

# The chiral anomaly and the eta-prime in vacuum and at low temperatures

Stefan Leupold,  
Carl Niblaeus, Elisabetta Perotti, Bruno Strandberg

Department of Physics and Astronomy  
Uppsala University

Reimei workshop, February 2022



# Table of Contents

- 1 Introduction
- 2 Consequences of chiral anomaly in vacuum
- 3 What changes in a medium?
  - Thermal width of  $\eta'$  meson
- 4 Summary and outlook
  - Consequences for finite density



# Introduction

- neglecting quark masses smaller than  $\Lambda_{\text{QCD}}$



# Introduction

- neglecting quark masses smaller than  $\Lambda_{\text{QCD}}$
- ↪ chiral  $U_R(3) \times U_L(3)$  symmetry of QCD Lagrangian
- ↪ subgroups ( $U_R(3) \times U_L(3) = U_V(3) \times U_A(3)$ ):



# Introduction

- neglecting quark masses smaller than  $\Lambda_{\text{QCD}}$
- $\rightsquigarrow$  chiral  $U_R(3) \times U_L(3)$  symmetry of QCD Lagrangian
- $\hookrightarrow$  subgroups ( $U_R(3) \times U_L(3) = U_V(3) \times U_A(3)$ ):
  - $U_V(1)$ : baryon number conservation



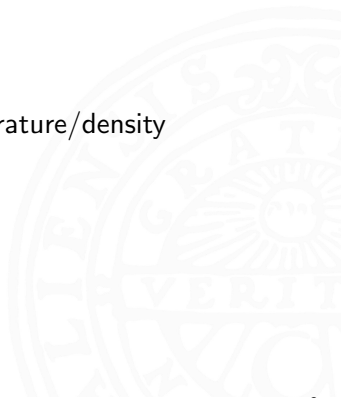
# Introduction

- neglecting quark masses smaller than  $\Lambda_{\text{QCD}}$
- ↪ chiral  $U_R(3) \times U_L(3)$  symmetry of QCD Lagrangian
- ↪ subgroups ( $U_R(3) \times U_L(3) = U_V(3) \times U_A(3)$ ):
  - $U_V(1)$ : baryon number conservation
  - $SU_V(3)$ : flavor multiplets



# Introduction

- neglecting quark masses smaller than  $\Lambda_{\text{QCD}}$
- ↪ chiral  $U_R(3) \times U_L(3)$  symmetry of QCD Lagrangian
- ↪ subgroups ( $U_R(3) \times U_L(3) = U_V(3) \times U_A(3)$ ):
  - $U_V(1)$ : baryon number conservation
  - $SU_V(3)$ : flavor multiplets
  - $SU_A(3)$ : spontaneously broken
- ↪ 8 Goldstone bosons in vacuum and chiral restoration at some finite temperature/density



# Introduction

- neglecting quark masses smaller than  $\Lambda_{\text{QCD}}$
- ↪ chiral  $U_R(3) \times U_L(3)$  symmetry of QCD Lagrangian
- ↪ subgroups ( $U_R(3) \times U_L(3) = U_V(3) \times U_A(3)$ ):
  - $U_V(1)$ : baryon number conservation
  - $SU_V(3)$ : flavor multiplets
  - $SU_A(3)$ : spontaneously broken
- ↪ 8 Goldstone bosons in vacuum and chiral restoration at some finite temperature/density
- $U_A(1)$ : would also be spontaneously broken (chiral condensate not invariant w.r.t.  $U_A(1)$ )



# Introduction

- neglecting quark masses smaller than  $\Lambda_{\text{QCD}}$
- ↪ chiral  $U_R(3) \times U_L(3)$  symmetry of QCD Lagrangian
- ↪ subgroups ( $U_R(3) \times U_L(3) = U_V(3) \times U_A(3)$ ):
  - $U_V(1)$ : baryon number conservation
  - $SU_V(3)$ : flavor multiplets
  - $SU_A(3)$ : spontaneously broken
- ↪ 8 Goldstone bosons in vacuum and chiral restoration at some finite temperature/density
- $U_A(1)$ : would also be spontaneously broken (chiral condensate not invariant w.r.t.  $U_A(1)$ )
- ↪ would lead to 9th Goldstone boson (flavor singlet with quantum numbers of  $\eta, \eta'$ )

# No symmetry of *Quantum* Chromodynamics

- chiral anomaly:  $U_A(1)$  is not symmetry of quantized theory



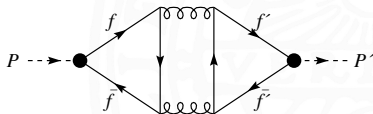
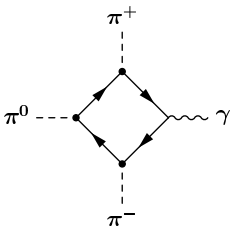
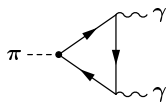
# No symmetry of *Quantum Chromodynamics*

- chiral anomaly:  $U_A(1)$  is **not** symmetry of **quantized** theory
  - NB: actually one can decide whether to break Lorentz invariance or  $U_A(1)$  by quantization
- ↪ path integral measure  $\mathcal{D}\bar{\psi} \mathcal{D}\psi$  breaks  $U_A(1)$ ,  
 $\mathcal{D}\psi^\dagger \mathcal{D}\psi$  breaks Lorentz invariance

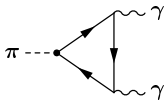


# No symmetry of *Quantum Chromodynamics*

- chiral anomaly:  $U_A(1)$  is **not** symmetry of **quantized** theory
- NB: actually one can decide whether to break Lorentz invariance or  $U_A(1)$  by quantization
- ↪ path integral measure  $\mathcal{D}\bar{\psi} \mathcal{D}\psi$  breaks  $U_A(1)$ ,  $\mathcal{D}\psi^\dagger \mathcal{D}\psi$  breaks Lorentz invariance
- ↪ experimental fact: nature seems to prefer Lorentz invariance
- ↪ consequences ...



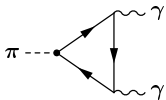
# Consequences of $U_A(1)$ anomaly in vacuum



- parameter-free prediction of  $\pi^0 \rightarrow 2\gamma$  in terms of pion decay constant  $f_\pi$

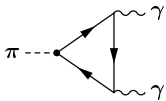


# Consequences of $U_A(1)$ anomaly in vacuum



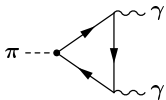
- parameter-free prediction of  $\pi^0 \rightarrow 2\gamma$  in terms of pion decay constant  $f_\pi$
- agreement with experiment within error bars  $< 1\%$  (PrimEx)

# Consequences of $U_A(1)$ anomaly in vacuum



- parameter-free prediction of  $\pi^0 \rightarrow 2\gamma$  in terms of pion decay constant  $f_\pi$
- agreement with experiment within error bars  $< 1\%$  (PrimEx)
- prediction strictly valid in the chiral limit ;-)

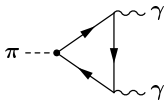
# Consequences of $U_A(1)$ anomaly in vacuum



- parameter-free prediction of  $\pi^0 \rightarrow 2\gamma$  in terms of pion decay constant  $f_\pi$
- agreement with experiment within error bars  $< 1\%$  (PrimEx)
- prediction strictly valid in the chiral limit ;-)
- ↪ coefficient of  $\frac{1}{f_\pi} \varepsilon^{\mu\nu\alpha\beta} \pi^0 F_{\mu\nu} F_{\alpha\beta}$  fixed by anomaly
- ↪ corrections suppressed  $\sim m_\pi^2 \varepsilon^{\mu\nu\alpha\beta} \pi^0 F_{\mu\nu} F_{\alpha\beta}$

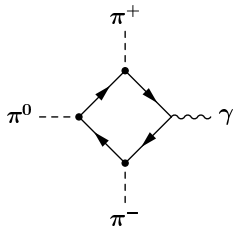


# Consequences of $U_A(1)$ anomaly in vacuum



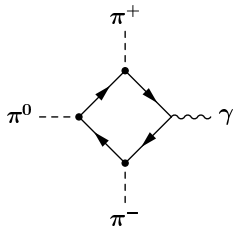
- parameter-free prediction of  $\pi^0 \rightarrow 2\gamma$  in terms of pion decay constant  $f_\pi$
  - agreement with experiment within error bars  $< 1\%$  (PrimEx)
  - prediction strictly valid in the chiral limit ;-)
  - ↪ coefficient of  $\frac{1}{f_\pi} \varepsilon^{\mu\nu\alpha\beta} \pi^0 F_{\mu\nu} F_{\alpha\beta}$  fixed by anomaly
  - ↪ corrections suppressed  $\sim m_\pi^2 \varepsilon^{\mu\nu\alpha\beta} \pi^0 F_{\mu\nu} F_{\alpha\beta}$
- in power counting of chiral perturbation theory:
- anomaly  $\sim O(q^4)$
  - otherwise  $\sim O(q^6)$
- with generic momentum  $q \sim m_\pi$

# Consequences of $U_A(1)$ anomaly in vacuum II



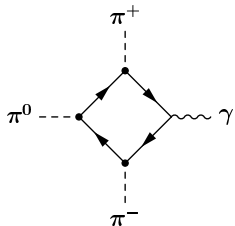
- coupling constant of  $\varepsilon^{\mu\nu\alpha\beta} \partial_\mu \pi^0 \partial_\nu \pi^+ \partial_\alpha \pi^- A_\beta \sim O(q^4)$   
fixed by anomaly

# Consequences of $U_A(1)$ anomaly in vacuum II



- coupling constant of  $\varepsilon^{\mu\nu\alpha\beta} \partial_\mu \pi^0 \partial_\nu \pi^+ \partial_\alpha \pi^- A_\beta \sim O(q^4)$   
fixed by anomaly
- corrections (=dominant in absence of anomaly)  $\sim O(q^6)$

# Consequences of $U_A(1)$ anomaly in vacuum II



- coupling constant of  $\varepsilon^{\mu\nu\alpha\beta} \partial_\mu \pi^0 \partial_\nu \pi^+ \partial_\alpha \pi^- A_\beta \sim O(q^4)$   
fixed by anomaly
- corrections (=dominant in absence of anomaly)  $\sim O(q^6)$
- ↪ predictive power for reactions  $e^+ e^- \rightarrow 3\pi$  and/or  $\gamma + \pi \rightarrow 2\pi$   
close to threshold?  $\rightsquigarrow$  spare slides

# Consequences of $U_A(1)$ anomaly in vacuum III

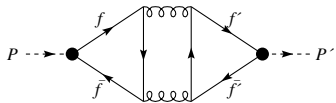


figure from Klabucar/Kekez/Scadron,  
hep-ph/0012267

- eta singlet is not a Goldstone boson

# Consequences of $U_A(1)$ anomaly in vacuum III

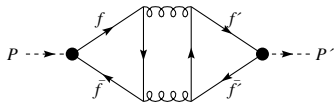


figure from Klabucar/Kekez/Scadron,  
hep-ph/0012267

- eta singlet is not a Goldstone boson
- Veneziano formula for mass of eta singlet:

$$m_{\eta_1}^2 = \frac{2N_f \tau}{f_\pi^2} \sim \frac{1}{N_c}$$

with topological susceptibility  $\tau = -i \int d^4x \langle 0 | \omega(x) \omega(0) | 0 \rangle$

and  $\omega \sim G_{\mu\nu}^a \tilde{G}_a^{\mu\nu}$

# Consequences of $U_A(1)$ anomaly in vacuum III

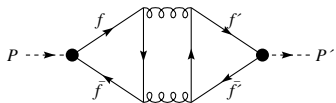


figure from Klabucar/Kekez/Scadron,  
hep-ph/0012267

- eta singlet is not a Goldstone boson
- Veneziano formula for mass of eta singlet:

$$m_{\eta_1}^2 = \frac{2N_f \tau}{f_\pi^2} \sim \frac{1}{N_c}$$

with topological susceptibility  $\tau = -i \int d^4x \langle 0 | \omega(x) \omega(0) | 0 \rangle$

and  $\omega \sim G_{\mu\nu}^a \tilde{G}_a^{\mu\nu}$

- note: nine light pseudoscalars in the combined chiral and large- $N_c$  limit ( $N_c =$  number of colors)

# Consequences of $U_A(1)$ anomaly in vacuum III

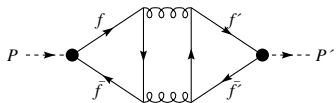


figure from Klabucar/Kekez/Scadron,  
hep-ph/0012267

- eta singlet is not a Goldstone boson
- Veneziano formula for mass of eta singlet:

$$m_{\eta_1}^2 = \frac{2N_f \tau}{f_\pi^2} \sim \frac{1}{N_c}$$

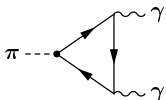
with topological susceptibility  $\tau = -i \int d^4x \langle 0 | \omega(x) \omega(0) | 0 \rangle$   
and  $\omega \sim G_{\mu\nu}^a \tilde{G}_a^{\mu\nu}$

- note: nine light pseudoscalars in the combined chiral and large- $N_c$  limit ( $N_c =$  number of colors)
- ↪ starting point of large- $N_c$  chiral perturbation theory ( $\chi$ PT)  
(Kaiser/Leutwyler, EPJ C 17, 623 (2000))



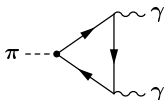
# What changes in a medium?

- $\pi$ - $\gamma$ - $\gamma$  coupling in vacuum  $\sim 1/f_\pi$



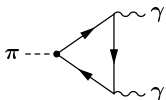
# What changes in a medium?

- $\pi$ - $\gamma$ - $\gamma$  coupling in vacuum  $\sim 1/f_\pi$
- ↪  $f_\pi$  is order parameter of  $SU(N_f)$  chiral symmetry breaking



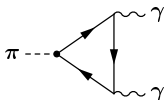
# What changes in a medium?

- $\pi$ - $\gamma$ - $\gamma$  coupling in vacuum  $\sim 1/f_\pi$
- $\hookrightarrow$   $f_\pi$  is order parameter of  $SU(N_f)$  chiral symmetry breaking
- $\rightsquigarrow$  enhanced decay  $\pi^0 \rightarrow 2\gamma$ ?



# What changes in a medium?

- $\pi$ - $\gamma$ - $\gamma$  coupling in vacuum  $\sim 1/f_\pi$
- ↪  $f_\pi$  is order parameter of  $SU(N_f)$  chiral symmetry breaking
- ↪ enhanced decay  $\pi^0 \rightarrow 2\gamma$ ?
- ↪ no! instead: decay decouples from anomaly, suppressed decay due to  $SU(N_f)$  chiral restoration (Pisarski/Trueman/Tytgat, PRD 56, 7077 (1997))



# Chiral multiplets

at chiral restoration of  $SU(3)$  (no statement about  $U_A(1)$  needed!):

- chiral multiplets  $(L, R)$  instead of just flavor multiplets

thanks to D. Jido for pointing this out to me, Phys.Rev.C85 (2012) 032201

# Chiral multiplets

at chiral restoration of  $SU(3)$  (no statement about  $U_A(1)$  needed!):

- chiral multiplets  $(L, R)$  instead of just flavor multiplets
- in particular look at spin 0:  $\bar{q}_{La} q_{Rb}$  (with flavor indices  $a, b$ )

thanks to D. Jido for pointing this out to me, Phys.Rev.C85 (2012) 032201

# Chiral multiplets

at chiral restoration of  $SU(3)$  (no statement about  $U_A(1)$  needed!):

- chiral multiplets  $(L, R)$  instead of just flavor multiplets
- in particular look at spin 0:  $q_{L_a} \bar{q}_{R_b}$  (with flavor indices  $a, b$ )

$$\rightarrow (\bar{3}, 1) \times (1, 3) = (\bar{3}, 3) \rightsquigarrow \text{nonet}$$

thanks to D. Jido for pointing this out to me, Phys.Rev.C85 (2012) 032201

# Chiral multiplets

at chiral restoration of SU(3) (no statement about  $U_A(1)$  needed!):

- chiral multiplets  $(L, R)$  instead of just flavor multiplets
- in particular look at spin 0:  $q_{La} q_{Rb}$  (with flavor indices  $a, b$ )
- ↪  $(\bar{3}, 1) \times (1, 3) = (\bar{3}, 3) \rightsquigarrow$  nonet
- parity commutes with QCD Hamiltonian and changes  $(\bar{3}, 3)$  to  $(3, \bar{3})$
- ↪ at and above transition: 18plet of degenerate states, with quantum numbers of  $\pi, K, \eta, \eta', a_0, \kappa, 2 f_0$ 's

thanks to D. Jido for pointing this out to me, Phys.Rev.C85 (2012) 032201



# Chiral multiplets

at chiral restoration of SU(3) (no statement about  $U_A(1)$  needed!):

- chiral multiplets  $(L, R)$  instead of just flavor multiplets
- in particular look at spin 0:  $\bar{q}_{La} q_{Rb}$  (with flavor indices  $a, b$ )
- $(\bar{3}, 1) \times (1, 3) = (\bar{3}, 3) \rightsquigarrow$  nonet
- parity commutes with QCD Hamiltonian and changes  $(\bar{3}, 3)$  to  $(3, \bar{3})$
- ↪ at and above transition: 18plet of degenerate states, with quantum numbers of  $\pi, K, \eta, \eta', a_0, \kappa, 2 f_0$ 's
- at and below transition: **Goldstone bosons** stay light

thanks to D. Jido for pointing this out to me, Phys.Rev.C85 (2012) 032201

# Chiral multiplets

at chiral restoration of SU(3) (no statement about  $U_A(1)$  needed!):

- chiral multiplets  $(L, R)$  instead of just flavor multiplets
- in particular look at spin 0:  $q_{La} q_{Rb}$  (with flavor indices  $a, b$ )
- ↪  $(\bar{3}, 1) \times (1, 3) = (\bar{3}, 3) \rightsquigarrow$  nonet
- parity commutes with QCD Hamiltonian and changes  $(\bar{3}, 3)$  to  $(3, \bar{3})$
- ↪ at and above transition: 18plet of degenerate states, with quantum numbers of  $\pi, K, \eta, \eta', a_0, \kappa, 2 f_0$ 's
- at and below transition: **Goldstone bosons** stay light
- ↪  $\eta'$  should become light at transition

thanks to D. Jido for pointing this out to me, Phys.Rev.C85 (2012) 032201

## A light $\eta'$ meson?

- might imply that  $\eta'$  becomes 9th Goldstone boson, effective restoration of  $U_A(1)$ ?



## A light $\eta'$ meson?

- might imply that  $\eta'$  becomes 9th Goldstone boson, effective restoration of  $U_A(1)$ ?
- or might imply that  $\eta'$  decouples from anomaly, i.e. change of decay constant of  $\eta'$



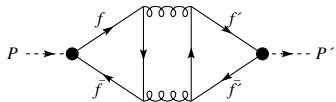
## A light $\eta'$ meson?

- might imply that  $\eta'$  becomes 9th Goldstone boson, effective restoration of  $U_A(1)$ ?
- or might imply that  $\eta'$  decouples from anomaly, i.e. change of decay constant of  $\eta'$
- ↪ would modify decays  $\eta' \rightarrow \gamma\gamma, \gamma\pi^+\pi^-$



# A light $\eta'$ meson?

- might imply that  $\eta'$  becomes 9th Goldstone boson, effective restoration of  $U_A(1)$
  - or might imply that  $\eta'$  decouples from anomaly, i.e. change of decay constant of  $\eta'$
  - ↪ would modify decays  $\eta' \rightarrow \gamma\gamma, \gamma\pi^+\pi^-$
- 
- caveats: only good argument if not strong first-order transition and **if states survive as quasi-particles**



# Does $\eta'$ remain as a quasi-particle?

task: determine in-medium width of  $\eta'$  meson



# Does $\eta'$ remain as a quasi-particle?

task: determine in-medium width of  $\eta'$  meson

- experiment: ask, e.g., Volker Metag





# Does $\eta'$ remain as a quasi-particle?

task: determine in-medium width of  $\eta'$  meson

- experiment: ask, e.g., Volker Metag
- our theory approach: try to be as model independent as possible



# Does $\eta'$ remain as a quasi-particle?

task: determine in-medium width of  $\eta'$  meson

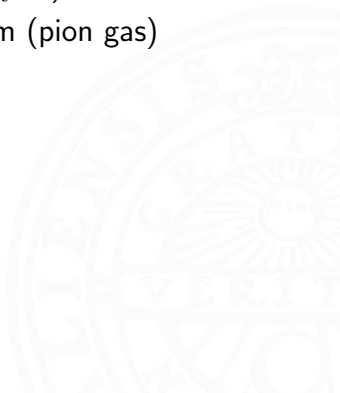
- experiment: ask, e.g., Volker Metag
  - our theory approach: try to be as model independent as possible
- ↪ use large- $N_c$  chiral perturbation theory ( $\chi$ PT) + resonances



# Does $\eta'$ remain as a quasi-particle?

task: determine in-medium width of  $\eta'$  meson

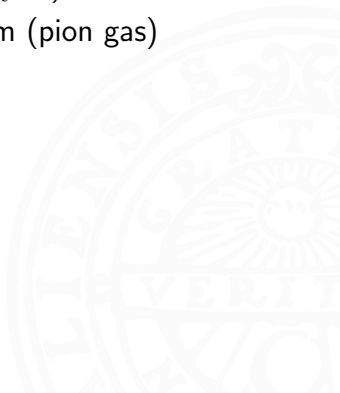
- experiment: ask, e.g., Volker Metag
- our theory approach: try to be as model independent as possible
- ↪ use large- $N_c$  chiral perturbation theory ( $\chi$ PT) + resonances
- start with simpler case of thermal medium (pion gas)
  - ↪ in this talk



# Does $\eta'$ remain as a quasi-particle?

task: determine in-medium width of  $\eta'$  meson

- experiment: ask, e.g., Volker Metag
- our theory approach: try to be as model independent as possible
- ↪ use large- $N_c$  chiral perturbation theory ( $\chi$ PT) + resonances
  - start with simpler case of thermal medium (pion gas)
    - ↪ in this talk
  - outlook: consequences for finite density



# Why is this useful?

- Why a **low-temperature** calculation?
- Aren't we interested in **extreme** conditions?



# Why is this useful?

- Why a **low-temperature** calculation?
  - Aren't we interested in **extreme** conditions?
- ↪ chiral perturbation theory allows for **systematic**, model independent calculations at low temperatures



# Why is this useful?

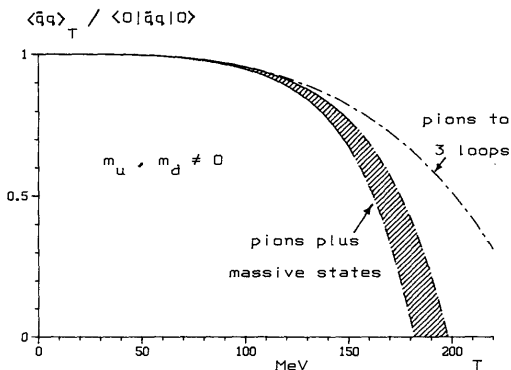
- Why a **low-temperature** calculation?
- Aren't we interested in **extreme** conditions?
- ↪ chiral perturbation theory allows for **systematic**, model independent calculations at low temperatures
- provides excellent description of onset of chiral restoration for chiral condensate (and for pion decay constant)  $\rightsquigarrow$  next slide

# Why is this useful?

- Why a **low-temperature** calculation?
- Aren't we interested in **extreme** conditions?
- ↪ chiral perturbation theory allows for **systematic**, model independent calculations at low temperatures
- provides excellent description of onset of chiral restoration for chiral condensate (and for pion decay constant)  $\rightsquigarrow$  next slide
- ↪ chiral perturbation theory has something to say about not too low temperatures



# Drop of quark condensate



- uses interacting pions and resonance gas
- produces realistic transition temperature

Gerber/Leutwyler, Nucl.Phys.B 321, 387 (1989)

# Role of resonances

- chiral perturbation theory is model independent (contains only Goldstone bosons)



# Role of resonances

- chiral perturbation theory is model independent (contains only Goldstone bosons)
- but resonances are also important



# Role of resonances

- chiral perturbation theory is model independent (contains only Goldstone bosons)
- but resonances are also important
- in particular:
  - large number of colors  $N_c$ :

$$\frac{1}{\sqrt{N_c}} \sim m_{\eta'} \ll m_R \sim N_c^0$$



# Role of resonances

- chiral perturbation theory is model independent (contains only Goldstone bosons)
- but resonances are also important
- in particular:
  - large number of colors  $N_c$ :

$$\frac{1}{\sqrt{N_c}} \sim m_{\eta'} \ll m_R \sim N_c^0$$

- while for  $N_c = 3$ :

$$m_{\eta'} \stackrel{!}{\approx} m_R$$

for resonances  $R = a_0, f_0, \kappa, K^*$

# Role of resonances

- chiral perturbation theory is model independent (contains only Goldstone bosons)
- but resonances are also important
- in particular:
  - large number of colors  $N_c$ :

$$\frac{1}{\sqrt{N_c}} \sim m_{\eta'} \ll m_R \sim N_c^0$$

- while for  $N_c = 3$ :

$$m_{\eta'} \stackrel{!}{\approx} m_R$$

for resonances  $R = a_0, f_0, \kappa, K^*$

- ↪ include resonances such that formal low-energy and large- $N_c$  limit fits to chiral perturbation theory

# Role of resonances

- chiral perturbation theory is model independent (contains only Goldstone bosons)
- but resonances are also important
- in particular:
  - large number of colors  $N_c$ :

$$\frac{1}{\sqrt{N_c}} \sim m_{\eta'} \ll m_R \sim N_c^0$$

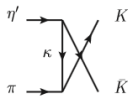
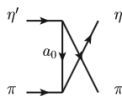
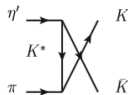
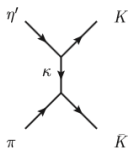
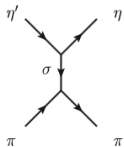
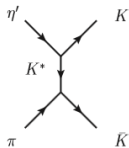
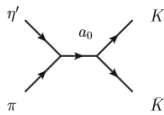
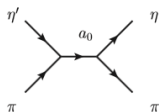
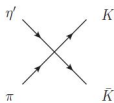
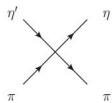
- while for  $N_c = 3$ :

$$m_{\eta'} \overset{!}{\approx} m_R$$

for resonances  $R = a_0, f_0, \kappa, K^*$

- ↪ include resonances such that formal low-energy and large- $N_c$  limit fits to chiral perturbation theory
- ↪ inclusion of resonances as model independent as possible

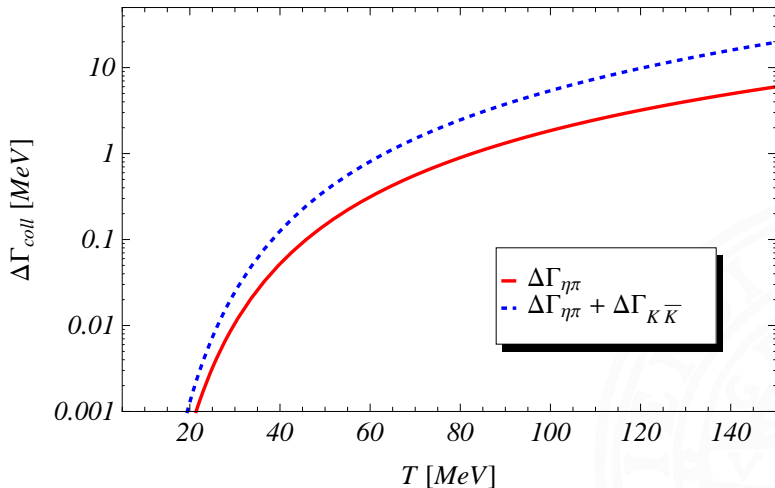
# Processes/diagrams for $\pi\eta'$ scattering



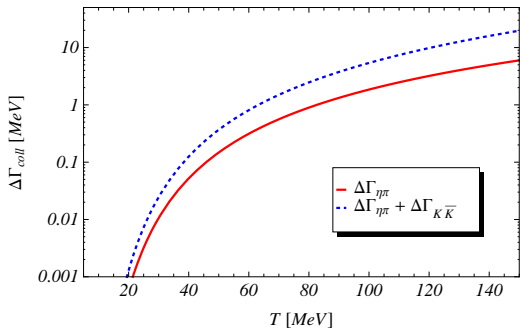
(note: loops are suppressed in the large- $N_c$  limit)



# Results: collisional width of $\eta'$ meson

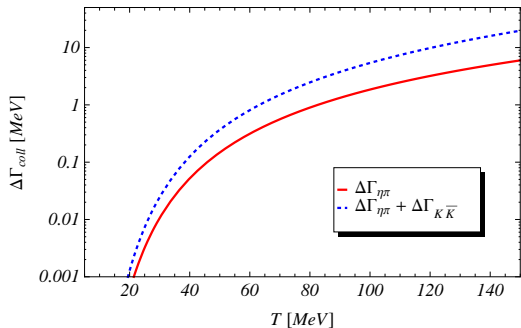


# Results: collisional width of $\eta'$ meson



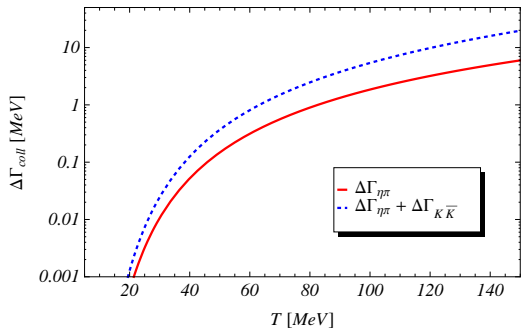
- $\eta'$  remains a rather narrow state

# Results: collisional width of $\eta'$ meson



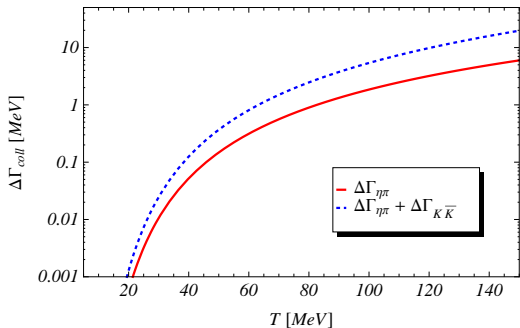
- $\eta'$  remains a rather narrow state
- width comparable to life time of fireball in heavy-ion collision

# Results: collisional width of $\eta'$ meson



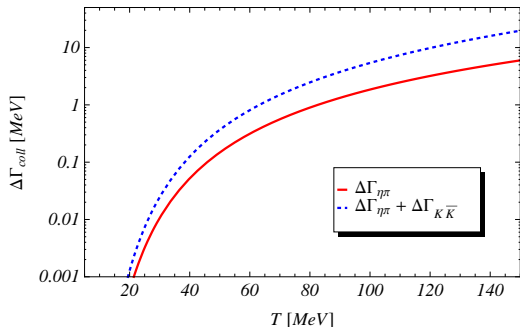
- $\eta'$  remains a rather narrow state
- width comparable to life time of fireball in heavy-ion collision
- kaons are important

# Results: collisional width of $\eta'$ meson



- $\eta'$  remains a rather narrow state
- width comparable to life time of fireball in heavy-ion collision
- kaons are important
- how to measure this?

# Results: collisional width of $\eta'$ meson



- $\eta'$  remains a rather narrow state
  - width comparable to life time of fireball in heavy-ion collision
  - kaons are important
  - how to measure this?
- ↪ maybe via anomaly driven processes  $\eta' \rightarrow \gamma\gamma, \gamma\pi^+\pi^-$

# Summary

- ongoing: experimental checks of consequences of chiral anomaly in **vacuum** (with support from theory)



# Summary

- ongoing: experimental checks of consequences of chiral anomaly in **vacuum** (with support from theory)
- **thermal** modifications? (chiral symmetry, deconfinement, fate of anomaly, ...)





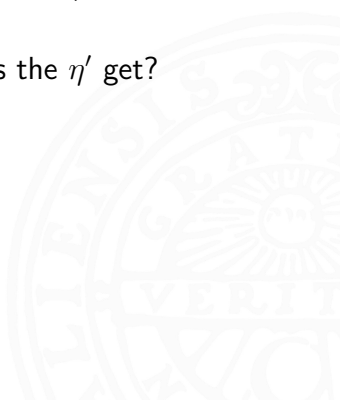
# Summary

- ongoing: experimental checks of consequences of chiral anomaly in **vacuum** (with support from theory)
- **thermal** modifications? (chiral symmetry, deconfinement, fate of anomaly, . . .)
- suggestive to have a downward mass shift for  $\eta'$  (18plet of states at chiral restoration)



# Summary

- ongoing: experimental checks of consequences of chiral anomaly in **vacuum** (with support from theory)
- **thermal** modifications? (chiral symmetry, deconfinement, fate of anomaly, ...)
- suggestive to have a downward mass shift for  $\eta'$  (18plet of states at chiral restoration)
- complementary question: how broad does the  $\eta'$  get?



# Summary

- ongoing: experimental checks of consequences of chiral anomaly in **vacuum** (with support from theory)
- **thermal** modifications? (chiral symmetry, deconfinement, fate of anomaly, ...)
- suggestive to have a downward mass shift for  $\eta'$  (18plet of states at chiral restoration)
- complementary question: how broad does the  $\eta'$  get?
- quantitative answer from large- $N_c$  chiral perturbation theory + resonances

# Summary

- ongoing: experimental checks of consequences of chiral anomaly in **vacuum** (with support from theory)
  - **thermal** modifications? (chiral symmetry, deconfinement, fate of anomaly, ...)
  - suggestive to have a downward mass shift for  $\eta'$  (18plet of states at chiral restoration)
  - complementary question: how broad does the  $\eta'$  get?
  - quantitative answer from large- $N_c$  chiral perturbation theory + resonances
- ↪  $\eta'$  survives as a quasi-particle

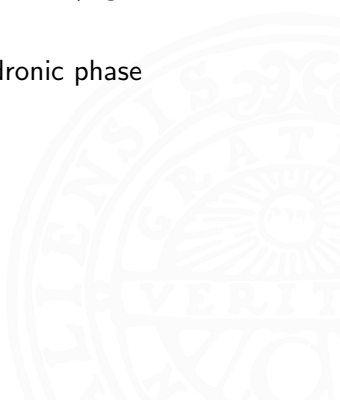
# Outlook

- **density** modifications? (chiral symmetry, deconfinement, fate of anomaly, ...)
- suggestive to have a downward mass shift for  $\eta'$  (18plet of states at chiral restoration)
- complementary question: how broad does the  $\eta'$  get?



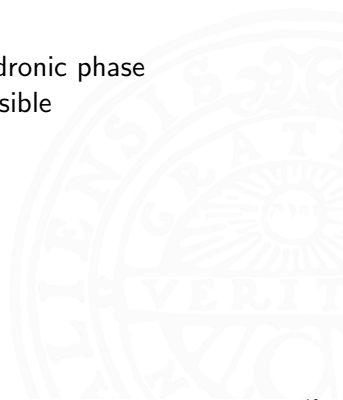
# Outlook

- density modifications? (chiral symmetry, deconfinement, fate of anomaly, ...)
- suggestive to have a downward mass shift for  $\eta'$  (18plet of states at chiral restoration)
- complementary question: how broad does the  $\eta'$  get?
- translation from thermal calculations:
  - use hadronic degrees of freedom in hadronic phase



# Outlook

- density modifications? (chiral symmetry, deconfinement, fate of anomaly, ...)
- suggestive to have a downward mass shift for  $\eta'$  (18plet of states at chiral restoration)
- complementary question: how broad does the  $\eta'$  get?
- translation from thermal calculations:
  - use hadronic degrees of freedom in hadronic phase
  - try to be as model independent as possible



# Outlook

- density modifications? (chiral symmetry, deconfinement, fate of anomaly, ...)
- suggestive to have a downward mass shift for  $\eta'$  (18plet of states at chiral restoration)
- complementary question: how broad does the  $\eta'$  get?
- translation from thermal calculations:
  - use hadronic degrees of freedom in hadronic phase
  - try to be as model independent as possible
  - let pions come from pion cloud of nucleon  $\rightsquigarrow$  figures

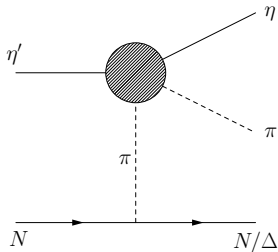
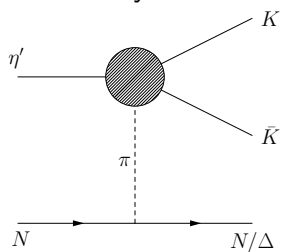


# Outlook

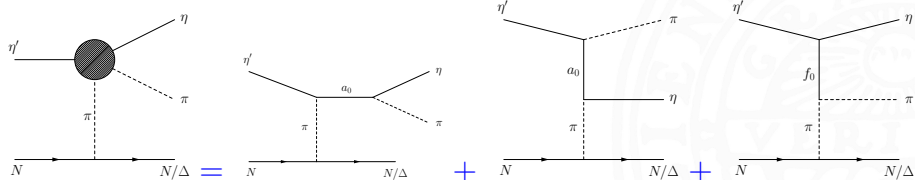
- density modifications? (chiral symmetry, deconfinement, fate of anomaly, ...)
- suggestive to have a downward mass shift for  $\eta'$  (18plet of states at chiral restoration)
- complementary question: how broad does the  $\eta'$  get?
- translation from thermal calculations:
  - use hadronic degrees of freedom in hadronic phase
  - try to be as model independent as possible
  - let pions come from pion cloud of nucleon  $\rightsquigarrow$  figures
- ↪ suggests that three-body final states are important!  
(while most calculations implicitly assume dominance of two-body final states)
- ↪ still work left to do for theory

# Important processes at finite density

three-body final states!



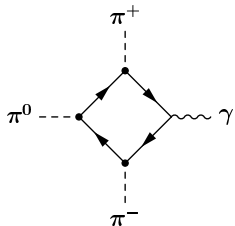
in more detail:



# Spare slides

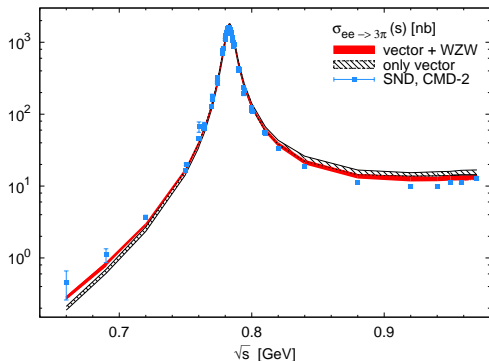


# Consequences of $U_A(1)$ anomaly in vacuum II



- coupling constant of  $\varepsilon^{\mu\nu\alpha\beta} \partial_\mu \pi^0 \partial_\nu \pi^+ \partial_\alpha \pi^- A_\beta \sim O(q^4)$   
fixed by anomaly
- corrections (=dominant in absence of anomaly)  $\sim O(q^6)$
- ↪ predictive power for reactions  $e^+ e^- \rightarrow 3\pi$  and/or  $\gamma + \pi \rightarrow 2\pi$   
close to threshold?
- ↪ How far is threshold away from idealized case of chiral limit?

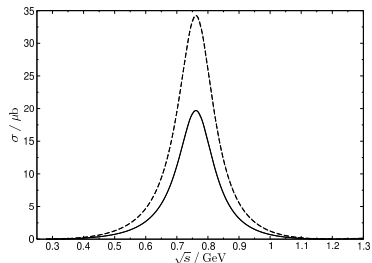
$$e^+ e^- \rightarrow 3\pi$$



C. Terschläsen, B. Strandberg, SL, F. Eichstädt,  
Eur.Phys.J.A 49 (2013) 116

- anomaly caused by Wess-Zumino-Witten (WZW) term
- data dominated by  $\omega$  vector meson peak
- ↪ anomaly has some effect, but does not dominate a region
- ↪ threshold at  $3 m_\pi$  already quite sizable  $\neq$  chiral limit

$$\gamma + \pi \rightarrow 2\pi$$

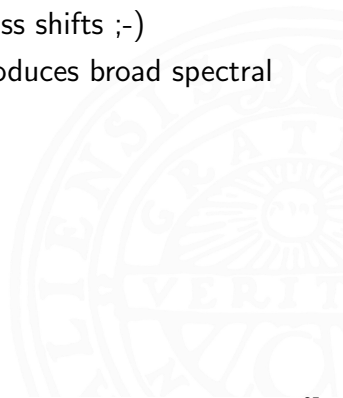


M. Hoferichter, B. Kubis, D. Sakas,  
Phys.Rev.D86, 116009 (2012)

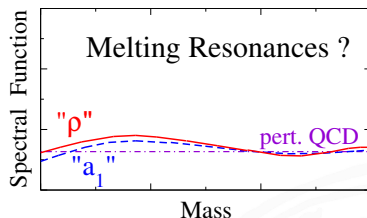
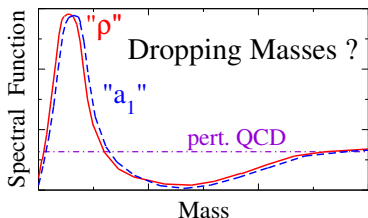
- expect future data from COMPASS@CERN (pion beam, Primakov effect)
  - calculation includes anomaly and pion-pion rescattering using dispersion theory
  - **solid line:** prediction from anomaly
  - **dashed line:** size of anomaly scaled up by about 30%
- ↪ can use whole range to pin down anomaly, not just threshold region

# Mass shift?

- systematic calculation of thermal modifications of properties of  $\eta'$  using large- $N_c$  chiral perturbation + resonances is interesting
- ↪ mass shift?
- ↪ check first: does  $\eta'$  survive at all in a thermal medium?
- personal history in scepticism against mass shifts ;-)
- hadronic many-body framework often produces broad spectral functions instead of dropping masses
- deconfinement shines out chiral effects



# Mass shift vs. broadening

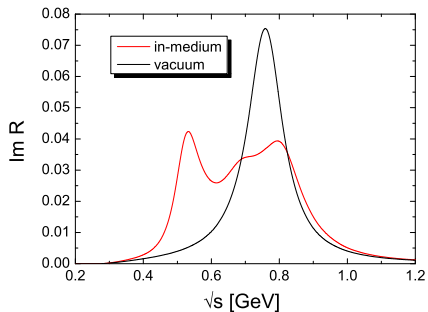


- chiral restoration demands degeneracy of spectra of chiral partners  
(btw: for spin 1: 16-plets, not 18-plets)
- melting of resonances might be hadronic precursor to deconfinement

figures from Rapp, Wambach, van Hees, arXiv:0901.3289 [hep-ph]



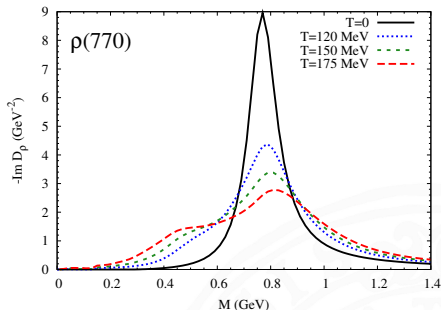
# Much celebrated example: $\rho$ meson



Post, SL, U. Mosel,

Nucl.Phys.A 741, 81 (2004)

M.



Rapp, J. Wambach, H. van Hees,

arXiv:0901.3289 [hep-ph]

R.

different groups (with different models) obtain  
broad  $\rho$  meson spectra, essentially no mass shift