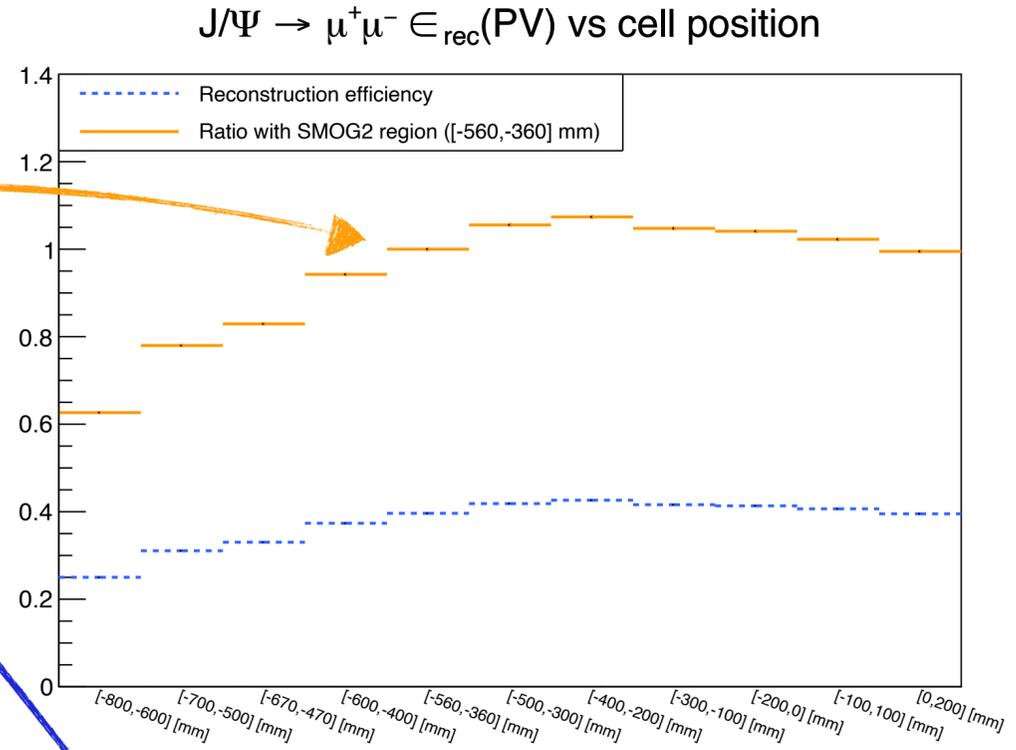
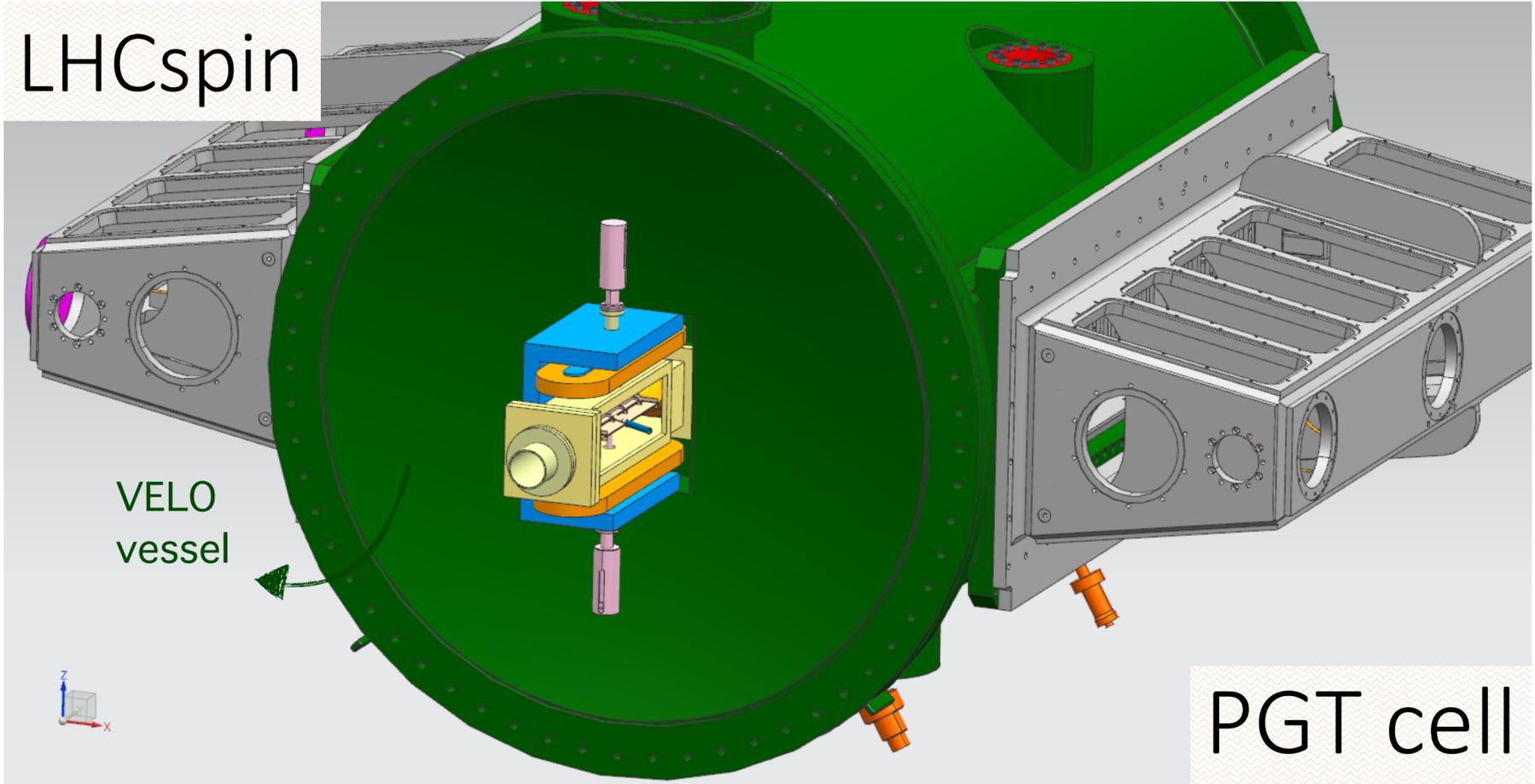
From here to here



Space available in front of LHCb

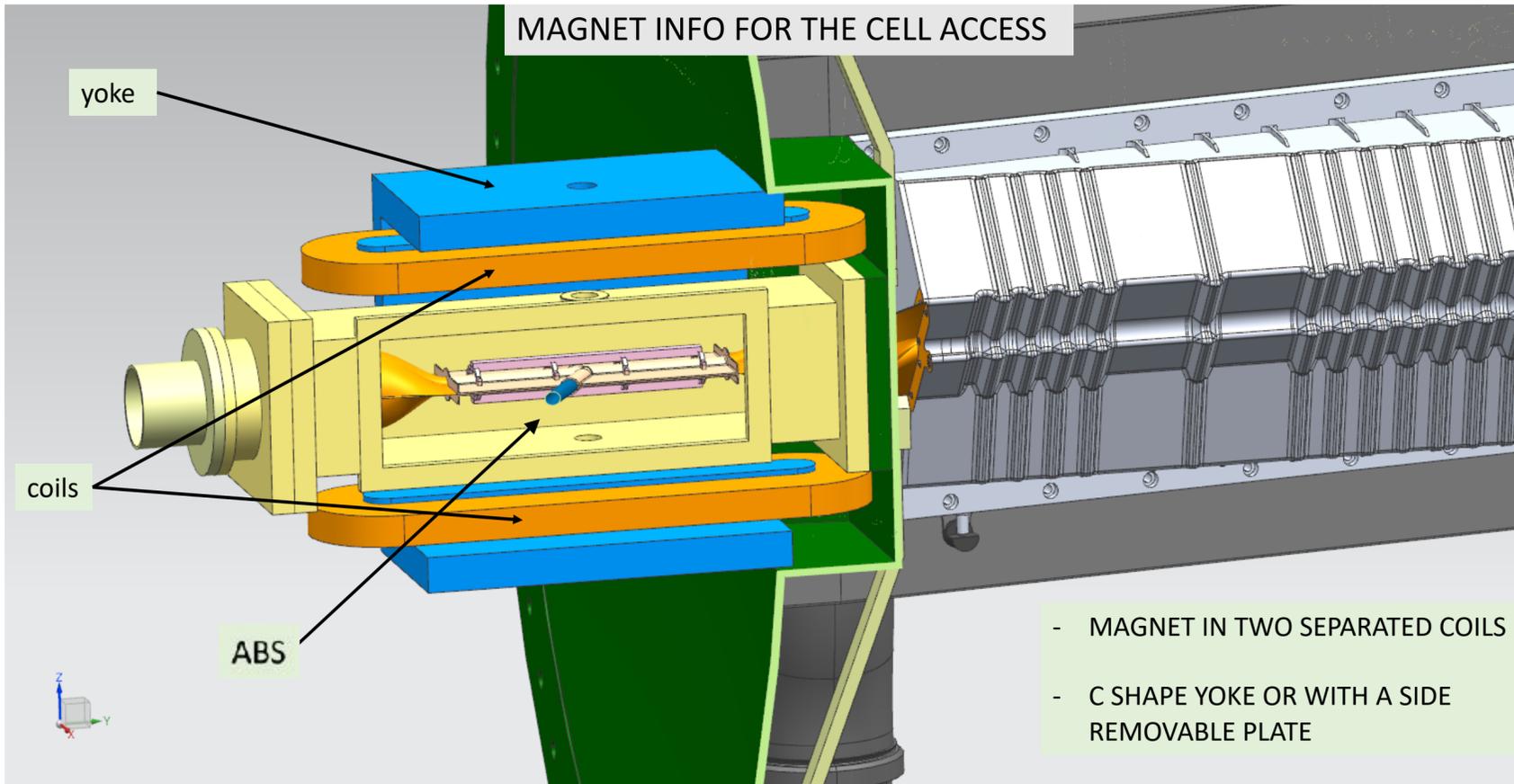
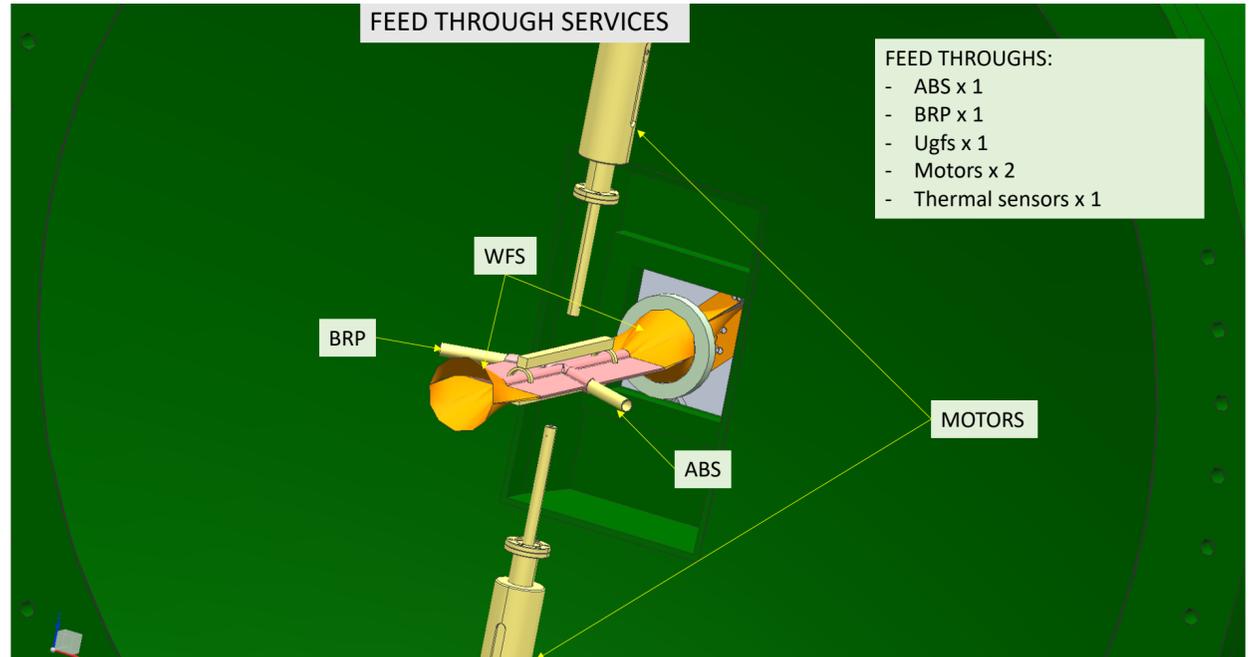
PGT implementation into LHCb

- Cylindrical target cell with SMOG2 dimensions: $L = 20$ cm and $D = 1$ cm
- Full LHCb simulations show broader kinematic acceptance & higher efficiency in the same position of the SMOG2 cell
- Work ongoing to develop dedicated trigger lines and to improve reconstruction algorithms for Run 3



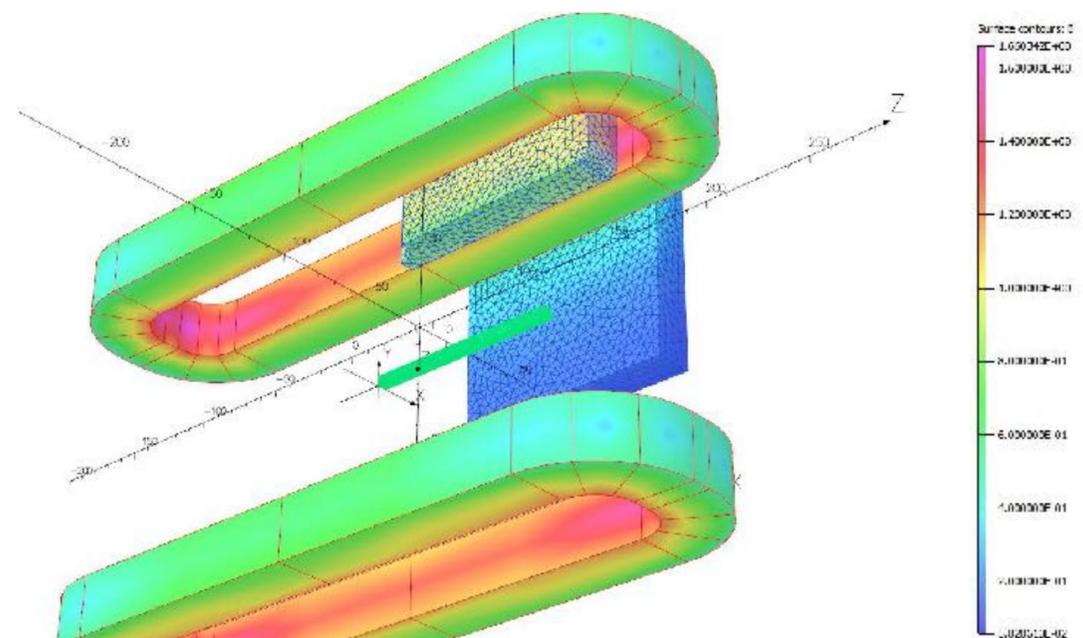
PGT implementation into LHCb

- Inject both polarised and unpolarised gases via ABS and UGFS

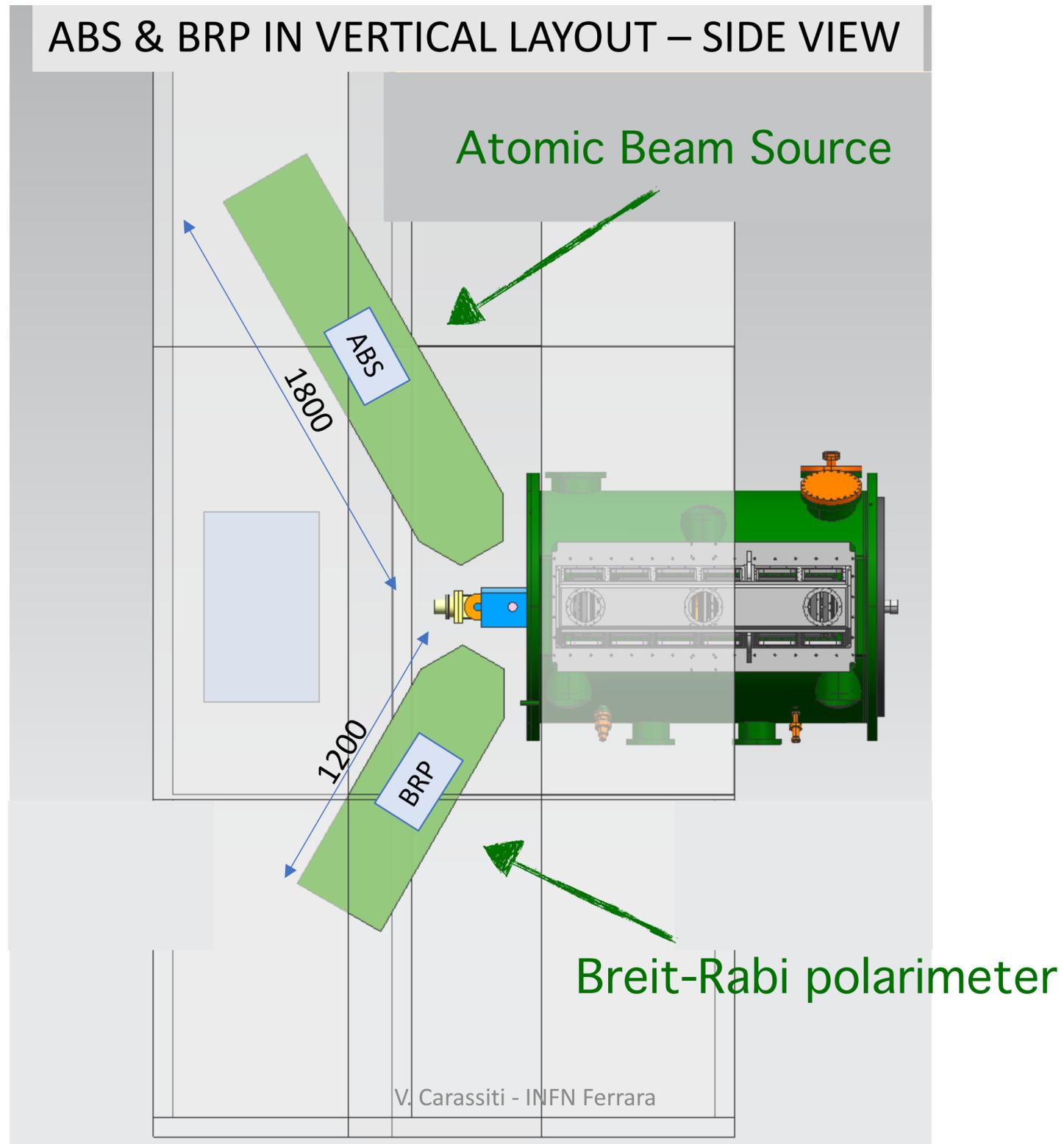


- Compact dipole magnet → static transverse field
- Superconductive coils + iron yoke configuration fits the space constraints
- $B = 300 \text{ mT}$ with polarity inversion, $\Delta B/B \approx 10\%$, suitable to avoid beam-induced depolarisation [[PoS \(SPIN2018\)](#)]

Possibility to switch to a solenoid and provide longitudinal polarisation (e.g. in Run 5)



ABS & BRP implementation into LHCb



- Reduce the size of both ABS and BRP to fit into the available space in the LHCb cavern: a challenging R&D!
- **No need for additional detectors** in LHCb: only a modification of the VELO flange is needed
- $P \simeq 85\%$ achieved at HERMES

Injected intensity of H-atoms:
 $\phi = 6.5 \times 10^{16} \text{ s}^{-1}$

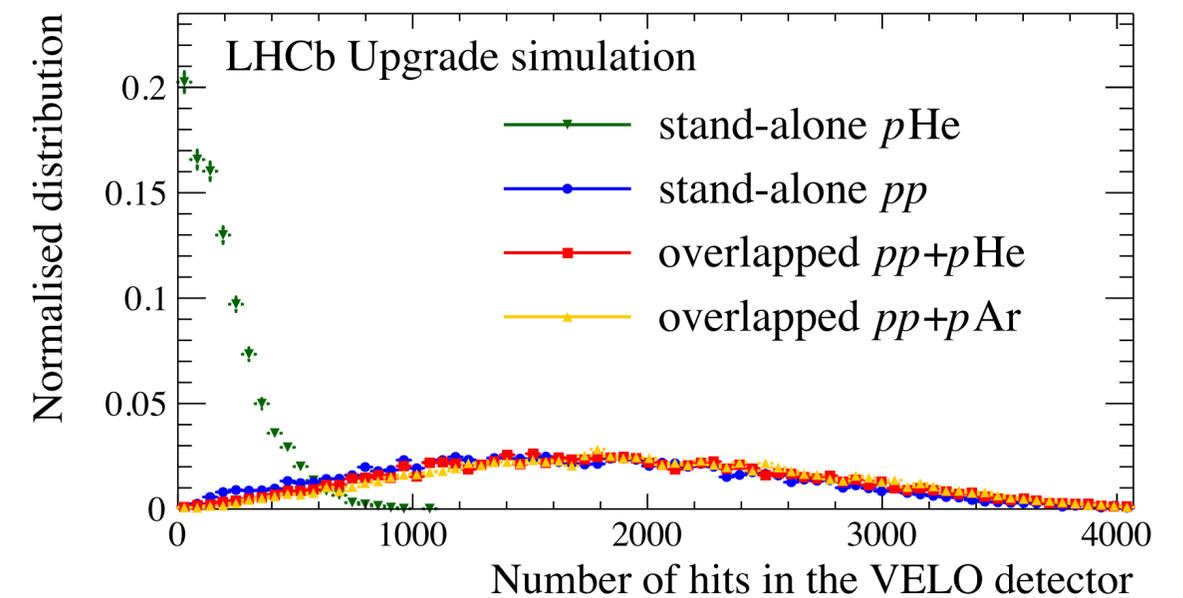
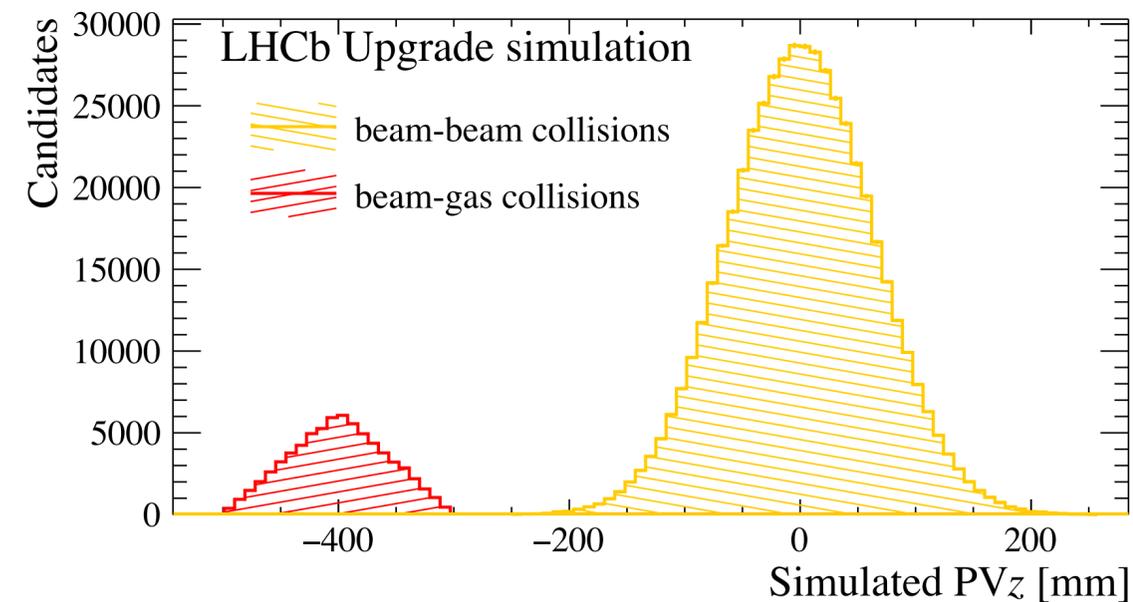
Achievable Luminosity (HL-LHC):
 $\sim 8 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$

- Backup solution is being investigated: a jet target provides lower density but higher polarisation degree

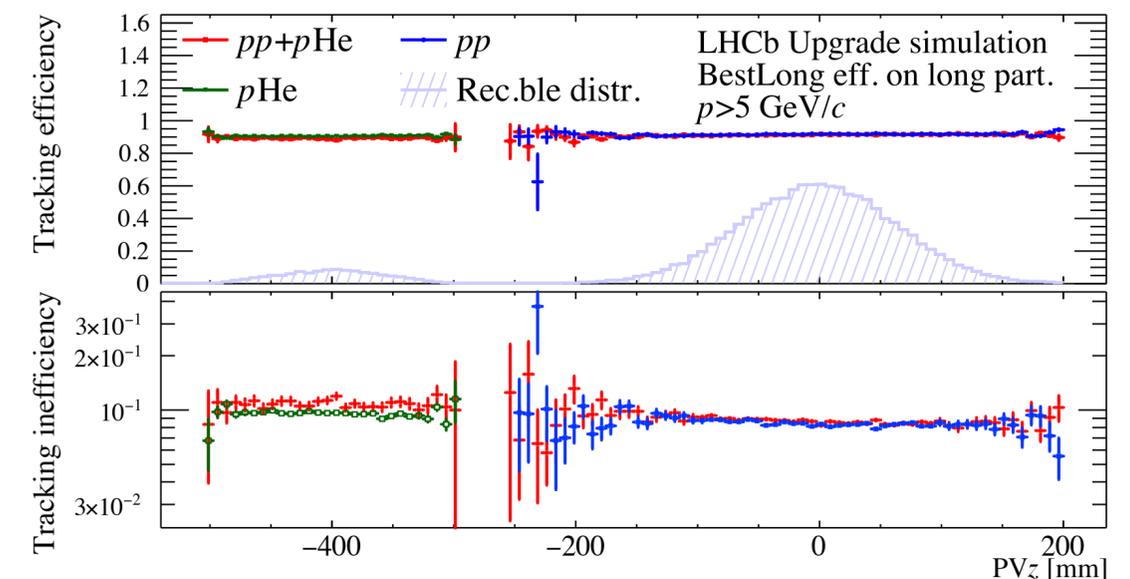
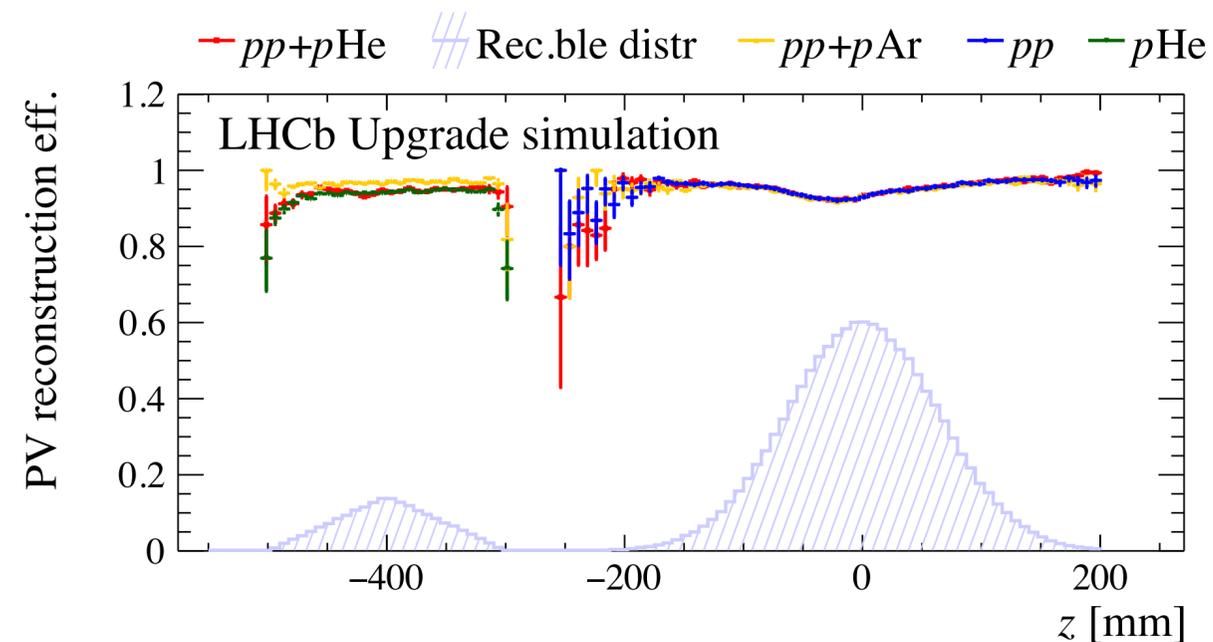
SMOG2/LHCspin performances

[LHCb-FIGURE-2022-002]

- beam-beam and beam-gas interaction regions are well detached
- Negligible increase of multiplicity: 1 – 3 % throughput decrease when adding beam-gas to the LHCb event reconstruction sequence



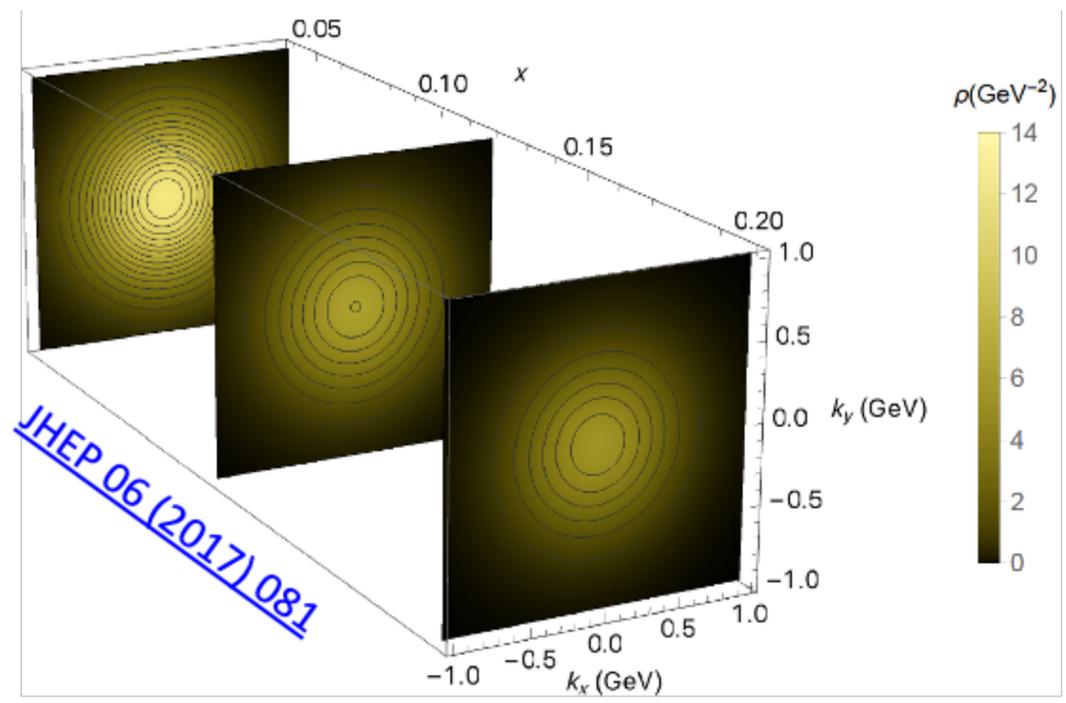
- Full reconstruction efficiency (PV & tracks) retained in the beam-gas region



LHCb is the only experiment able to run in collider and fixed-target mode simultaneously!

The physics goals of

- Multi-dimensional nucleon structure in a poorly explored kinematic domain
- Measure experimental observables sensitive to both **quarks and gluons TMDs**
- **Make use of new probes (charmed and beauty mesons)**
- Complement present and future SIDIS results
- Test non-trivial process dependence of quarks and (especially) gluons TMDs



quark pol.

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

nucleon pol.

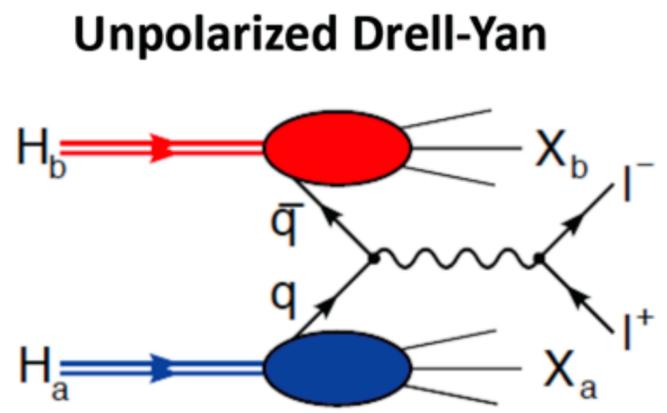
Theoretically cleanest hard h-h scattering process:

- LHCb has excellent $\mu - ID$ & reconstruction for $\mu^+\mu^-$
- dominant:** $\bar{q}(x_{beam}) + q(x_{target}) \rightarrow \mu^+\mu^-$
 - suppressed:** $q(x_{beam}) + \bar{q}(x_{target}) \rightarrow \mu^+\mu^-$

- Sensitive to unpol. and BM TMDs for $q_T \ll M_T$

$$d\sigma_{UU}^{DY} \propto f_1^{\bar{q}} \otimes f_1^q + \cos 2\phi h_1^{\perp, \bar{q}} \otimes h_1^{\perp, q}$$

- H & D targets allow to study the antiquark content of the nucleon
- SeaQuest (E906): $\bar{d}(x) > \bar{u}(x)$ \rightarrow proton sea is not flavour symmetric
- intrinsic heavy quarks?

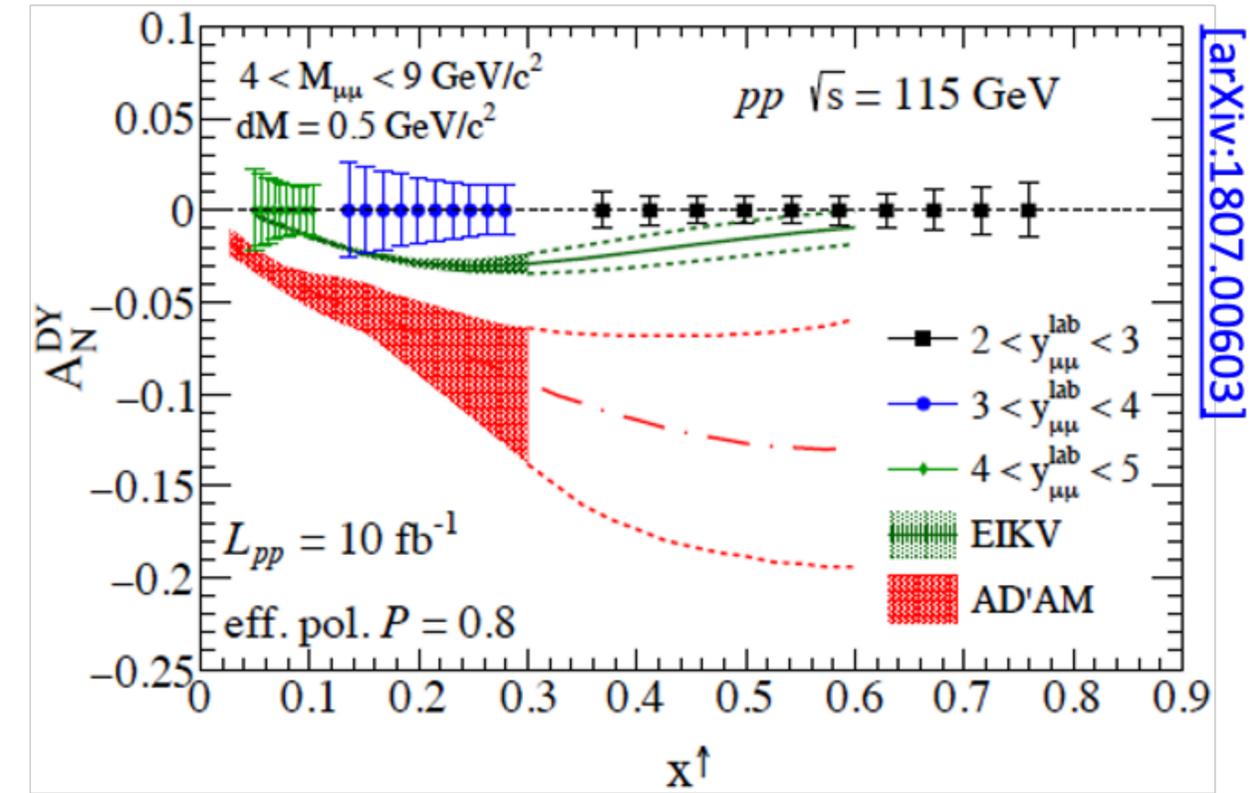
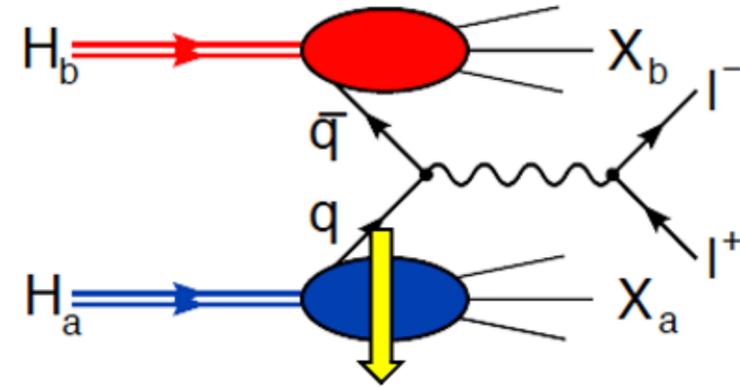


... still a lot to be understood and investigated

Quark TMDs

		quark pol.		
		U	L	T
nucleon pol.	U	f_1		h_1^\perp
	L		g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

Transv. polarized Drell-Yan



- Sensitive to quark TMDs through TSSAs

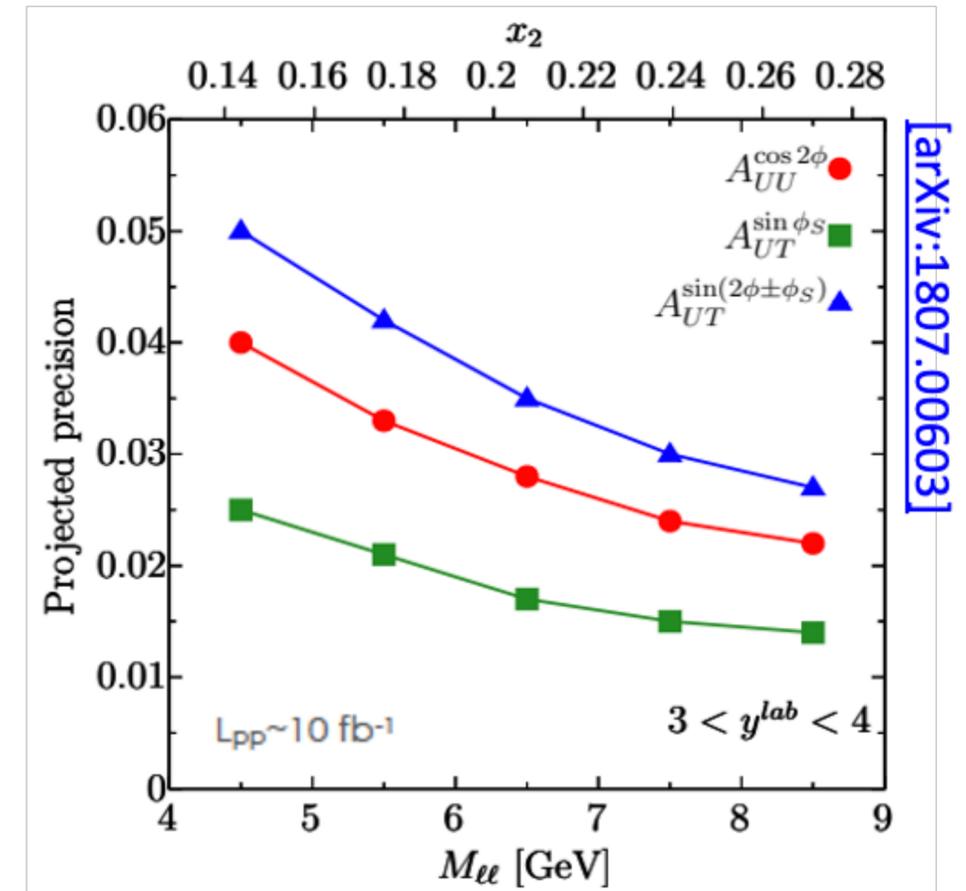
$$A_N^{DY} = \frac{1}{P} \frac{\sigma_{DY}^\uparrow - \sigma_{DY}^\downarrow}{\sigma_{DY}^\uparrow + \sigma_{DY}^\downarrow} \Rightarrow A_{UT}^{\sin\phi_S} \sim \frac{f_1^q \otimes f_{1T}^{\perp q}}{f_1^q \otimes f_1^q}, \quad A_{UT}^{\sin(2\phi - \phi_S)} \sim \frac{h_1^{\perp q} \otimes h_1^q}{f_1^q \otimes f_1^q}, \dots$$

(ϕ : azimuthal orientation of lepton pair in dilepton CM)

- Extraction of qTMDs from DY does not require knowledge of FF
- Verify sign change of Sivers function wrt SIDIS

$$f_{1T}^\perp|_{DY} = -f_{1T}^\perp|_{SIDIS}$$

- Test flavour sensitivity using both H and D targets

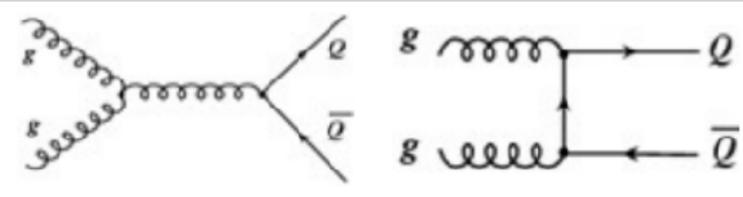


Probing the gTMDs

		gluon pol.		
		U	Circularly	Linearly
nucleon pol.	U	f_1^g		$h_1^{\perp g}$
	L		g_{1L}^g	$h_{1L}^{\perp g}$
	T	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g, h_{1T}^{\perp g}$

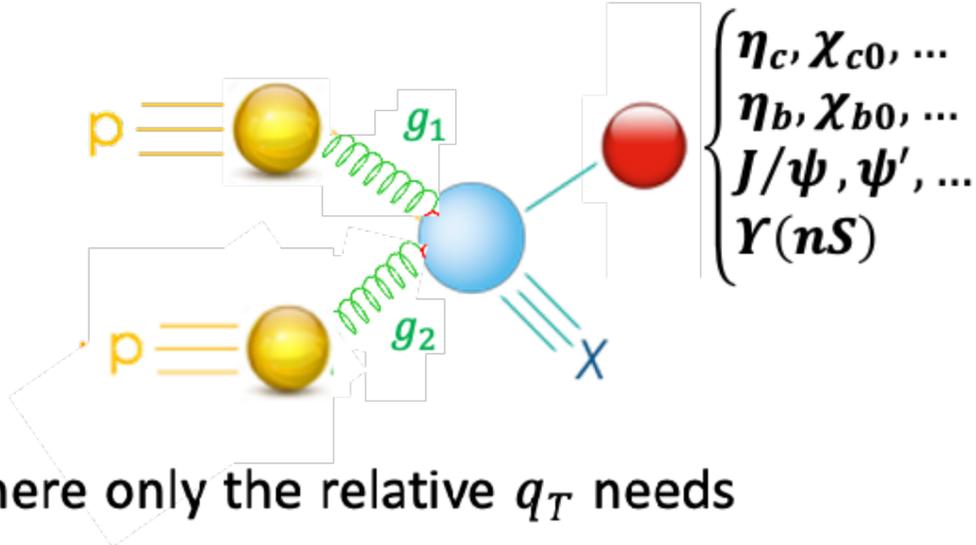
Theory framework well consolidated ...**but experimental access still extremely limited!**

In high-energy hadron collisions, heavy quarks are dominantly produced by gg fusion:



The most efficient way to access the gluon dynamics inside the proton at LHC is to **measure heavy-quark observables**

- **Inclusive quarkonia production in (un)polarized pp interaction** ($pp^{(\uparrow)} \rightarrow [Q\bar{Q}]X$) turns out to be an ideal observable to access gTMDs (assuming TMD factorization)



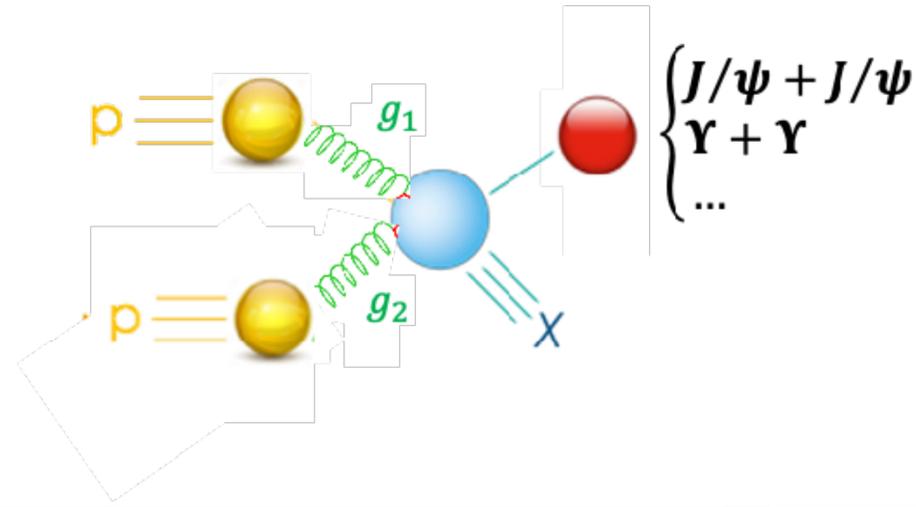
- TMD factorization requires $q_T(Q) \ll M_Q$. Can look at **associate quarkonia production**, where only the relative q_T needs to be small:

$$\text{E.g.: } pp^{(\uparrow)} \rightarrow J/\psi + J/\psi + X$$

- Due the larger masses this condition is more easily matched in the case of **bottomonium**, where TMD factorization can hold at larger q_T (although very challenging for experiments!)

Probing the gTMDs

$$\frac{d\sigma}{dM_{QQ}dY_{QQ}d^2P_{QQT}d\Omega} = \frac{\sqrt{M_{QQ}^2 - 4M_Q^2}}{(2\pi)^2 8s M_{QQ}^2} \left\{ \begin{aligned} &F_1(M_{QQ}, \theta_{CS}) C[f_1^g f_1^g](x_{1,2}, P_{QQT}) \\ &+ F_2(M_{QQ}, \theta_{CS}) C[w_2 h_1^{\perp g} h_1^{\perp g}](x_{1,2}, P_{QQT}) \\ &+ \left(F_3(M_{QQ}, \theta_{CS}) C[w_3 f_1^g h_1^{\perp g}](x_{1,2}, P_{QQT}) + F'_3(M_{QQ}, \theta_{CS}) C[w'_3 h_1^{\perp g} f_1^g](x_{1,2}, P_{QQT}) \right) \cos 2\phi_{CS} \\ &+ F_4(M_{QQ}, \theta_{CS}) C[w_4 h_1^{\perp g} h_1^{\perp g}](x_{1,2}, P_{QQT}) \cos 4\phi_{CS} \end{aligned} \right\}$$

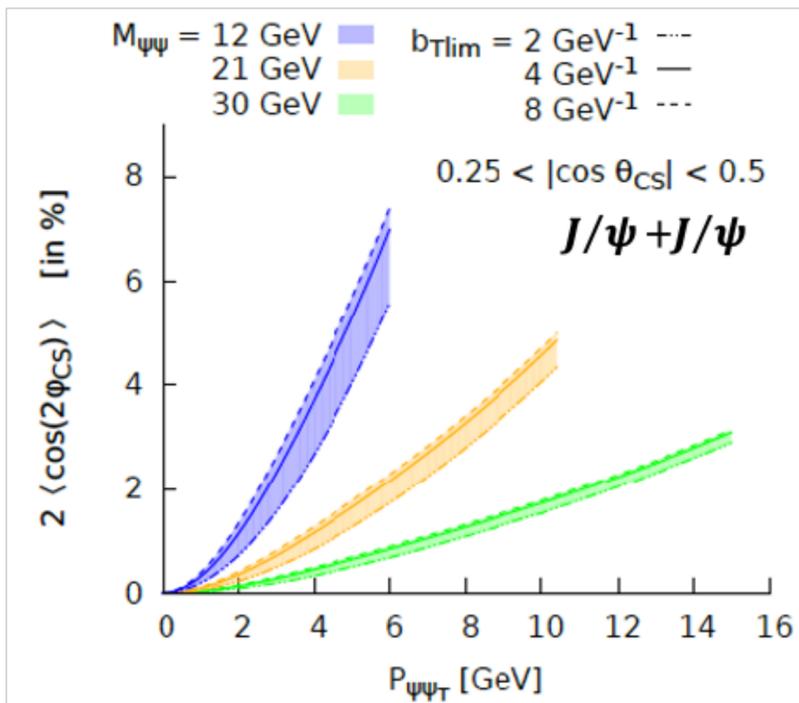


gluon pol.

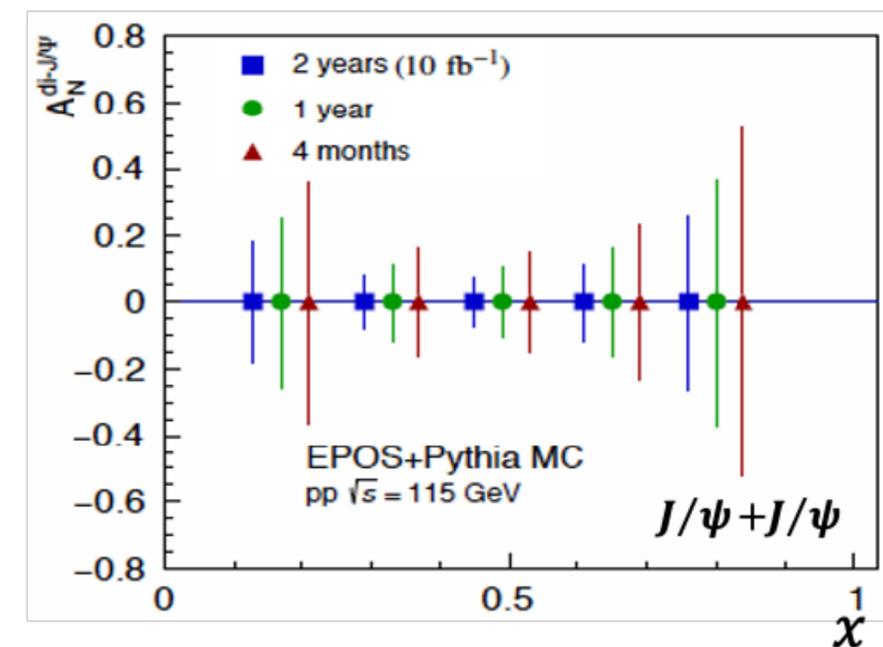
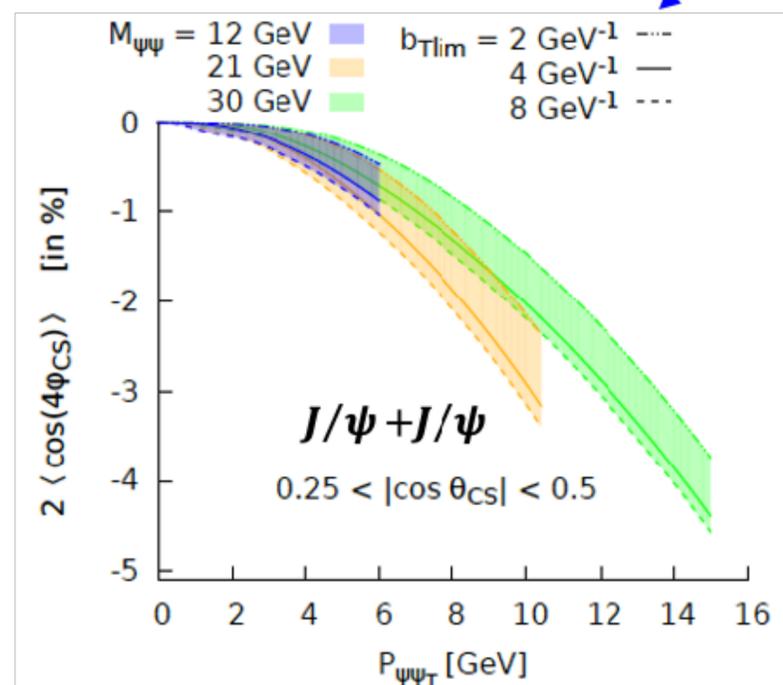
	U	Circularly	Linearly
U	f_1^g		$h_1^{\perp g}$
L		g_{1L}^g	$h_{1L}^{\perp g}$
T	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g, h_{1T}^{\perp g}$

nucleon pol.

Predictions based on CSM + TMD evolution for $x_1 \sim x_2 \sim 10^{-3}$ at forward rapidity [EPJ C 80, 87 (2020)]

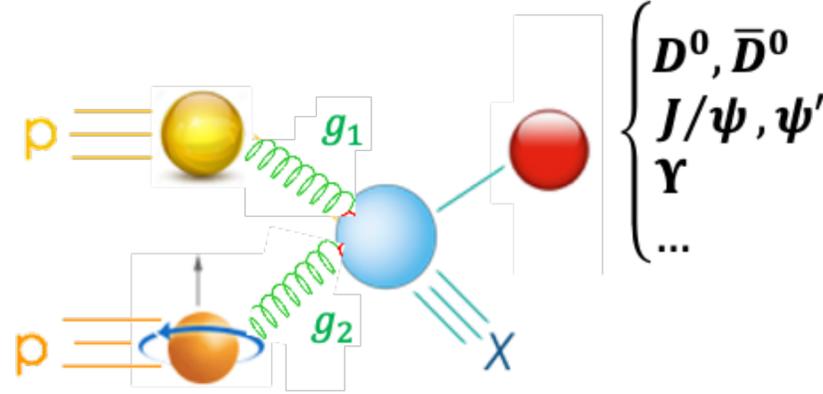


Azimuthal
amplitudes
~5%!



Probing the gluon Sivers function

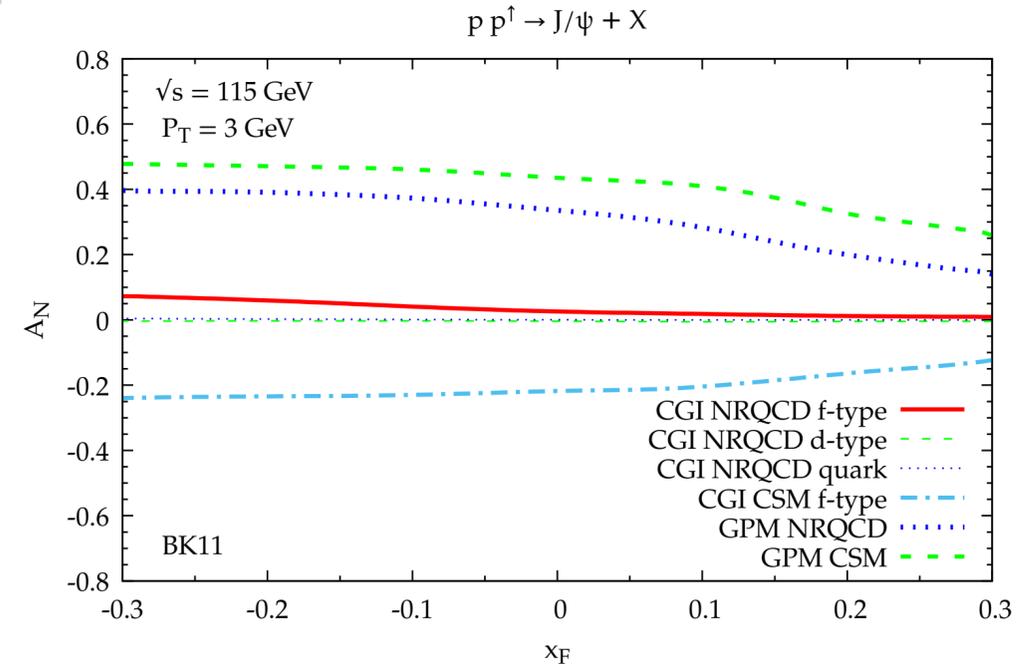
$$\Gamma_T^{\mu\nu}(x, \mathbf{p}_T) = \frac{x}{2} \left\{ g_T^{\mu\nu} \frac{\epsilon_T^{\rho\sigma} p_{T\rho} S_{T\sigma}}{M_p} f_{1T}^{\perp g}(x, \mathbf{p}_T^2) + \dots \right\}$$



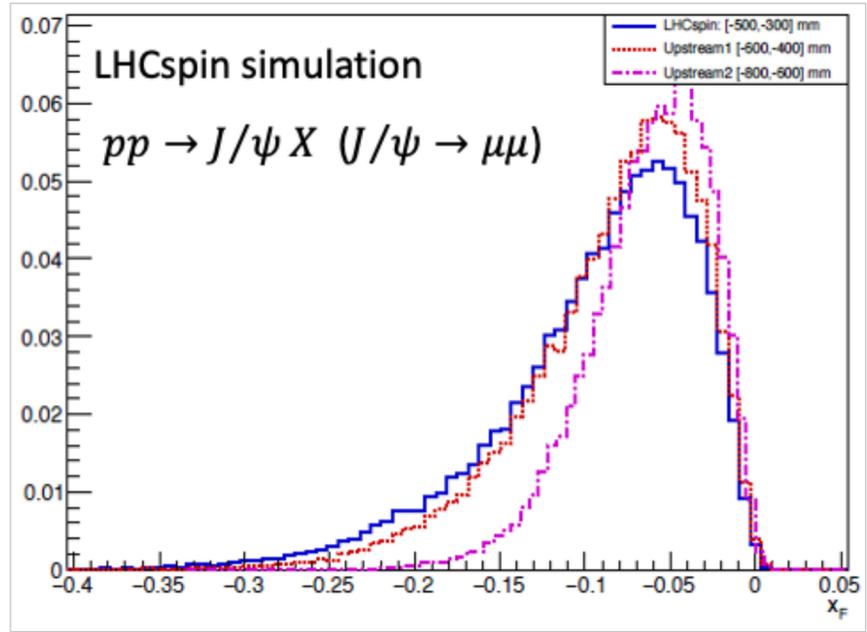
		gluon pol.		
		U	Circularly	Linearly
nucleon pol.	U	f_1^g		$h_1^{\perp g}$
	L		g_{1L}^g	$h_{1L}^{\perp g}$
	T	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g, h_{1T}^{\perp g}$

- Sheds light on spin-orbit correlations of unpol. gluons inside a transv. pol. proton
- sensitive to color exchange among IS and FS and gluon OAM
- expected to be small (quasi-saturation of Burkardt sum rule by $f_{1T}^{\perp g}$ and QCD predictions in large- N_c limit)
- can be accessed through the Fourier decomposition of the TSSAs for **inclusive heavy meson production**

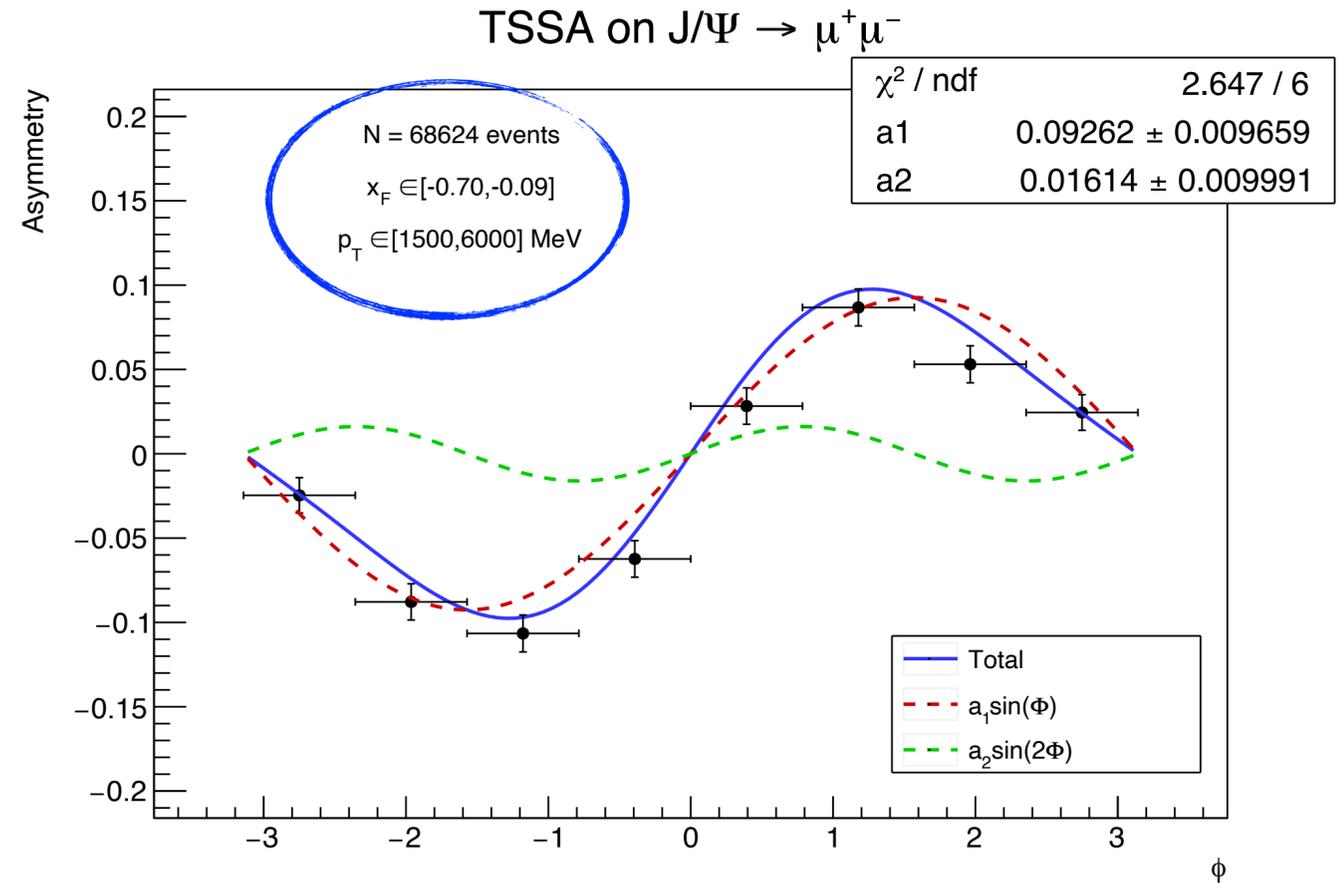
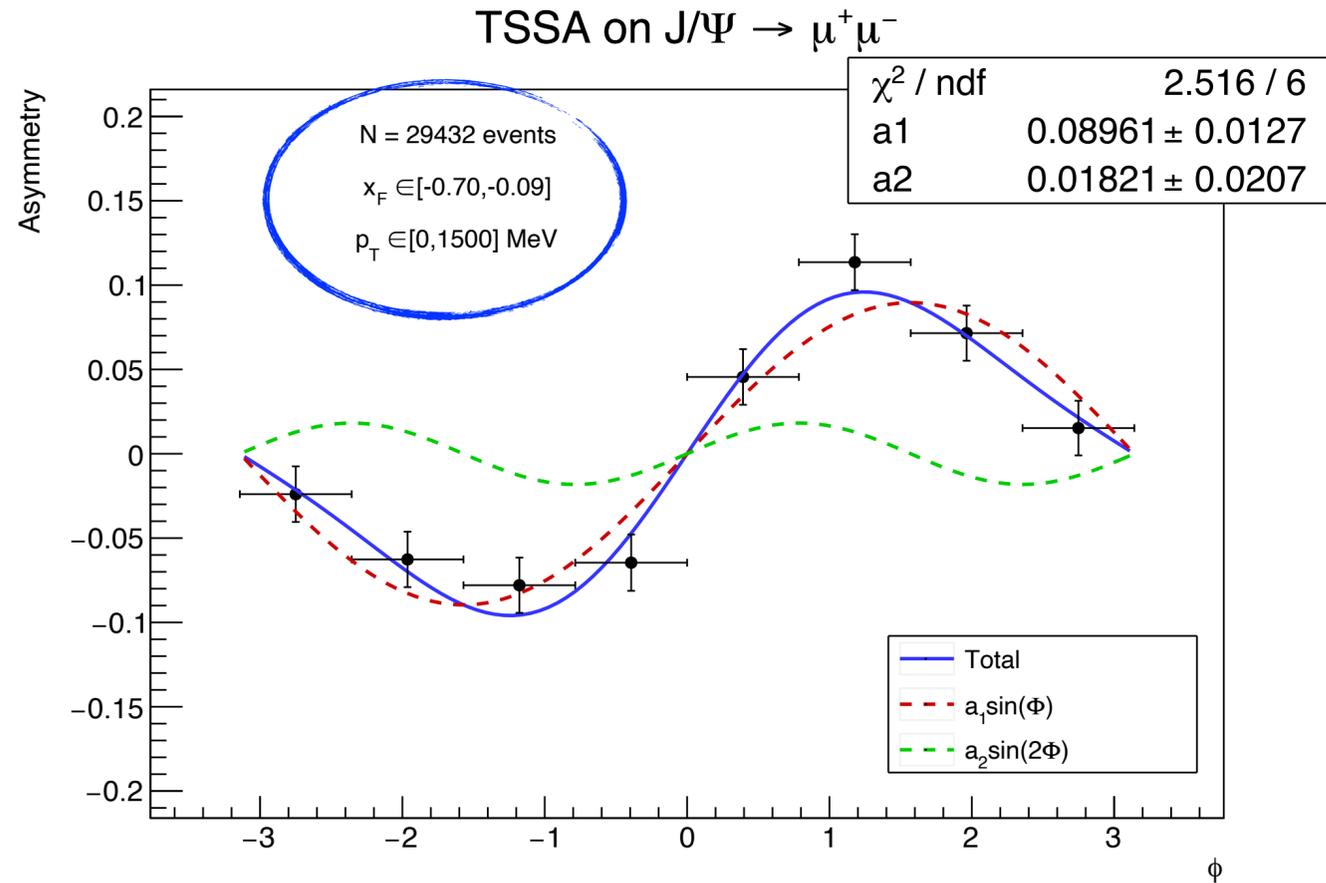
$$A_N = \frac{1}{P} \frac{\sigma^\uparrow - \sigma^\downarrow}{\sigma^\uparrow + \sigma^\downarrow} \propto [f_{1T}^{\perp g}(x_a, k_{\perp a}) \otimes f_g(x_b, k_{\perp b}) \otimes d\sigma_{gg \rightarrow QQg}] \sin \phi_S + \dots$$



- **Predictions for pol. FT meas. at LHC (LHCspin-like)**
- Phys. Rev. D 102, 094011 (2020)
- $pp^\uparrow \rightarrow J/\psi + X$
- **based on GPM & CGI-GPM**
- **Expected amplitudes could reach 5-10% in the $x_F < 0$ region**



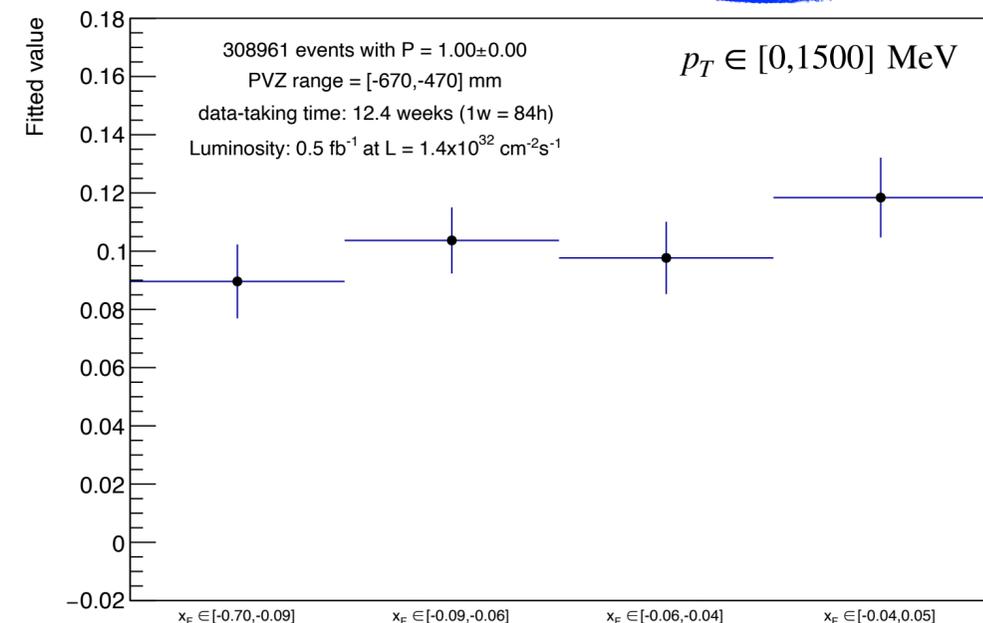
A TSSA analysis at LHCspin with $J/\Psi \rightarrow \mu^+\mu^-$ events



[[JHEP 12 \(2020\) 010](#)]

- Full LHCb simulations of $J/\Psi \rightarrow \mu^+\mu^-$ in pP collisions \rightarrow emulate the target polarisation by assigning a $\uparrow \downarrow$ tag according to a given model. In this example: 10% asymmetry on $\sin \phi$, 2% on $\sin 2\phi$ + mild x_F, p_T dependence
- Fit the polarised data with the sum of two Fourier amplitudes (a_1, a_2) in $4 x_F \times 2 p_T \times 8 \phi$ bins
- Within this statistics, corresponding to ~ 3 months of data-taking, $A_N \sim 0.1 \pm 0.01$

$J/\Psi \rightarrow \mu^+\mu^-$: fit results for parameter a_1



Knowledge of the polarisation degree

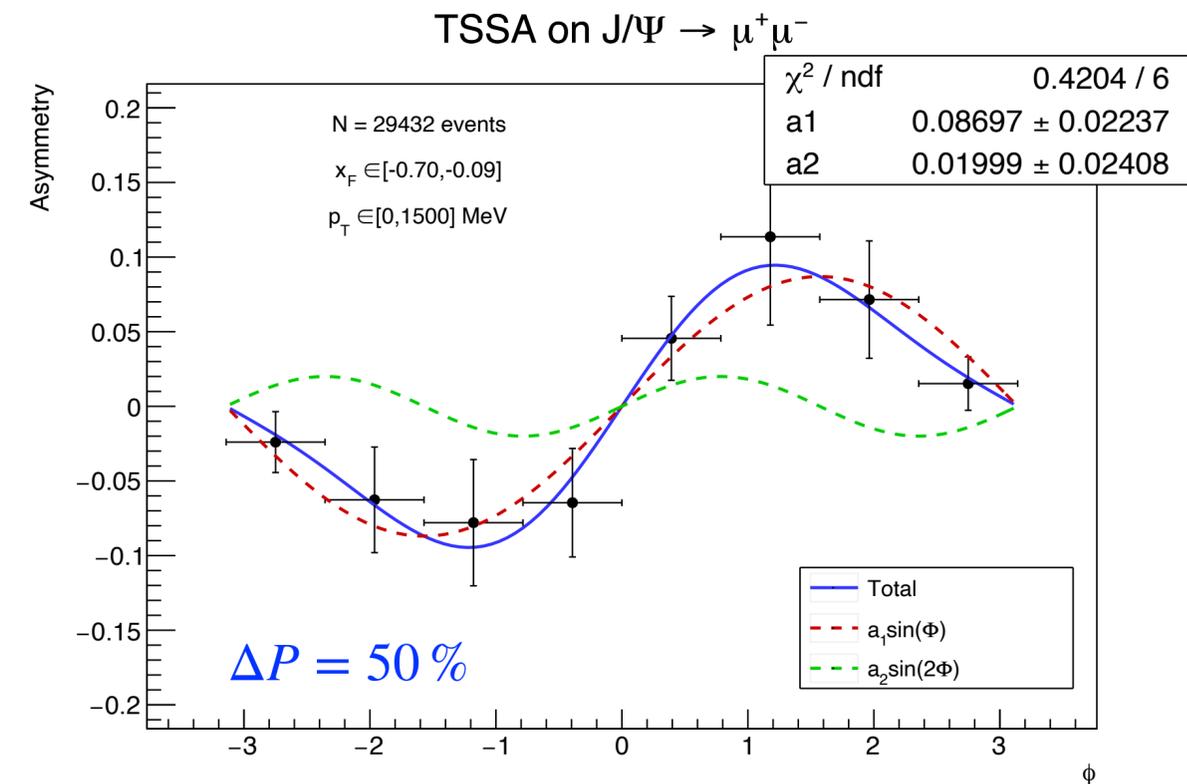
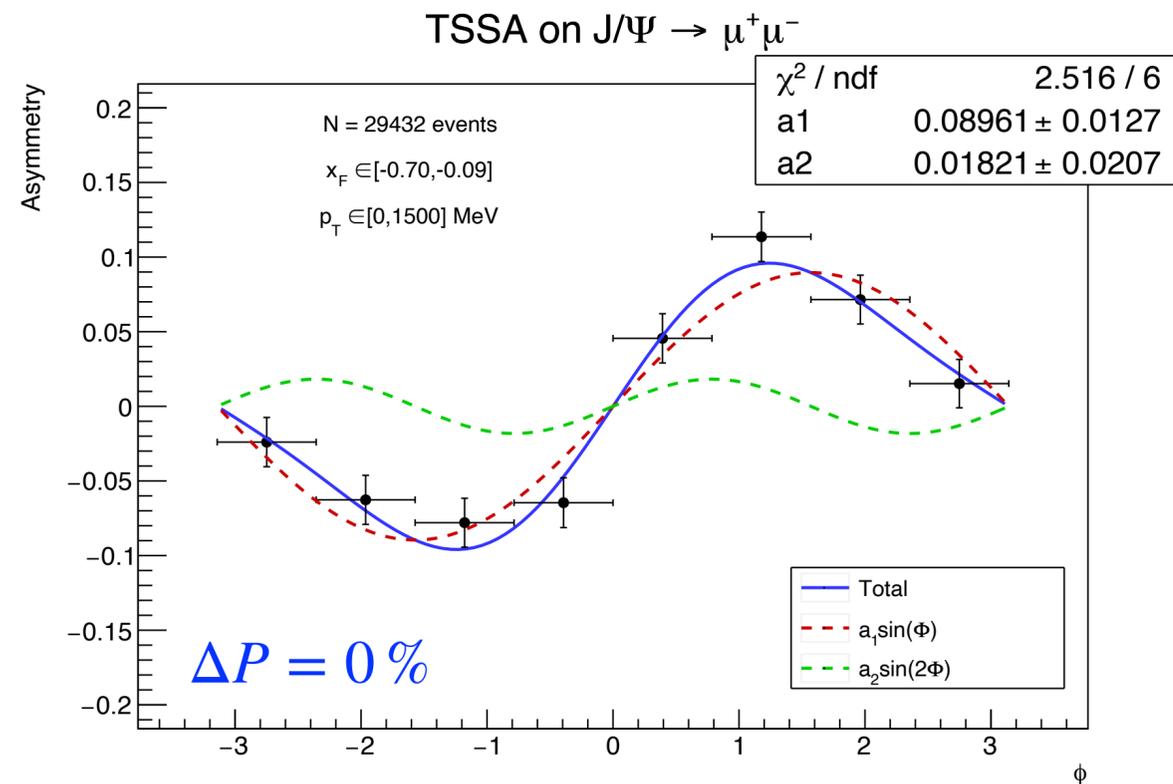
- To estimate the systematic error due to the measurement of the polarisation degree, the analysis is repeated with different ΔP
- Very relevant for the R&D (e.g. cell vs jet target). With the shown analysis* :
- 5% error (realistic value) \rightarrow negligible effect
- 20% error \rightarrow 30-40% of the stat. error
- 50% error \rightarrow syst. dominated

$\Delta P = 5\%$

p_T (MeV)	x_F	a_1
[0,1500]	[-0.70,-0.09]	0.089 ± 0.013
[0,1500]	[-0.09,-0.06]	0.104 ± 0.012
[0,1500]	[-0.06,-0.04]	0.098 ± 0.013
[0,1500]	[-0.04,0.05]	0.117 ± 0.014
[1500,6000]	[-0.70,-0.09]	0.092 ± 0.010
[1500,6000]	[-0.09,-0.06]	0.108 ± 0.011
[1500,6000]	[-0.06,-0.04]	0.105 ± 0.012
[1500,6000]	[-0.04,0.05]	0.105 ± 0.012

$\Delta P = 20\%$

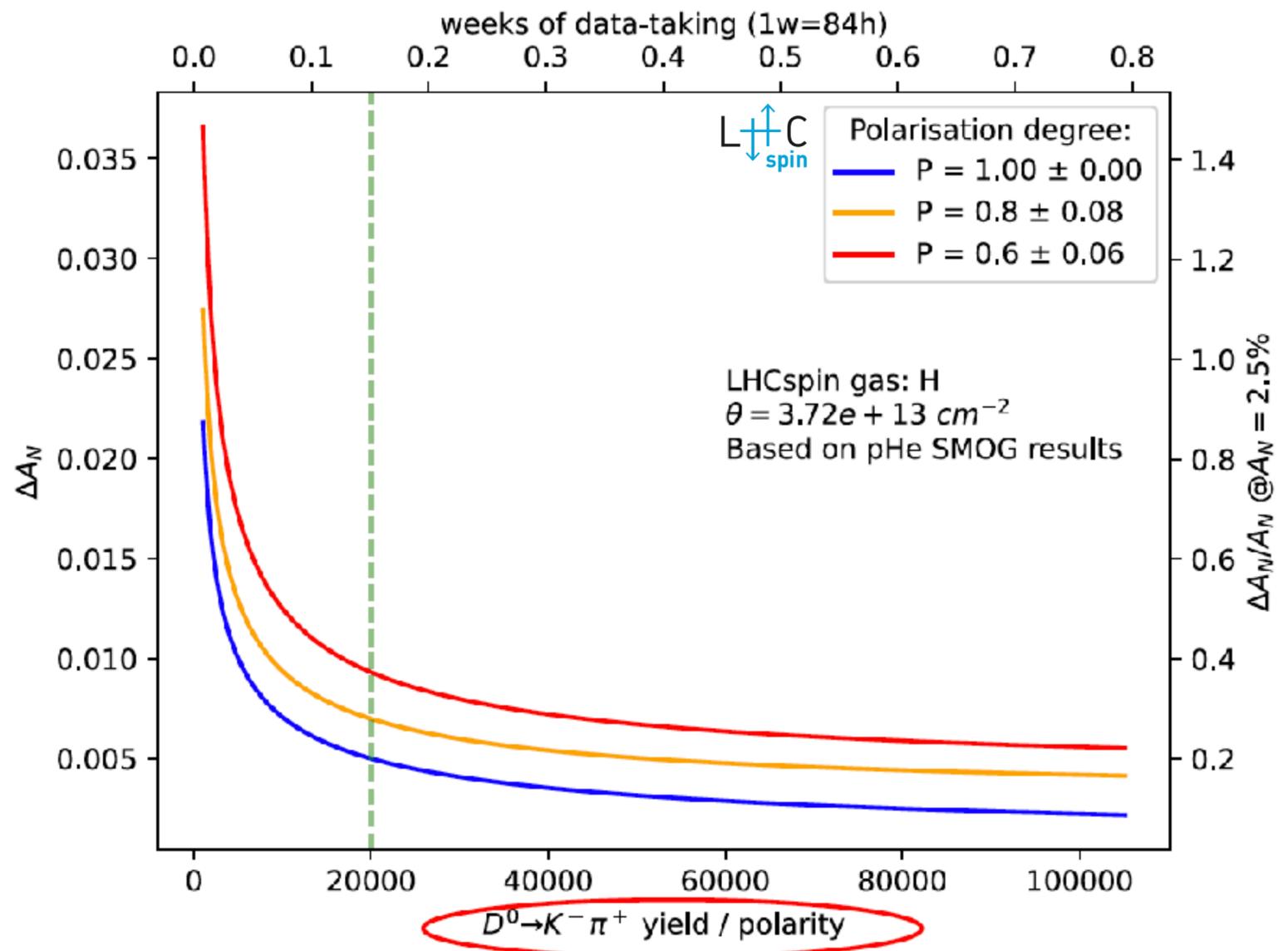
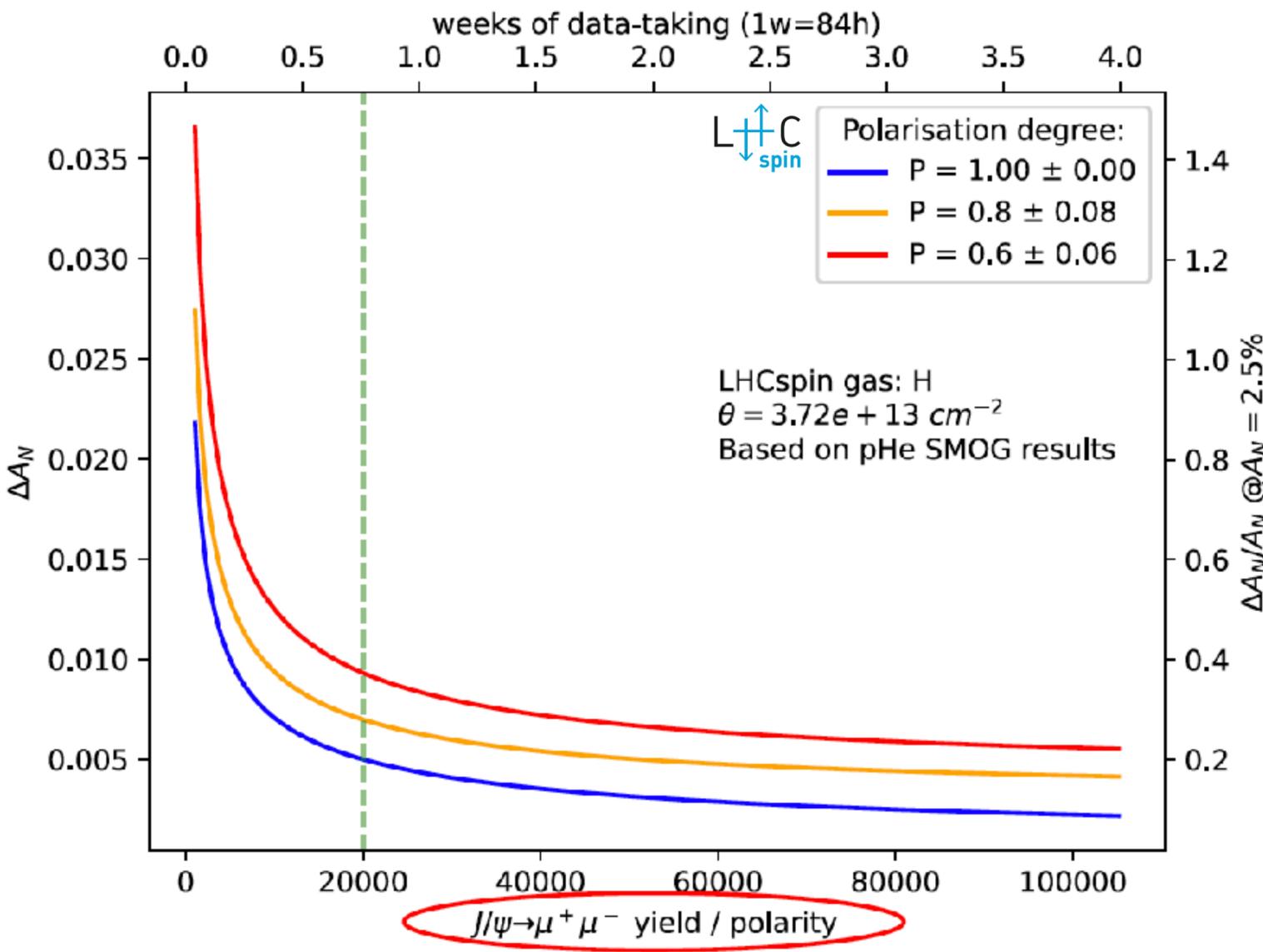
p_T (MeV)	x_F	a_1
[0,1500]	[-0.70,-0.09]	0.087 ± 0.014
[0,1500]	[-0.09,-0.06]	0.103 ± 0.016
[0,1500]	[-0.06,-0.04]	0.097 ± 0.016
[0,1500]	[-0.04,0.05]	0.114 ± 0.017
[1500,6000]	[-0.70,-0.09]	0.090 ± 0.013
[1500,6000]	[-0.09,-0.06]	0.108 ± 0.015
[1500,6000]	[-0.06,-0.04]	0.104 ± 0.015
[1500,6000]	[-0.04,0.05]	0.102 ± 0.015



* i.e. \sim 3 months of data-taking with this example model, channel and kinematic binning

LHCspin event rates

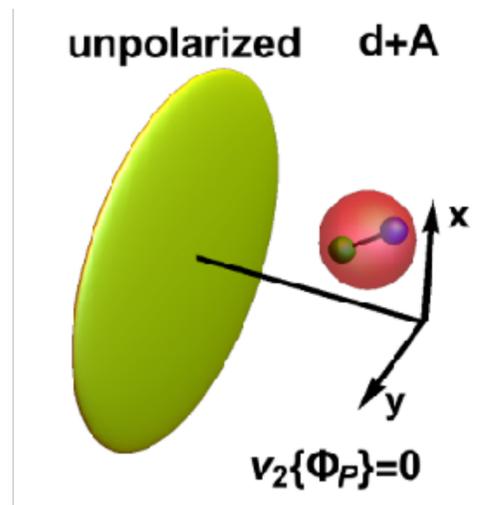
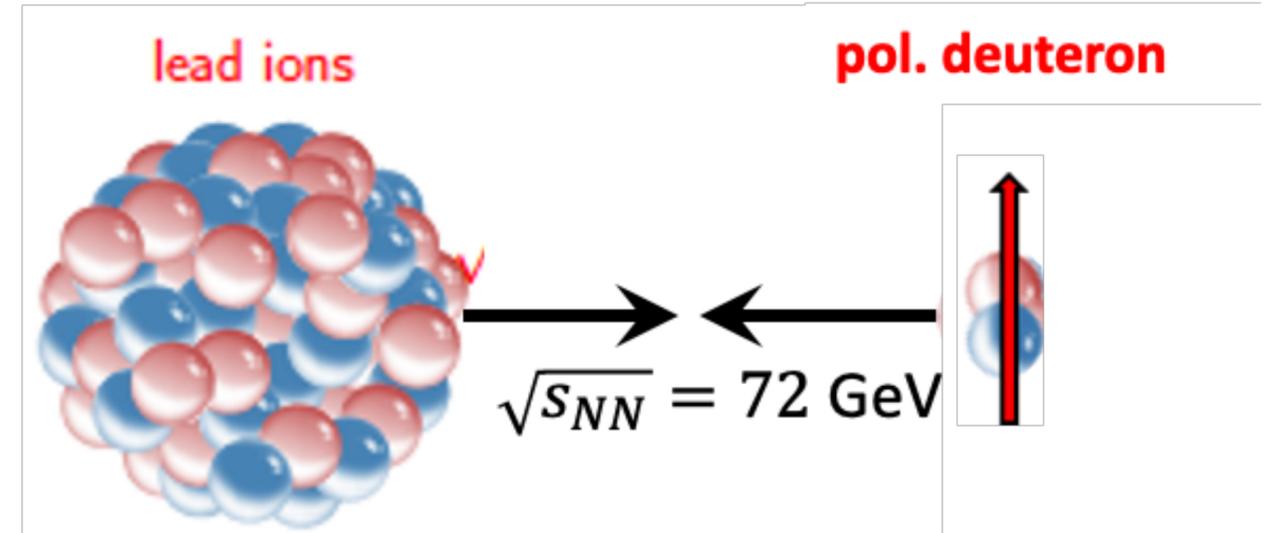
Precise spin asymmetry on $J/\Psi \rightarrow \mu^+ \mu^-$ and $D^0 \rightarrow K^- \pi^+$ for pH^\uparrow collisions in just few weeks with Run3 luminosity!
 Statistics further enhanced by a factor 3-5 in LHCb upgrade II



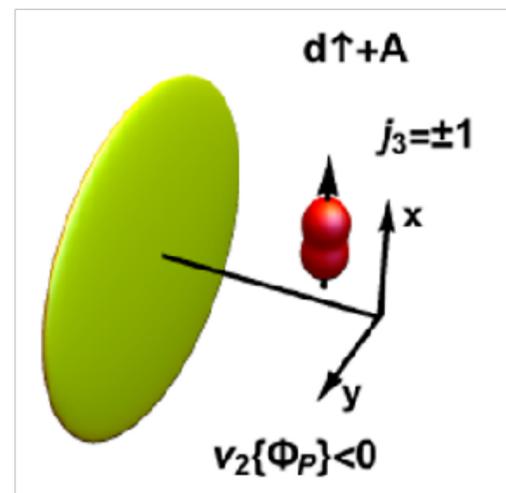
reconstructed particles

Spin physics in heavy-ion collisions

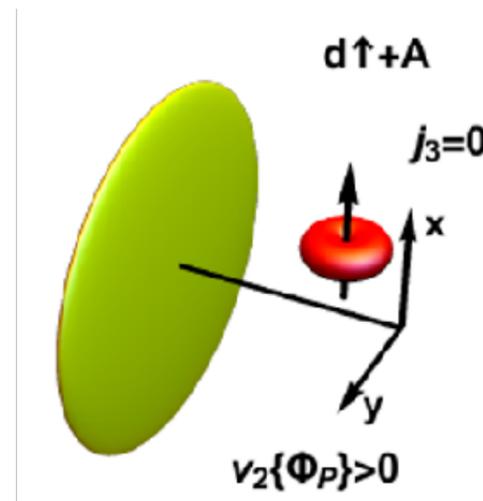
- probe collective phenomena in heavy-light systems through **ultra-relativistic collisions of heavy nuclei with trasv. pol. deuterons**
- polarized light target nuclei offer a unique opportunity to control the orientation of the formed fireball by measuring the **elliptic flow** relative to the polarization axis (**ellipticity**).



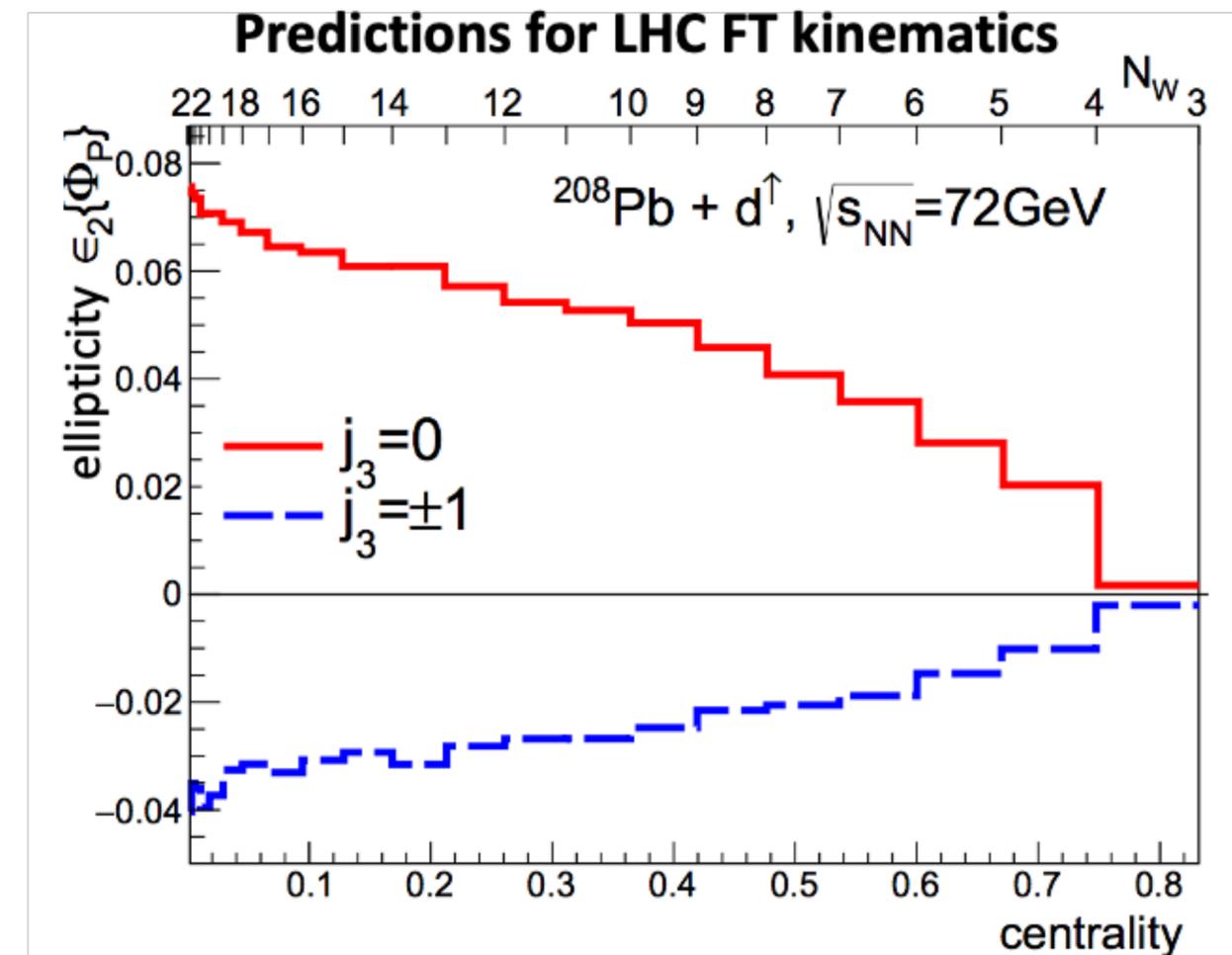
Unpol. deuterons: the fireball is azimuthally symmetric and $v_2 \approx 0$.



$j_3 = \pm 1 \rightarrow$ prolate fireball stretched along the pol. axis, corresponds to $v_2 < 0$



$j_3 = 0 \rightarrow$ oblate fireball corresponds to $v_2 > 0$



[PRC 101 (2020) 024901]

International framework and feedback

Several experiments dedicated to spin physics, but with many limitations:

very low energy, no rare probes, no ion beam, ...

➡ LHCspin is unique in this respect

LHCspin is complementary to EIC

[D. Boer: [arXiv:1611.06089](https://arxiv.org/abs/1611.06089)]

unpolarized gluon TMD

	DIS	DY	SIDIS	$pA \rightarrow \gamma \text{ jet } X$	$ep \rightarrow e' Q \bar{Q} X$ $ep \rightarrow e' j_1 j_2 X$	$pp \rightarrow \eta_{c,b} X$ $pp \rightarrow H X$	$pp \rightarrow J/\psi \gamma X$ $pp \rightarrow \Upsilon \gamma X$
$f_1^g^{[+,+]}$ (WW)	×	×	×	×	✓	✓	✓
$f_1^g^{[+,-]}$ (DP)	✓	✓	✓	✓	×	×	×

linearly polarized gluon TMD

	$pp \rightarrow \gamma \gamma X$	$pA \rightarrow \gamma^* \text{ jet } X$	$ep \rightarrow e' Q \bar{Q} X$ $ep \rightarrow e' j_1 j_2 X$	$pp \rightarrow \eta_{c,b} X$ $pp \rightarrow H X$	$pp \rightarrow J/\psi \gamma X$ $pp \rightarrow \Upsilon \gamma X$
$h_1^{\perp g [+,+]}$ (WW)	✓	×	✓	✓	✓
$h_1^{\perp g [+, -]}$ (DP)	×	✓	×	×	×

TMDs (Sivers)

[D. Boer: [arXiv:1611.06089](https://arxiv.org/abs/1611.06089), D. Boer et al. HEPJ 08 2016 001]

	DY	SIDIS	$p^\dagger A \rightarrow h X$	$p^\dagger A \rightarrow \gamma^{(*)} \text{ jet } X$	$p^\dagger p \rightarrow \gamma \gamma X$ $p^\dagger p \rightarrow J/\psi \gamma X$ $p^\dagger p \rightarrow J/\psi J/\psi X$	$ep^\dagger \rightarrow e' Q \bar{Q} X$ $ep^\dagger \rightarrow e' j_1 j_2 X$
$f_{1T}^{\perp g [+,+]}$ (WW)	×	×	×	×	✓	✓
$f_{1T}^{\perp g [+, -]}$ (DP)	✓	✓	✓	✓	×	×

$f_{1T}^{\perp g [+,+]}$ (Weizsacker-Williams type or "f-type") → antisymmetric colour structures

$f_{1T}^{\perp g [+, -]}$ (Dipole s type or "d-type") → symmetric colour structures

- ☐ Can be measured at the Electron Ion-Collider (EIC)
- ☐ Can be measured at LHCspin

"Ambitious and long term LHC-Fixed Target research program. The efforts of the existing LHC experiments to implement such a programme, including specific R&D actions on the collider, **deserve support**" (European Strategy for Particle Physics)

"This would be **unique and highly complementary** to existing and future measurements in lepton-proton collisions, because the asymmetries in question have a process dependence between pp and lp that is predicted by theory" (CERN Physics Beyond Collider)

Recognised relevance

solid target @

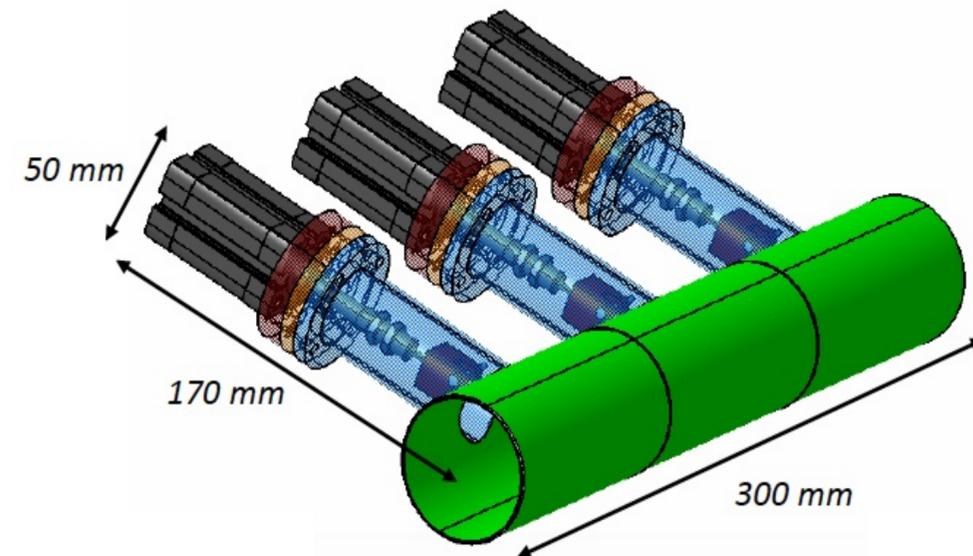


ALICE

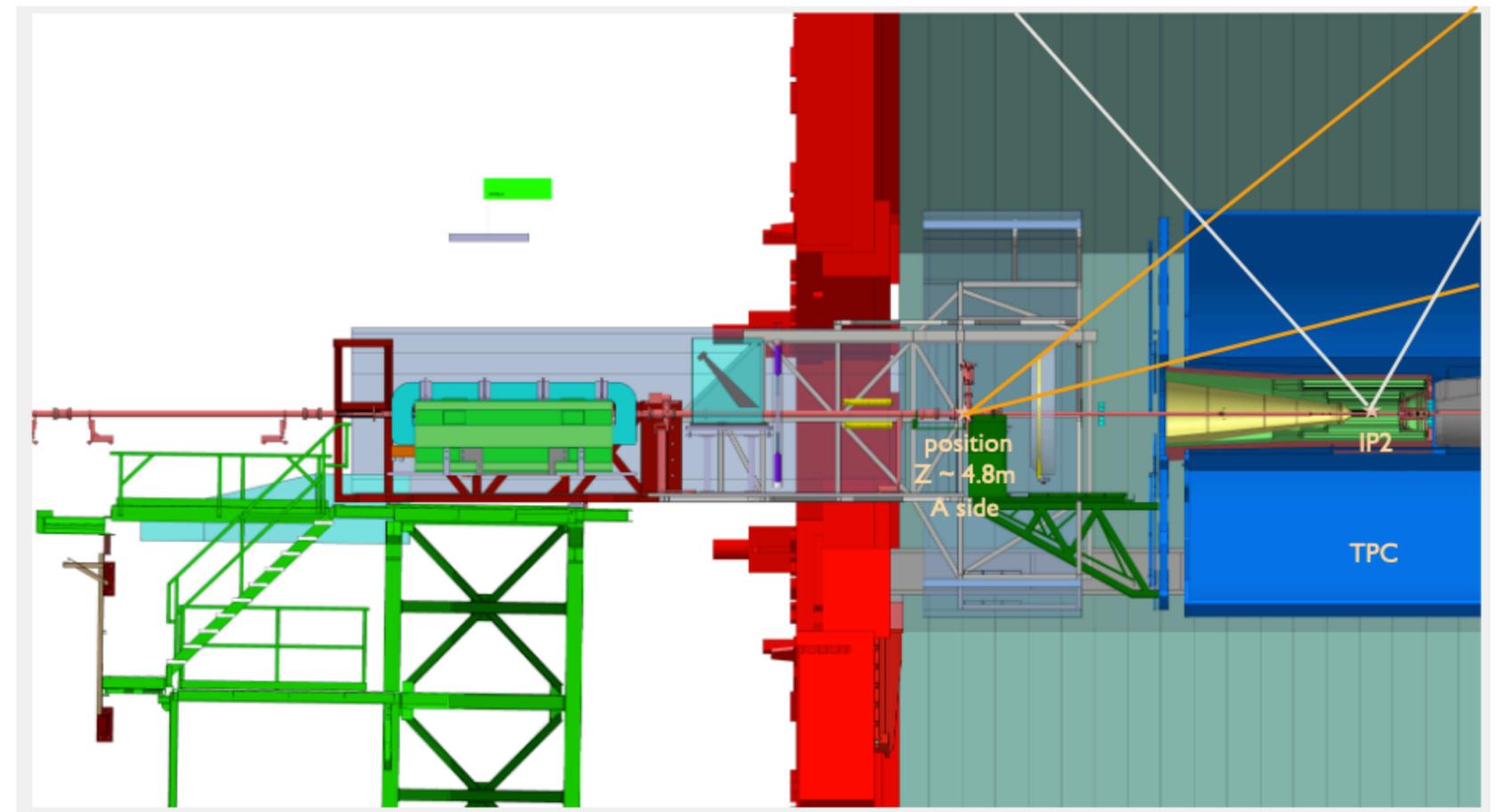
The ALICE unpolarised solid target

Two main physics goals:

- Advance our understanding of the large- x gluon, antiquark and heavy-quark content in the nucleon and nucleus (structure of nucleon and nuclei at large- x , gluon EMC effect in nuclei, intrinsic charm in nucleon)
 - Study heavy-ion collisions between SPS and RHIC energies towards large rapidities (longitudinal expansion of QGP formation, collectivity in small systems with heavy quarks, factorisation of CNM effects)
- Proton beam halo channelled with a bent crystal on a retractable solid target (C,W, Ti...)
 - Backward cms rapidity coverage with forward detectors in the lab thanks to the boost

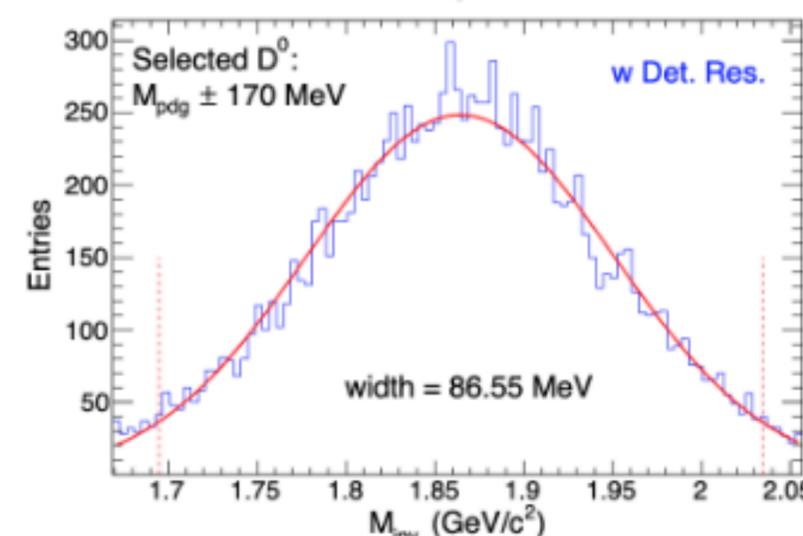
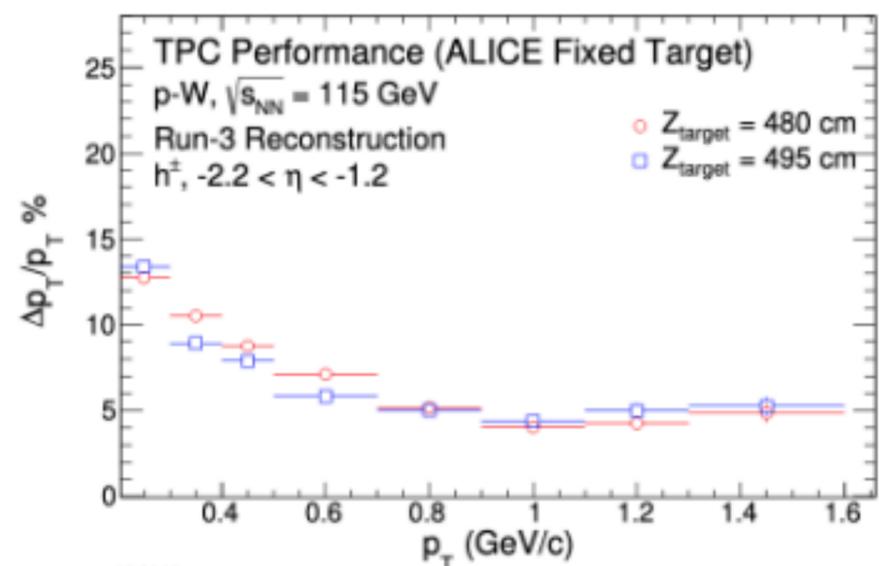
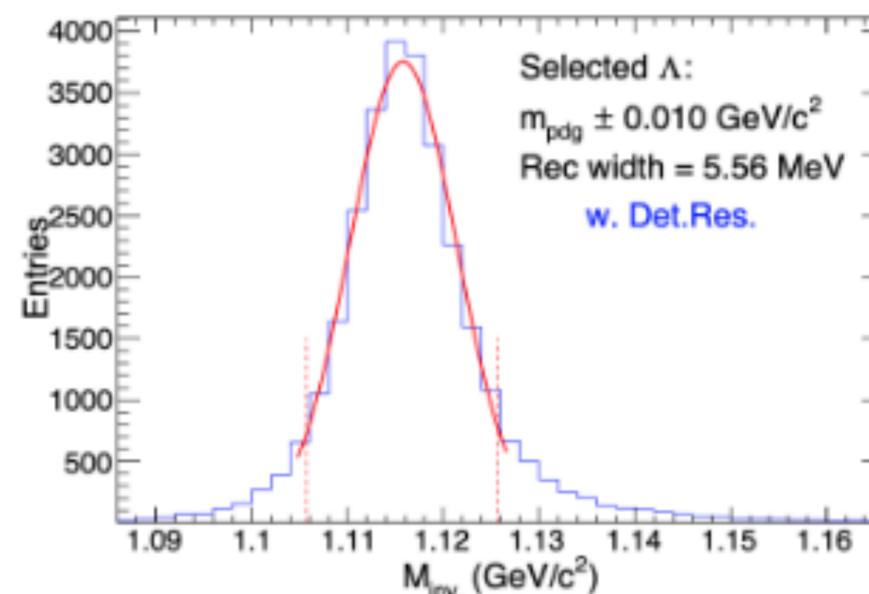
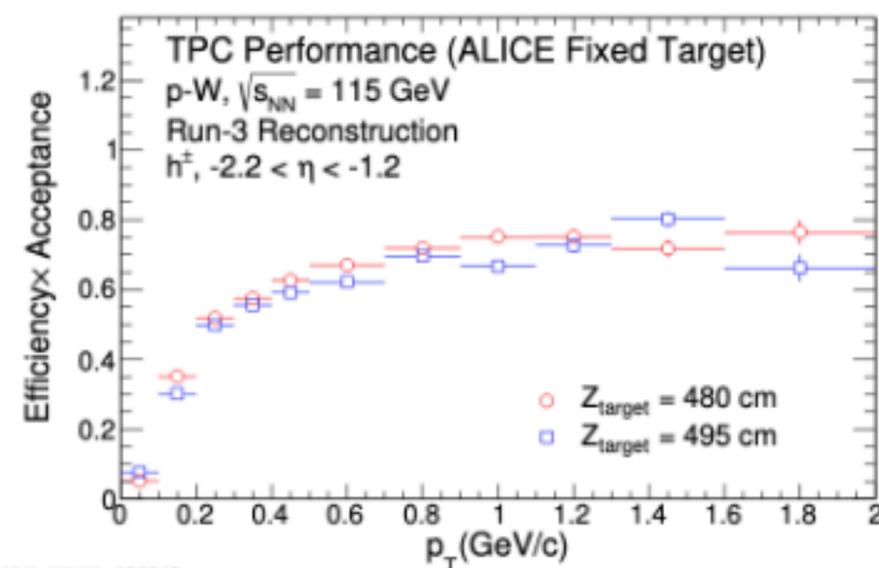


retractable solid target



Some of the performances

space availability at $z = 3259$ m



Proton beam collimation studies performed:
 loss maps, positioning of the crystal system and
 of the absorbers

LOI in ALICE (2022) —> aim for
 installation during LS3
 (2026-2028)

Some of the results achieved

Λ : efficiency and p_T resolution sufficient for analysis (without extra vertex detector)

D^0 : TPC vertex resolution not sufficient to use secondary vertex method for analysis. Investigating combinatorial background method, reduced target size and constraints on beam spot position for tracking

Integration solutions to comply with FOCAL and ITS motion constraints during EYETS

Physics performance with realistic detector conditions

Conclusions



Today **SMOG2**



Fixed target physics at LHC is an exiting reality



ALICE

has potentialities in the unpolarised case showing complementarity to LHCb



SMOG2

already operative and taking unpolarised data



is an innovative and unique project conceived to bring polarized physics at the LHC. It is extremely ambitious in terms of both physics reach and technical complexity. It could be installed in a realistic time schedule and costs