

# Phi meson properties in nuclear matter

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Talk at the “APCTP Focus Program in Nuclear Physics  
2021 Part I: Hadron properties in a nuclear medium from  
the quark and gluon degrees of freedom”,  
online  
July 15, 2021

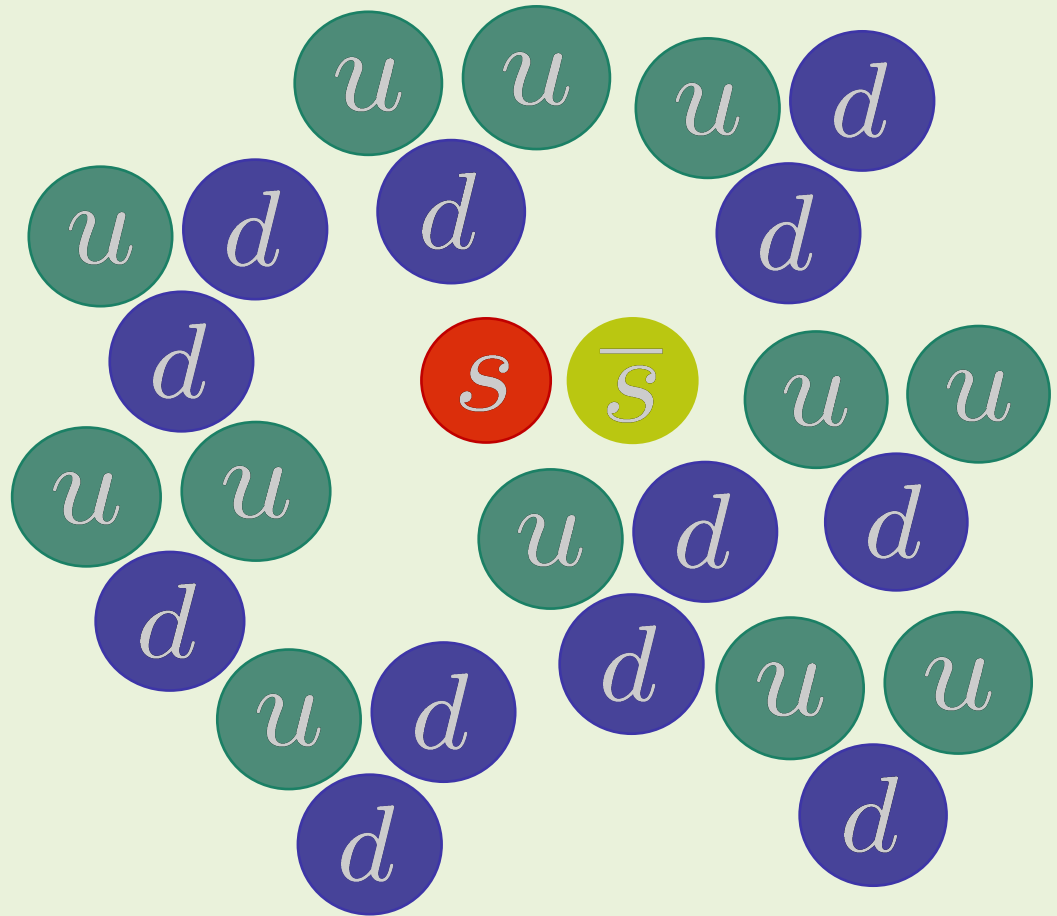
Work done in collaboration with  
Elena Bratkovskaya (Frankfurt U./GSI)  
HyungJoo Kim (APCTP/Yonsei U.)  
Keisuke Ohtani (Tokyo Tech)

$\phi$  meson



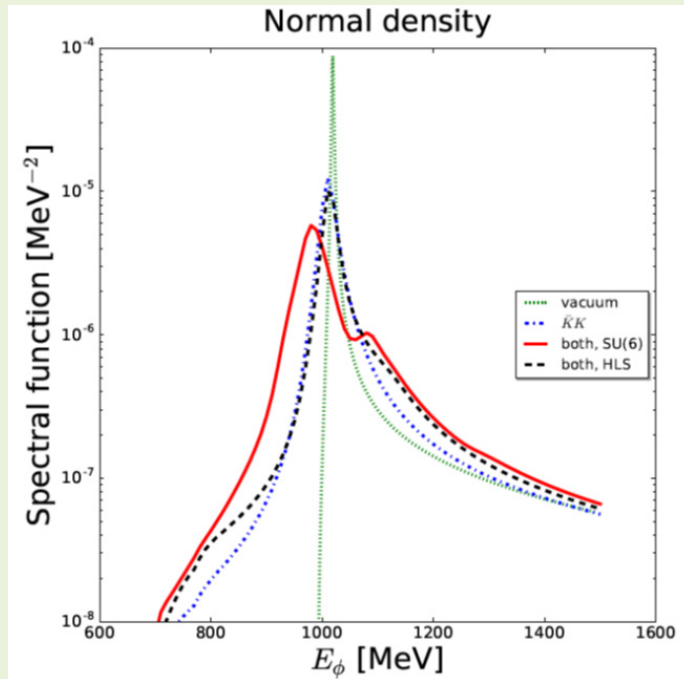
$$m_\phi = 1019 \text{ MeV}$$

$$\Gamma_\phi = 4.3 \text{ MeV}$$



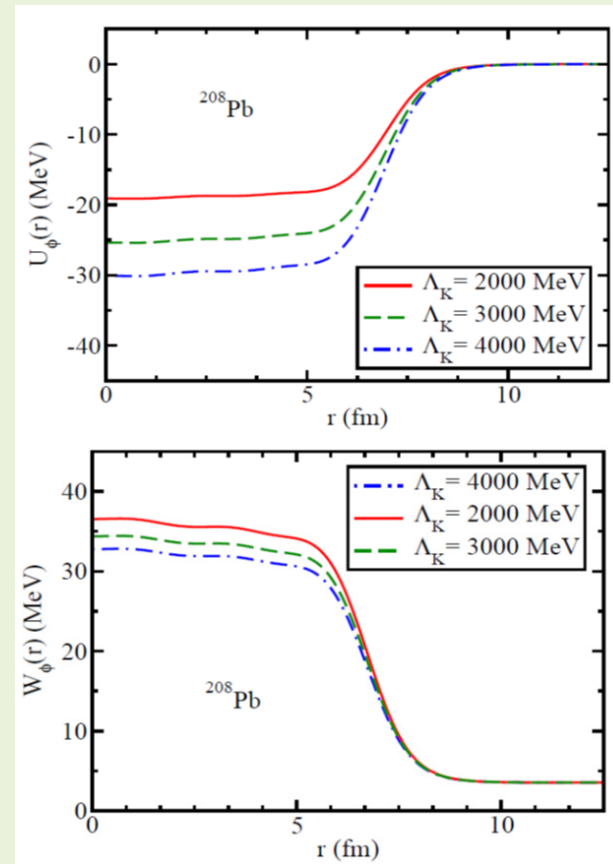
# Many theoretical works about the $\phi$ meson at finite density in recent years

## Spectral functions from hadronic models



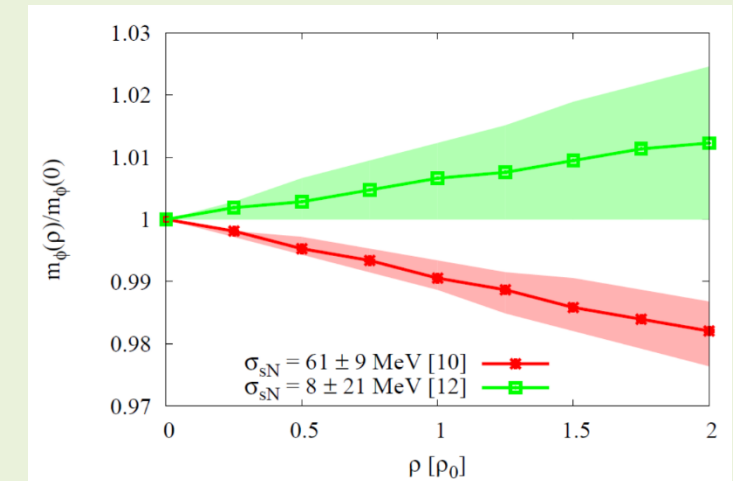
- P. Gubler and W. Weise, Phys. Lett. B **751**, 396 (2015).  
P. Gubler and W. Weise, Nucl. Phys. A **954**, 125 (2016).  
D. Cabrera *et al.*, Phys. Rev. C **95**, 015201 (2017).  
D. Cabrera *et al.*, Phys. Rev. C **96**, 034618 (2017).

## Possibility of $\phi$ -nucleus bound state



- J.J. Cobos-Martinez *et al.*, Phys. Lett. B **771**, 113 (2017).  
J.J. Cobos-Martinez *et al.*, Phys. Rev. C **96**, 035201 (2017).

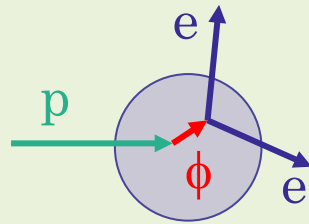
## Mass shift in nuclear matter from QCD sum rules



- P. Gubler and K. Ohtani, Phys. Rev. D **90**, 094002 (2014).  
H.J. Kim *et al.*, Phys. Lett. B **772**, 194 (2017).  
H.J. Kim and P. Gubler, Phys. Lett. B **805**, 135412 (2020).

# Previous experimental results

KEK  
E325



12 GeV  
pA-reaction

slow  $\varphi$ s

Pole mass:

$$\frac{m_\phi(\rho)}{m_\phi(0)} = 1 - k_1 \frac{\rho}{\rho_0}$$

$0.034 \pm 0.007$

intermediate  
 $\varphi$ s

Pole width:

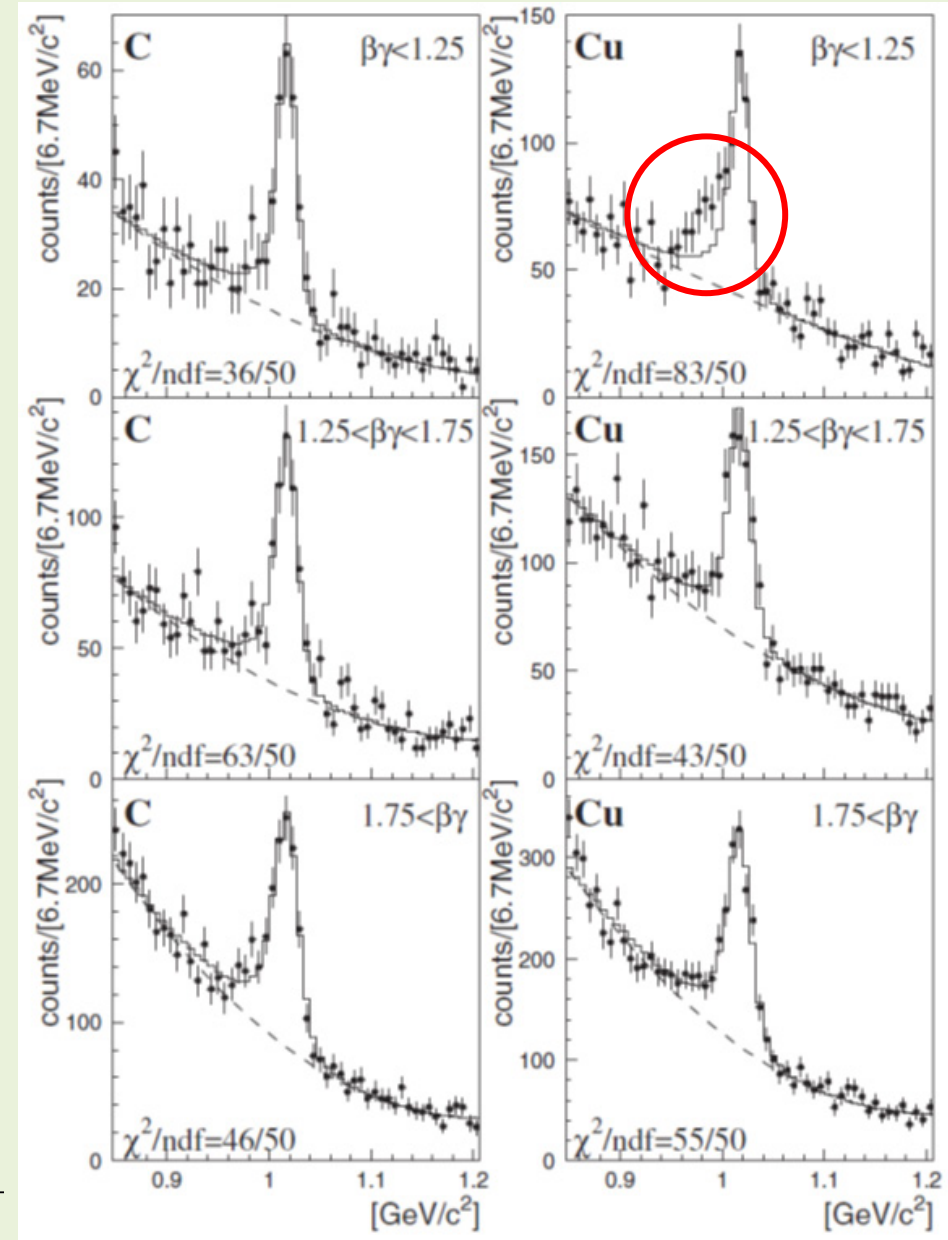
$$\frac{\Gamma_\phi(\rho)}{\Gamma_\phi(0)} = 1 + k_2 \frac{\rho}{\rho_0}$$

$2.6 \pm 1.5$

fast  $\varphi$ s

To be measured again at the J-PARC E16  
experiment with 100x increased statistics!


$$\beta\gamma = \frac{|\vec{p}|}{m_\phi}$$



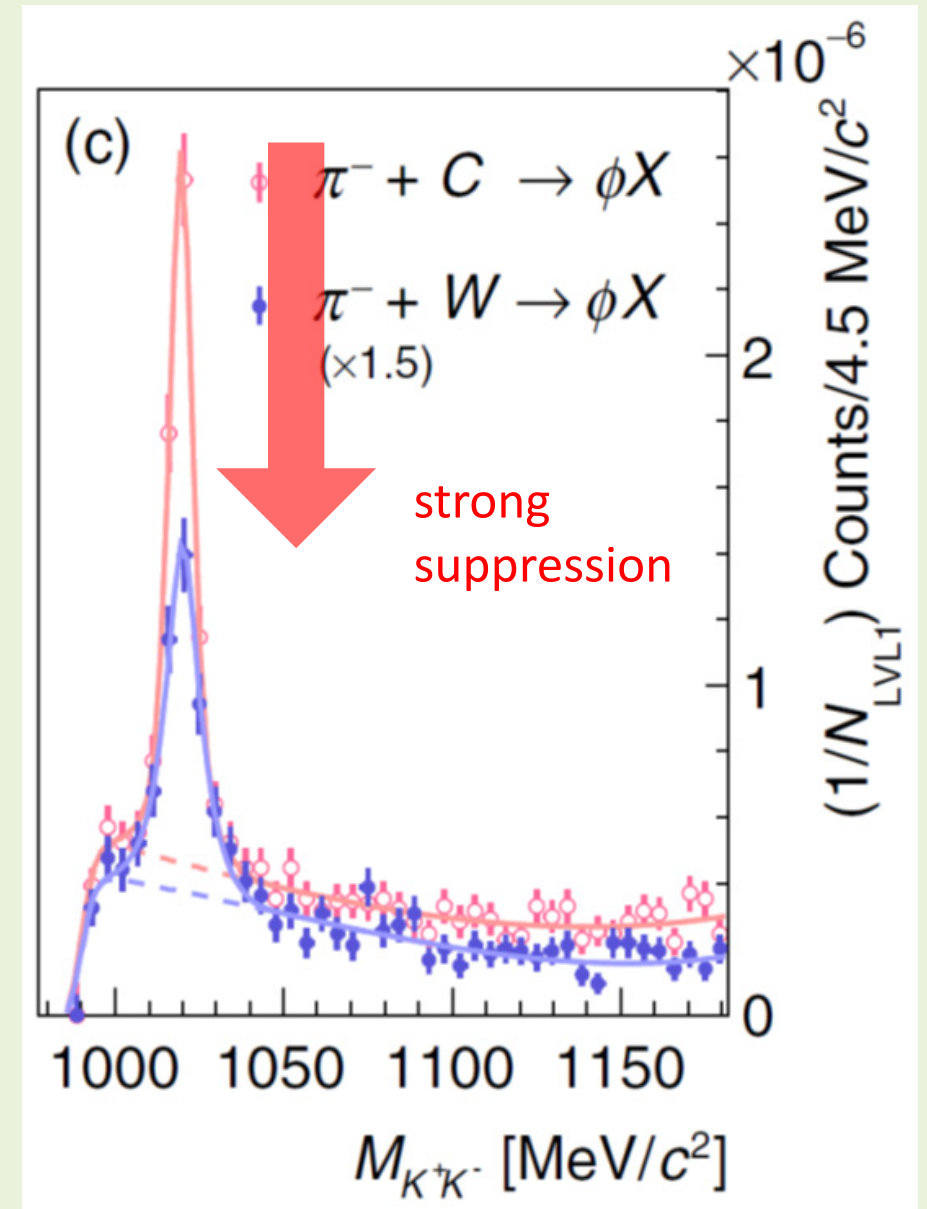


# Recent experimental results

HADES: 1.7 GeV  $\pi^-$ A-reaction

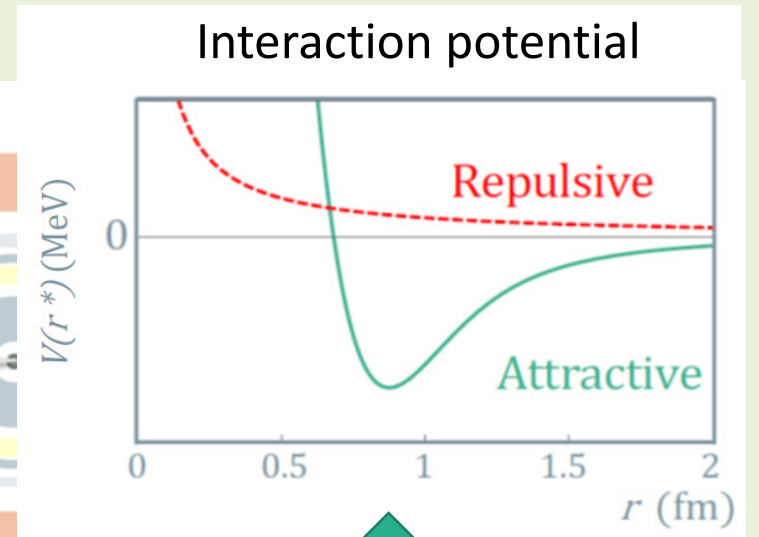
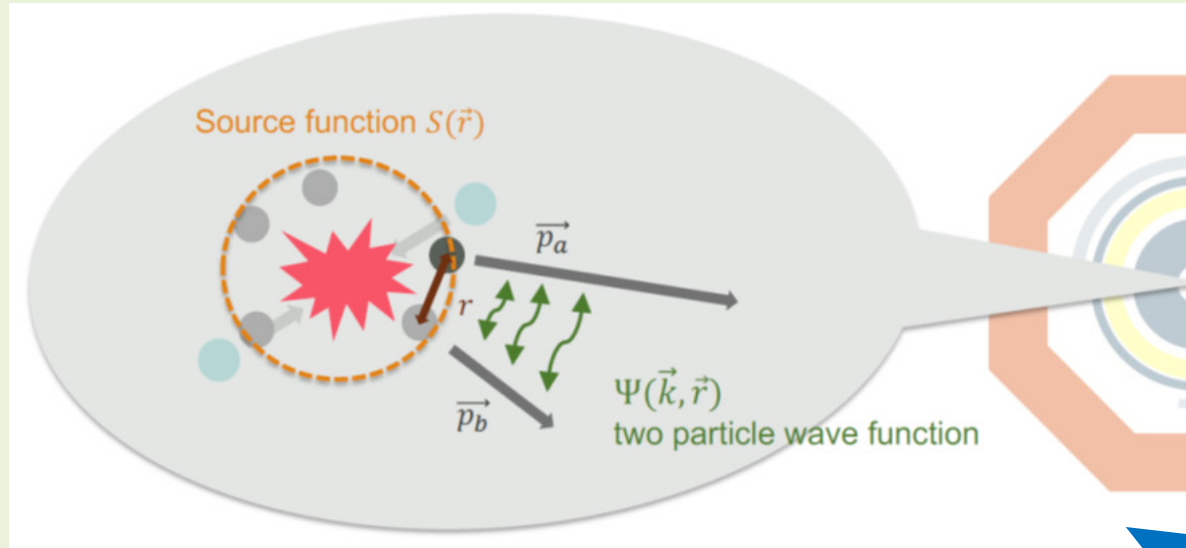
- ★ Larger suppression of  $K^-$  in the Tungsten target compared to the Carbon target
  - ★  $K^-/\phi$  ratio is similar for both Tungsten and Carbon targets
- 
- ★ Observation of large suppression (broadening?) of the  $\phi$  meson in large nuclei

$K^+K^-$  - invariant mass spectrum

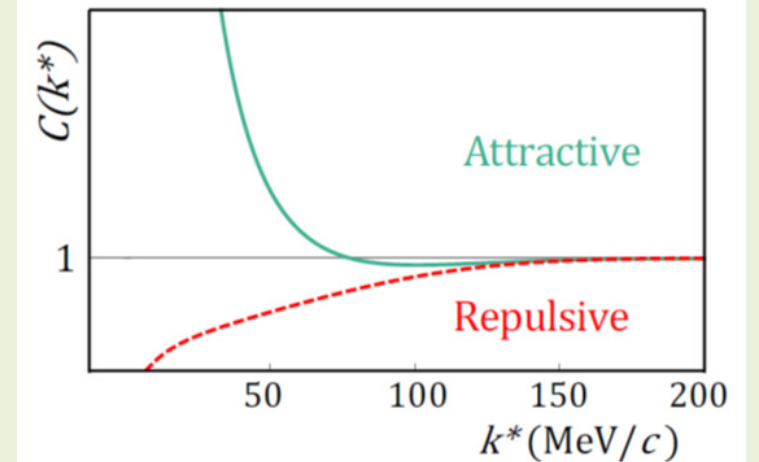


# New experimental results

## ALICE (Femtoscscopy)



Correlation function



The observable to be measured: the correlation function:

$$C(k) = \mathcal{N} \frac{N_{\text{Same}}}{N_{\text{Mixed}}} = \int S(\vec{r}) |\Psi(\vec{k}, \vec{r})|^2 d^3\vec{r}$$


Emission source  
(Gaussian)

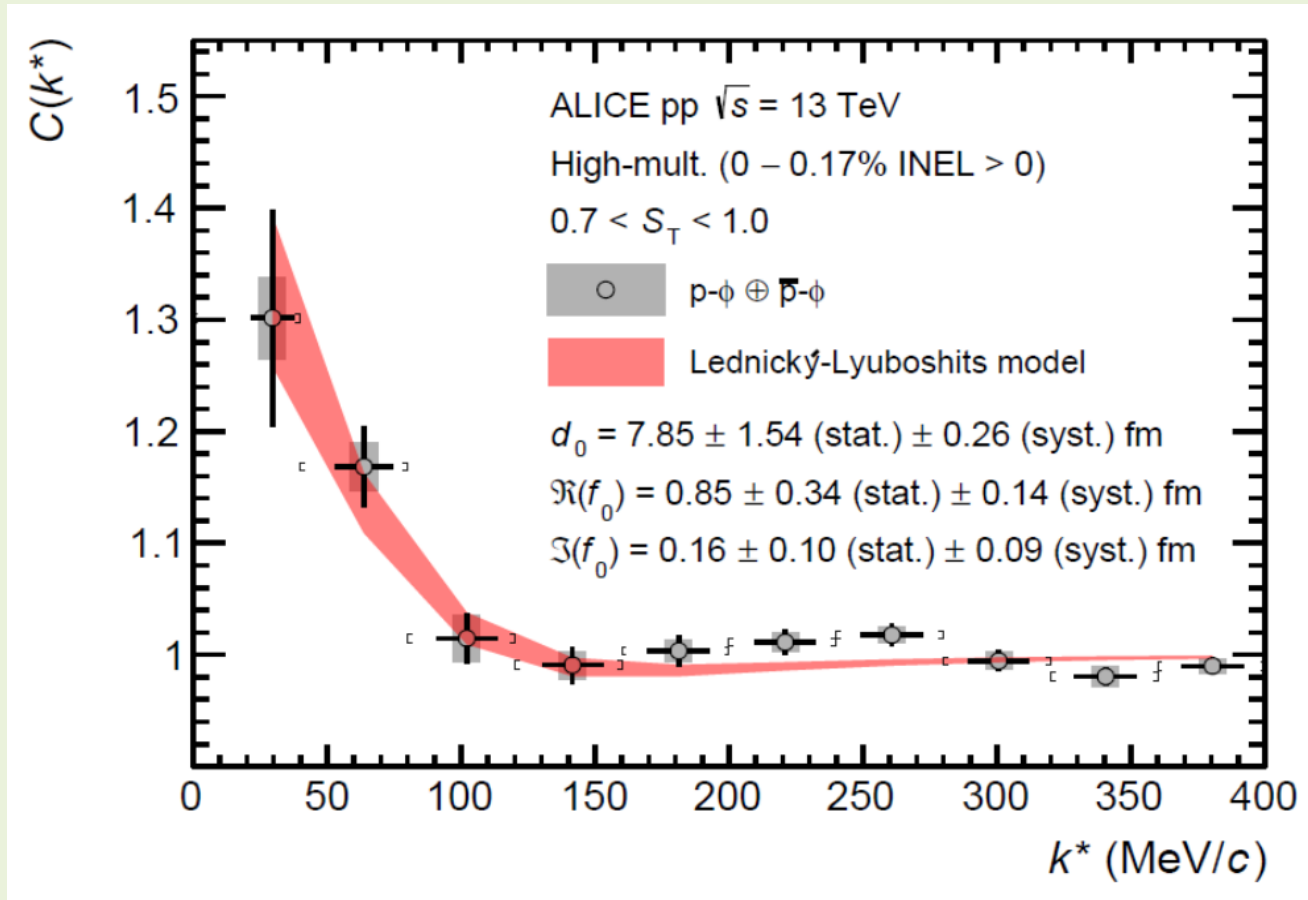
Relative momentum  
of the particle pair

# New experimental results

ALICE

Measurement of  $\varphi$ N the correlation function

Attraction!  




Information about the  $\varphi$ N scattering length

Real part:

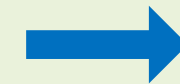
$$\Re(f_0) = 0.85 \pm 0.34(\text{stat.}) \pm 0.14(\text{syst.}) \text{ fm}$$



Attractive

Imaginary part:

$$\Im(f_0) = 0.16 \pm 0.10(\text{stat.}) \pm 0.09(\text{syst.}) \text{ fm}$$



Small absorption ?

# New experimental results

ALICE

Fit of the correlation function data to two simple phenomenological potentials

$$V_{\text{Yukawa}}(r) = -\frac{A}{r} e^{-\alpha r}$$

$$A = 0.021 \pm 0.009 \text{ (stat.)} \pm 0.006 \text{ (syst.)}$$

$$\alpha = 65.9 \pm 38.0 \text{ (stat.)} \pm 17.5 \text{ (syst.) MeV}$$

$$V_{\text{Gaussian}}(r) = -V_{\text{eff}} e^{-\mu r^2}$$

$$V_{\text{eff.}} = 2.5 \pm 0.9 \text{ (stat.)} \pm 1.4 \text{ (syst.) MeV}$$

$$\mu = 0.14 \pm 0.06 \text{ (stat.)} \pm 0.09 \text{ (syst.) fm}^{-2}$$

$$E_{\text{int}} = \int d^3\vec{r} \int d^3\vec{r}' \rho_N(\vec{r}) V(\vec{r} - \vec{r}') \rho_\phi(\vec{r}') \delta^{(3)}(\vec{r}')$$

$\rho_0$   $\delta^{(3)}(\vec{r}')$

$$E_{\text{int}} = -\frac{4\pi A \rho_0}{\alpha^2} = -79.3 \pm 108.8 \text{ MeV}$$

$$E_{\text{int}} = -\frac{\pi^{3/2} V_{\text{eff}} \rho_0}{\mu^{3/2}} = -45.2 \pm 61.5 \text{ MeV}$$

Larger attraction than what was observed at KEK 325, but large statistical and systematic uncertainties

Recent results from theory

# Results for the $\phi$ meson mass at rest (from QCD sum rules)

Most important parameter, that determines the behavior of the  $\phi$  meson mass at finite density:

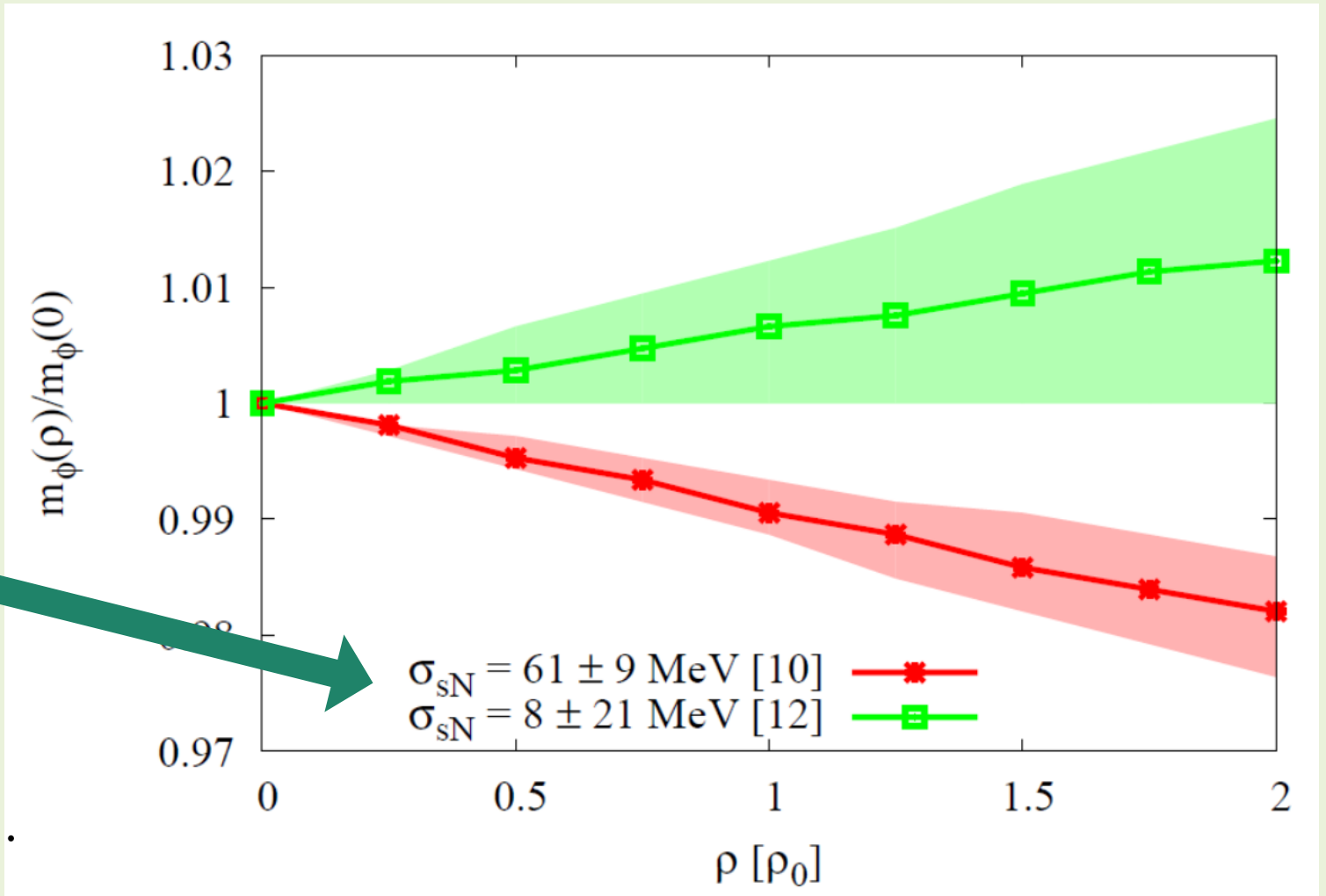
Strangeness content of the nucleon



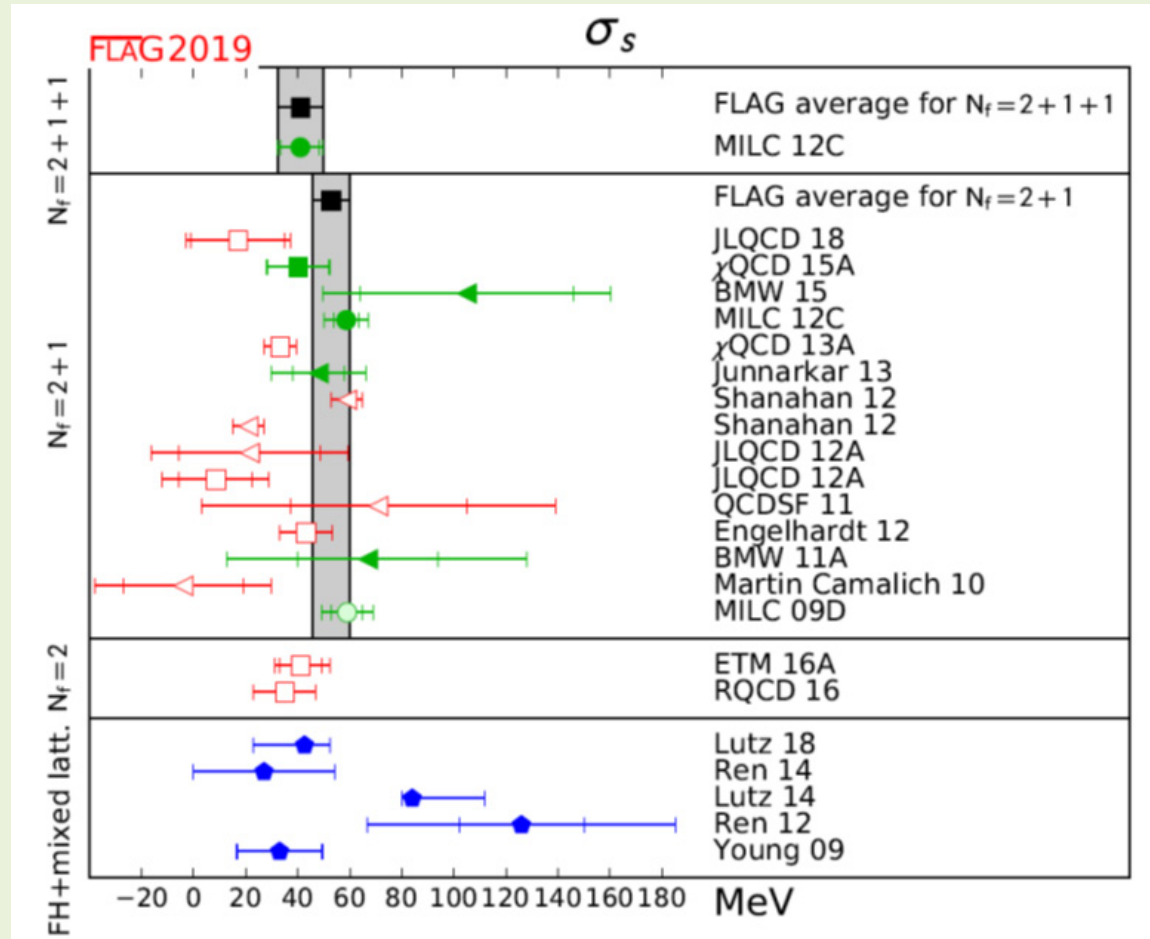
$$\sigma_{sN} = m_s \langle N | \bar{s}s | N \rangle$$



$$\langle \bar{s}s \rangle_\rho = \langle \bar{s}s \rangle_0 + \langle N | \bar{s}s | N \rangle \rho + \dots$$



# What does lattice QCD say about $\sigma_s$ ?



<http://flag.unibe.ch/2019/>

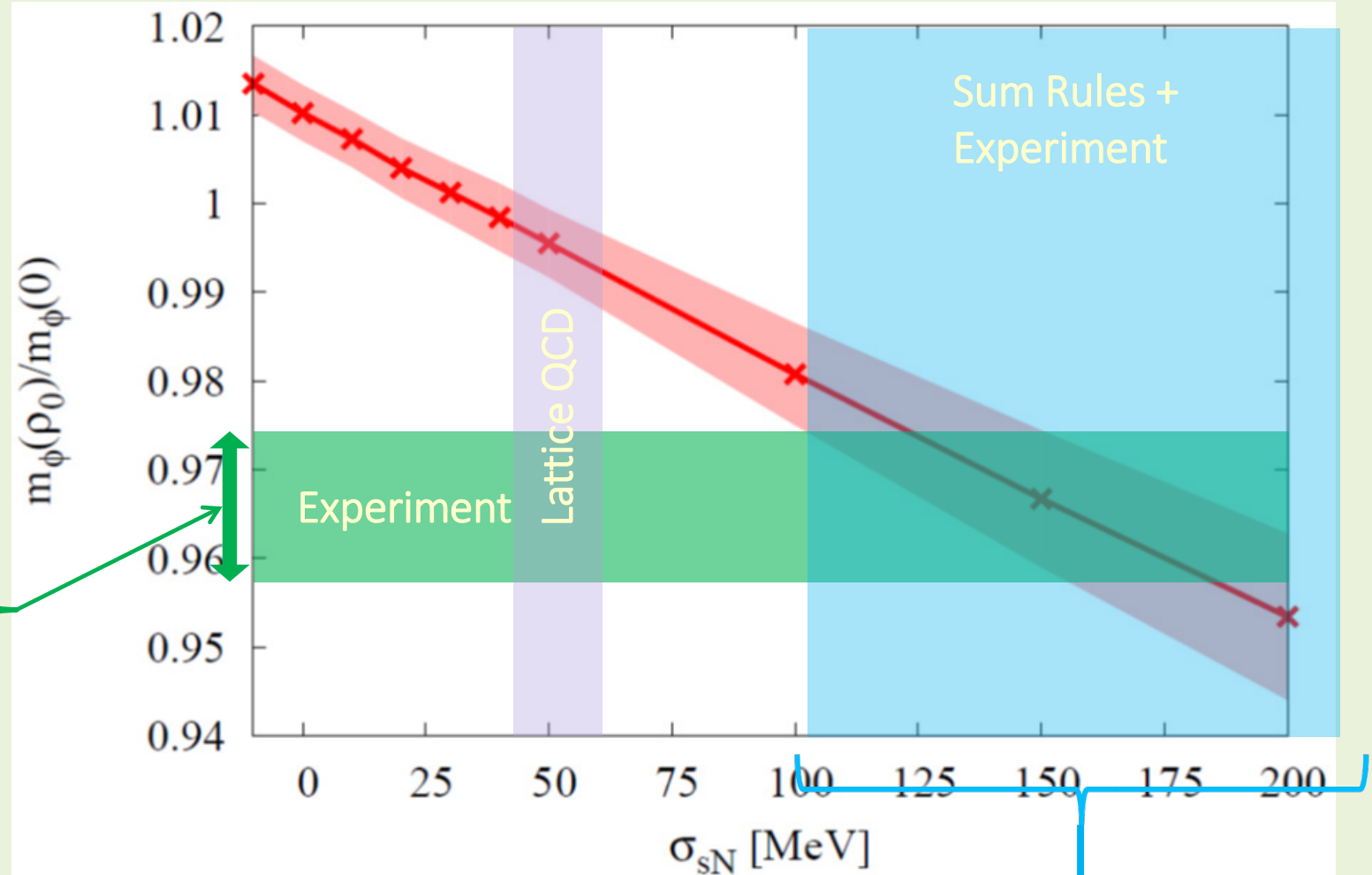
See also the most recent results of the BMW collaboration: Sz. Borsanyi et al., arXiv:2007.03319 [hep-lat].

# Compare Theory with Experiment

Not consistent?

Will be measured again with better statistics at the E16 experiment at J-PARC!

$$\frac{m_\phi(\rho)}{m_\phi(0)} = 0.966 \pm 0.007$$

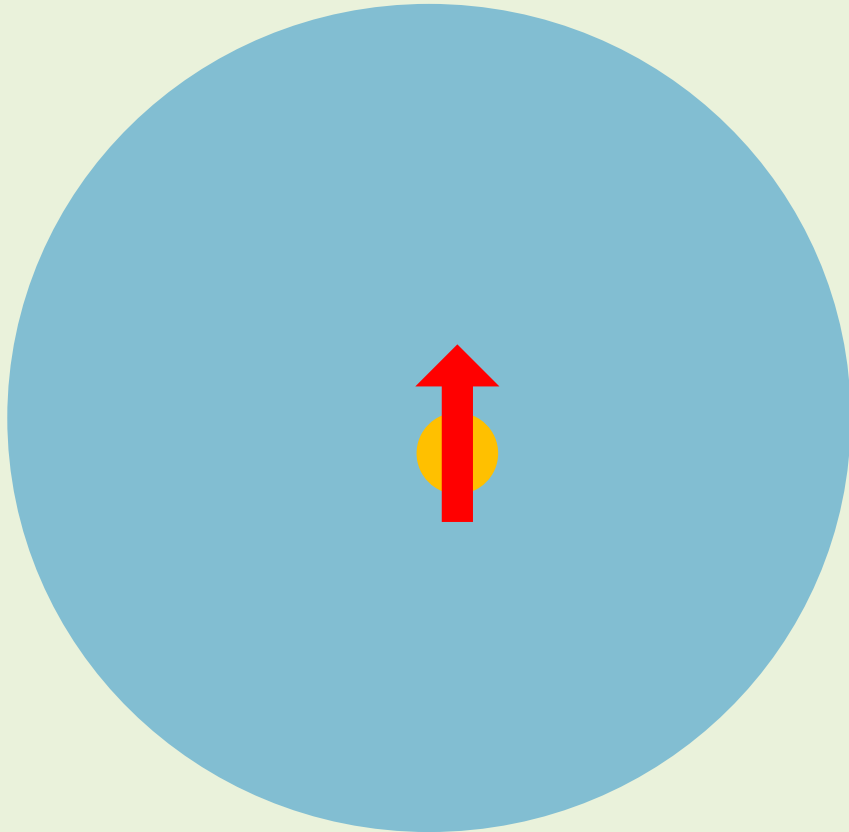


$$\sigma_{sN} \sim 160 \pm 50 \text{ MeV}$$

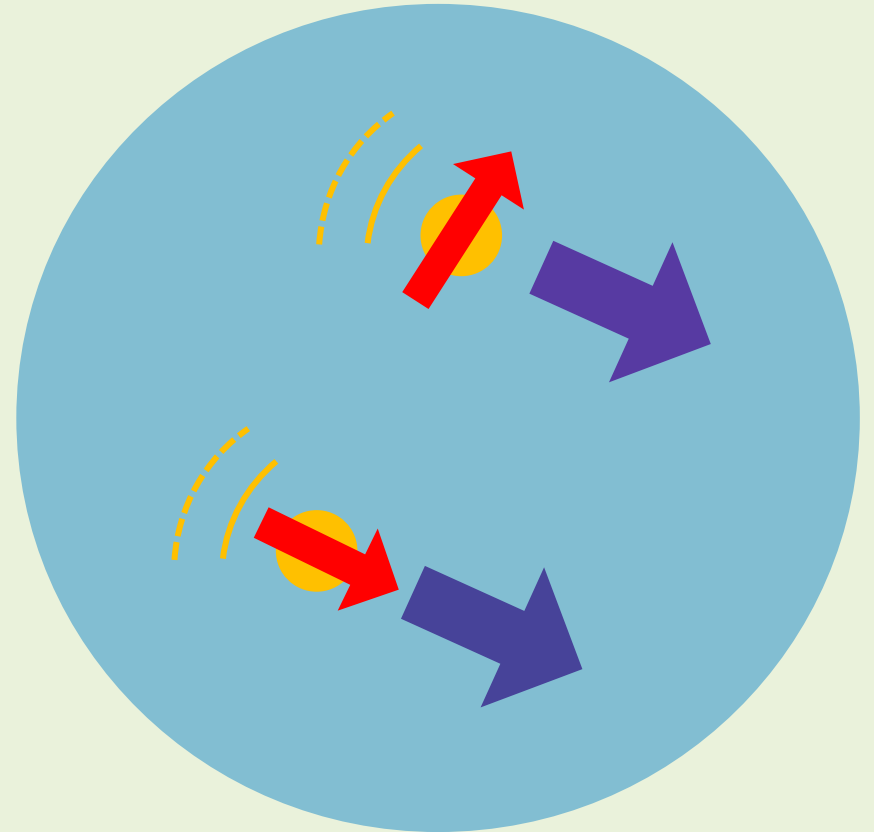


# Effect of momentum

vector meson at rest in nuclear matter



vector meson moving in nuclear matter

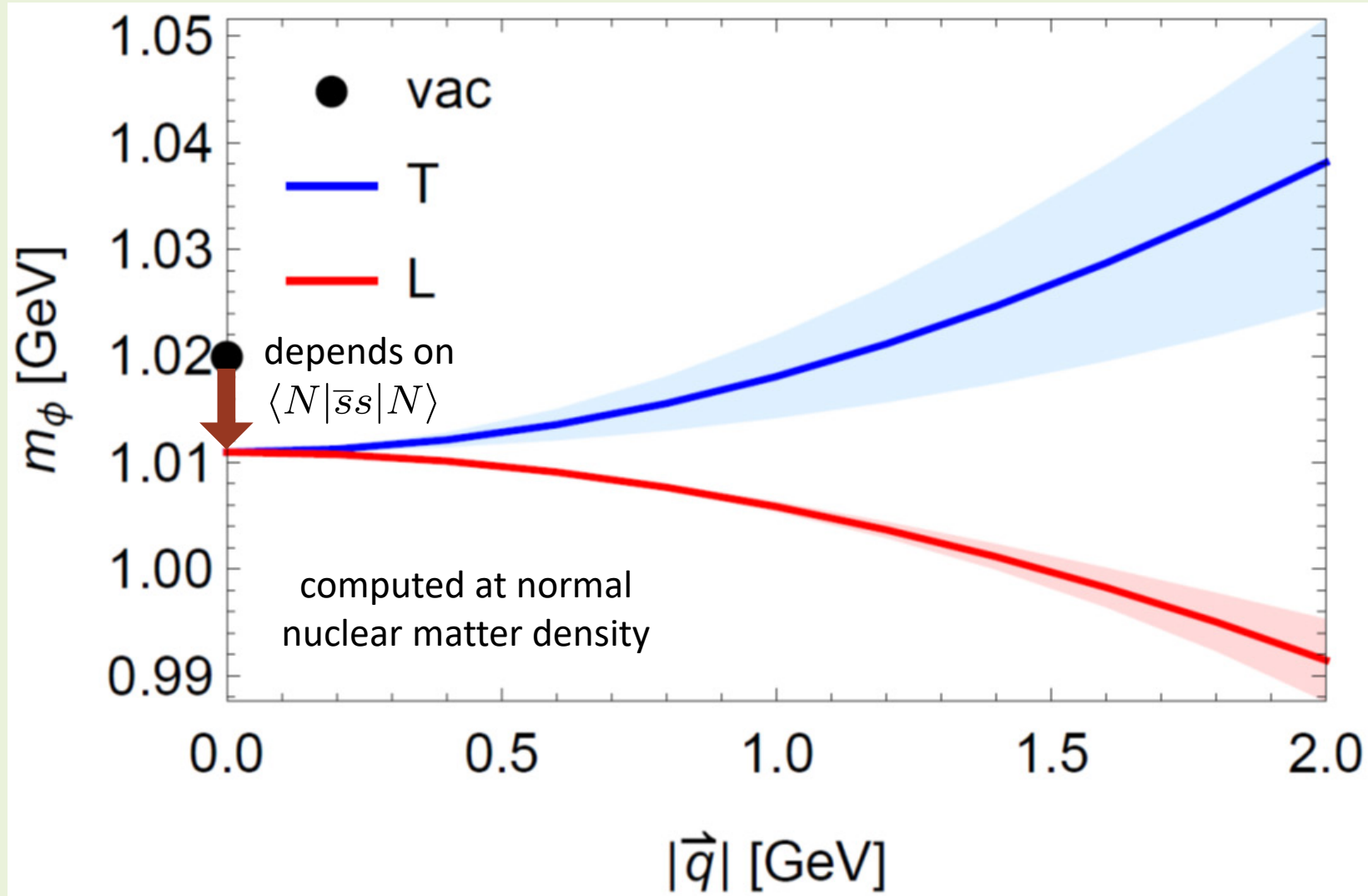


spin direction does not change physics  
(rotational symmetry)



spin direction changes physics  
(broken rotational symmetry)

# Results for the $\phi$ meson mass with non-zero momentum



caused by

$$\langle N | \mathcal{S} \mathcal{T} \bar{s} \gamma^\alpha i D^\beta s | N \rangle$$

+

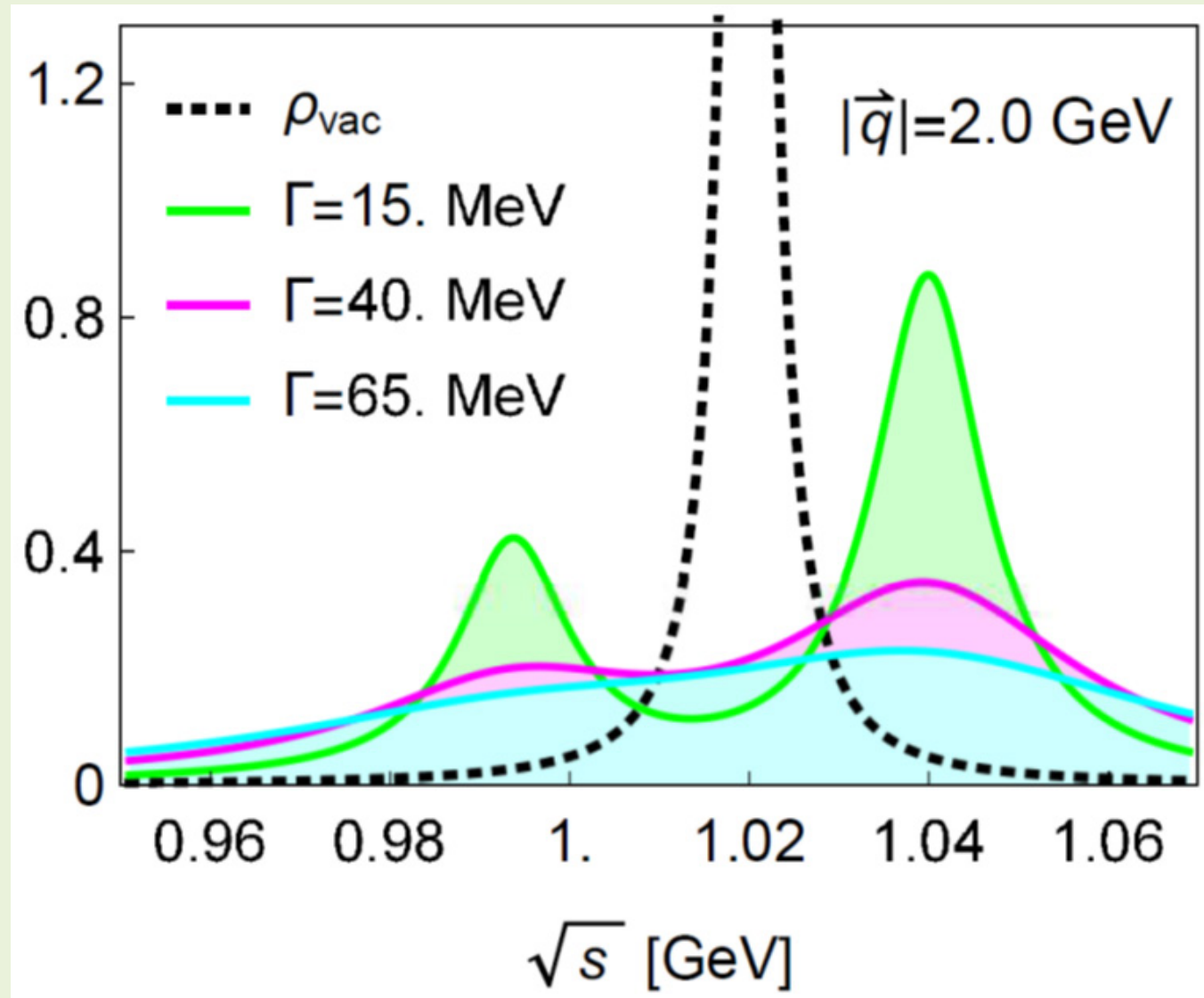
$$\langle N | \mathcal{S} \mathcal{T} G_\mu^{a\alpha} G^{a\mu\beta} | N \rangle$$

caused by

$$\langle N | \mathcal{S} \mathcal{T} G_\mu^{a\alpha} G^{a\mu\beta} | N \rangle$$

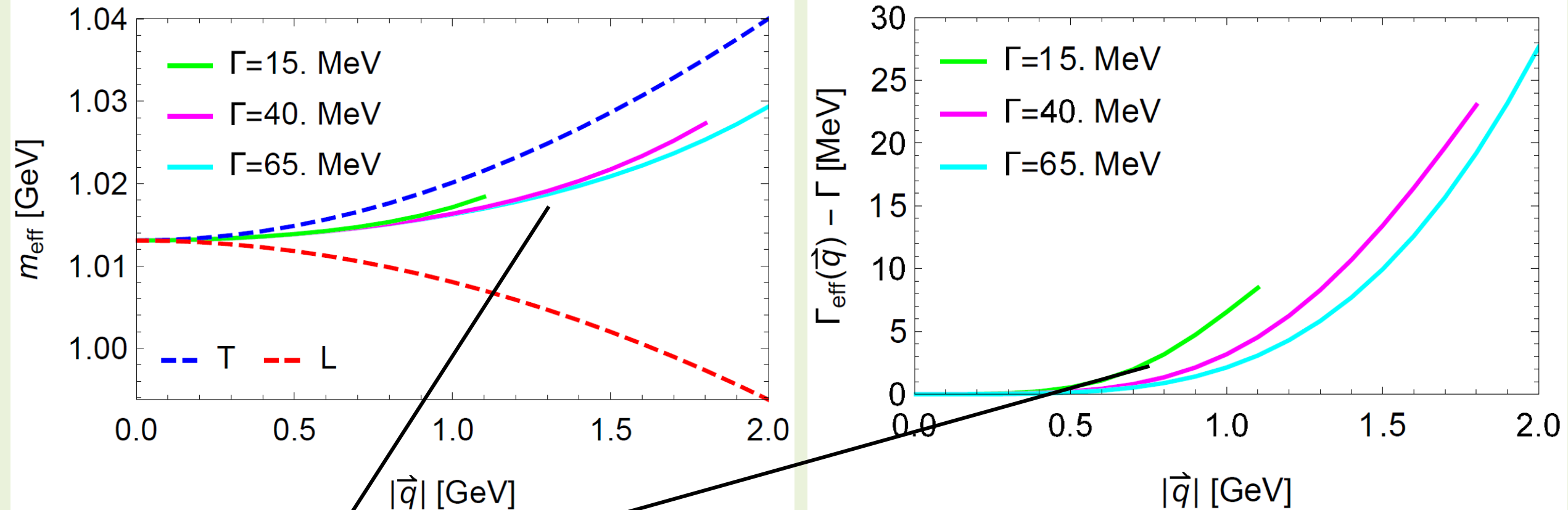
# The angle-averaged di-lepton spectrum

A double peak?



# The angle-averaged di-lepton spectrum

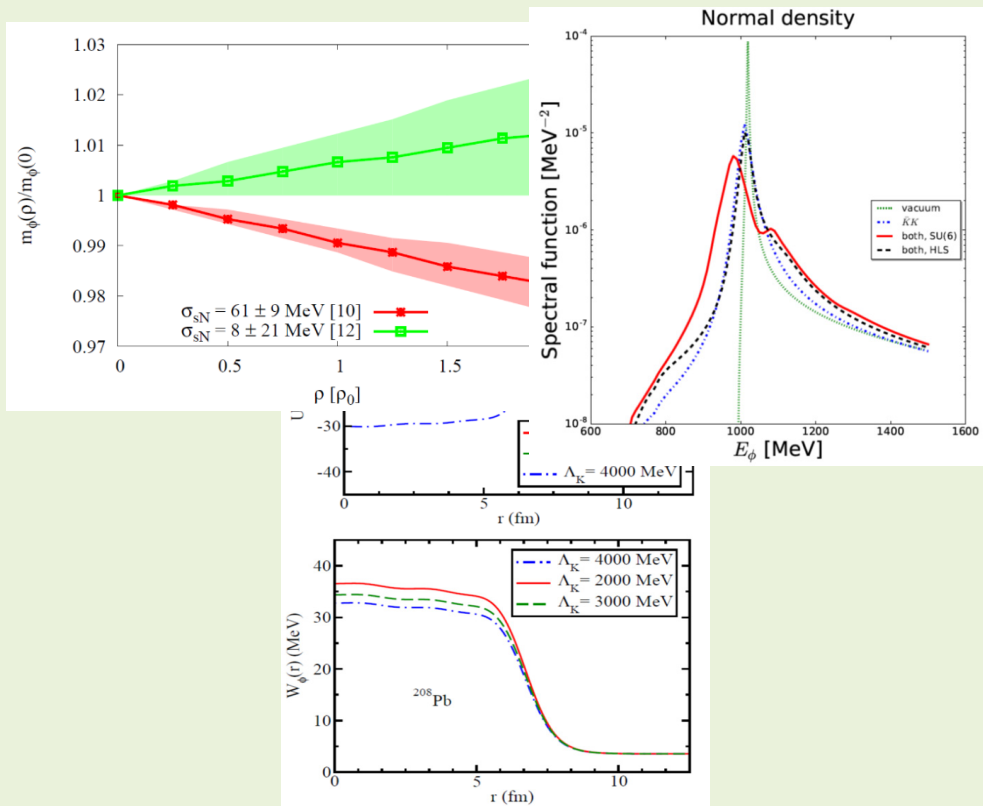
Momentum-dependent mass and width



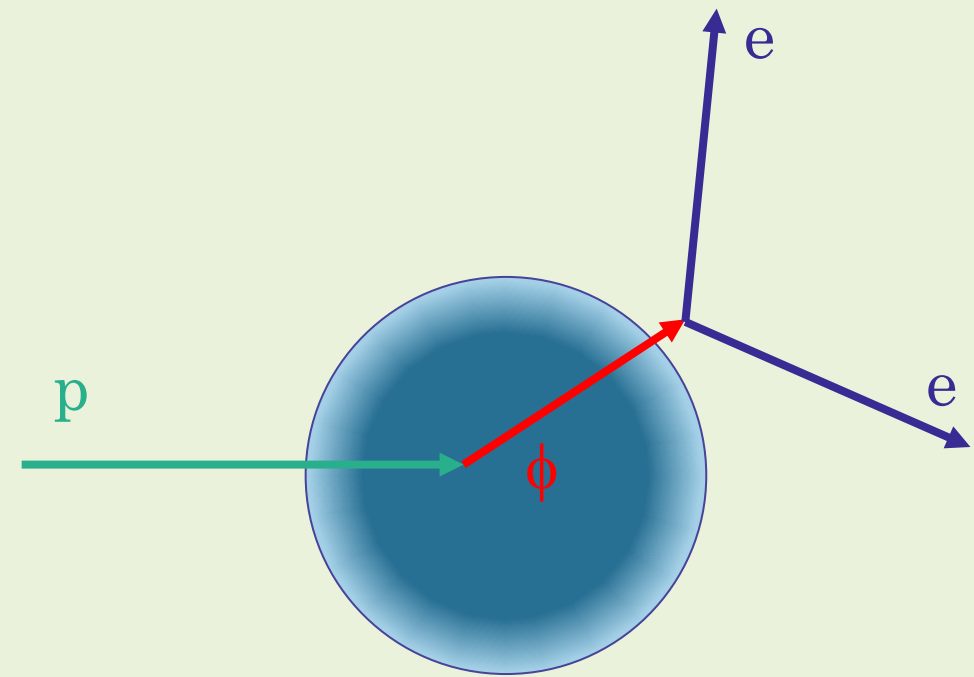
Results of one-peak fits

# How compare theory with experiment?

## Theory



## Experiment



Realistic simulation of pA reaction is needed!

# Our tool: a transport approach

## PHSD (Parton Hadron String Dynamics)

E.L. Bratkovskaya and W. Cassing, Nucl. Phys. A **807**, 214 (2008).

W. Cassing and E.L. Bratkovskaya, Phys. Rev. C **78**, 034919 (2008).

Most important feature for our purposes:

Off-shell dynamics of vector mesons

(dynamical modification of the vector meson spectral function during the simulated reaction)

off-shell terms

Testparticle approach:

$$\frac{d\vec{X}_i}{dt} = \frac{1}{1 - C_{(i)}} \frac{1}{2\varepsilon_i} \left[ 2\vec{P}_i + \vec{\nabla}_{P_i} \operatorname{Re} \Sigma_{(i)}^{\text{ret}} + \frac{\varepsilon_i^2 - \vec{P}_i^2 - M_0^2 - \operatorname{Re} \Sigma_{(i)}^{\text{ret}}}{\tilde{\Gamma}_{(i)}} \vec{\nabla}_{P_i} \tilde{\Gamma}_{(i)} \right],$$

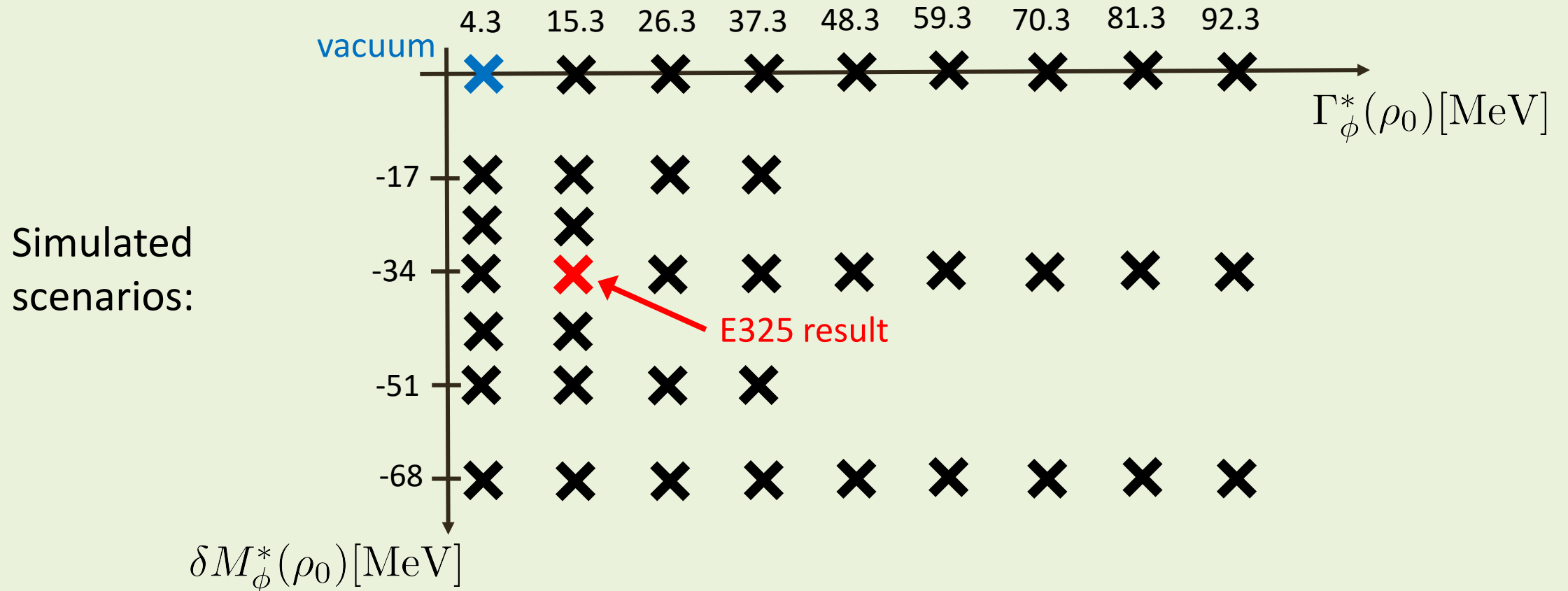
$$\frac{d\vec{P}_i}{dt} = -\frac{1}{1 - C_{(i)}} \frac{1}{2\varepsilon_i} \left[ \vec{\nabla}_{X_i} \operatorname{Re} \Sigma_{(i)}^{\text{ret}} + \frac{\varepsilon_i^2 - \vec{P}_i^2 - M_0^2 - \operatorname{Re} \Sigma_{(i)}^{\text{ret}}}{\tilde{\Gamma}_{(i)}} \vec{\nabla}_{X_i} \tilde{\Gamma}_{(i)} \right],$$

$$\frac{d\varepsilon_i}{dt} = \frac{1}{1 - C_{(i)}} \frac{1}{2\varepsilon_i} \left[ \frac{\partial \operatorname{Re} \Sigma_{(i)}^{\text{ret}}}{\partial t} + \frac{\varepsilon_i^2 - \vec{P}_i^2 - M_0^2 - \operatorname{Re} \Sigma_{(i)}^{\text{ret}}}{\tilde{\Gamma}_{(i)}} \frac{\partial \tilde{\Gamma}_{(i)}}{\partial t} \right],$$

# Advantage: vector meson spectra can be chosen freely

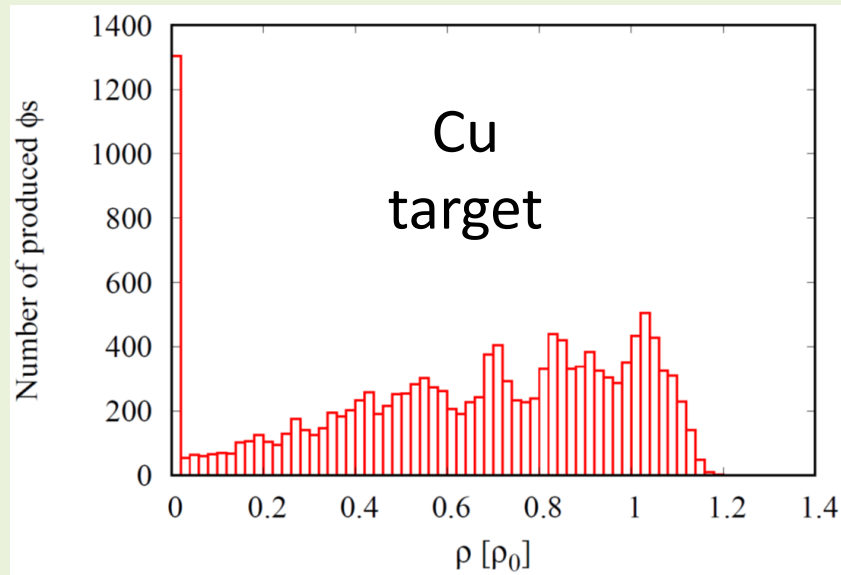
Our choice: a Breit-Wigner with density dependent mass and width

$$A_\phi(M, \rho) = C \frac{2}{\pi} \frac{M^2 \Gamma_\phi^*(M, \rho)}{[M^2 - M_\phi^{*2}(\rho)]^2 + M^2 \Gamma_\phi^{*2}(M, \rho)} \quad \text{with} \quad \begin{cases} M_\phi^*(\rho) = M_\phi^{\text{vac}} \left(1 - \alpha^\phi \frac{\rho}{\rho_0}\right), \\ \Gamma_\phi^*(M, \rho) = \Gamma_\phi^{\text{vac}} + \alpha_{\text{coll}}^\phi \frac{\rho}{\rho_0} \end{cases}$$

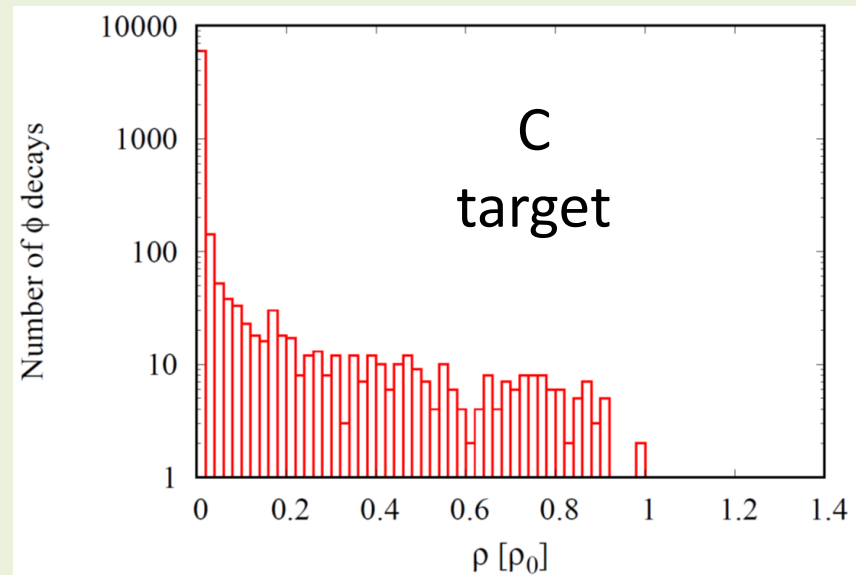
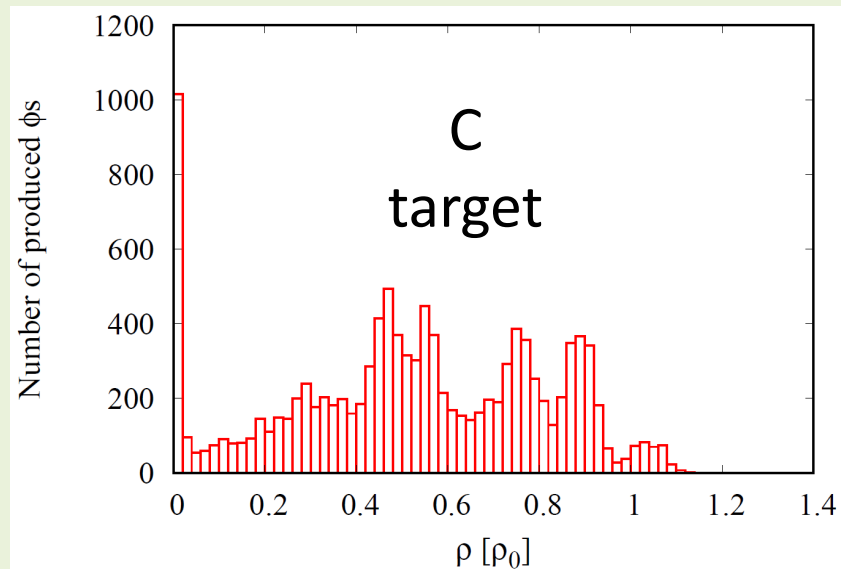
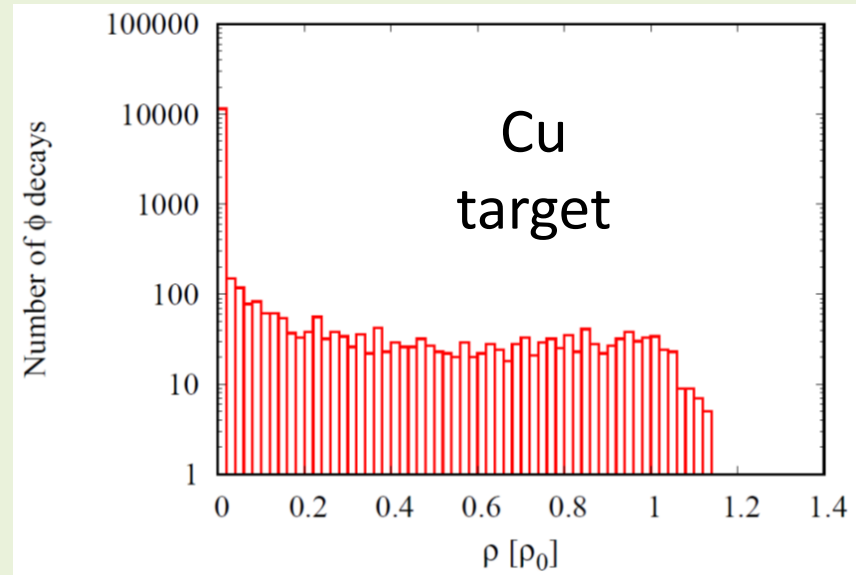


# What density does the $\phi$ feel in the reaction (p+Cu/C at 12 GeV)?

## Production

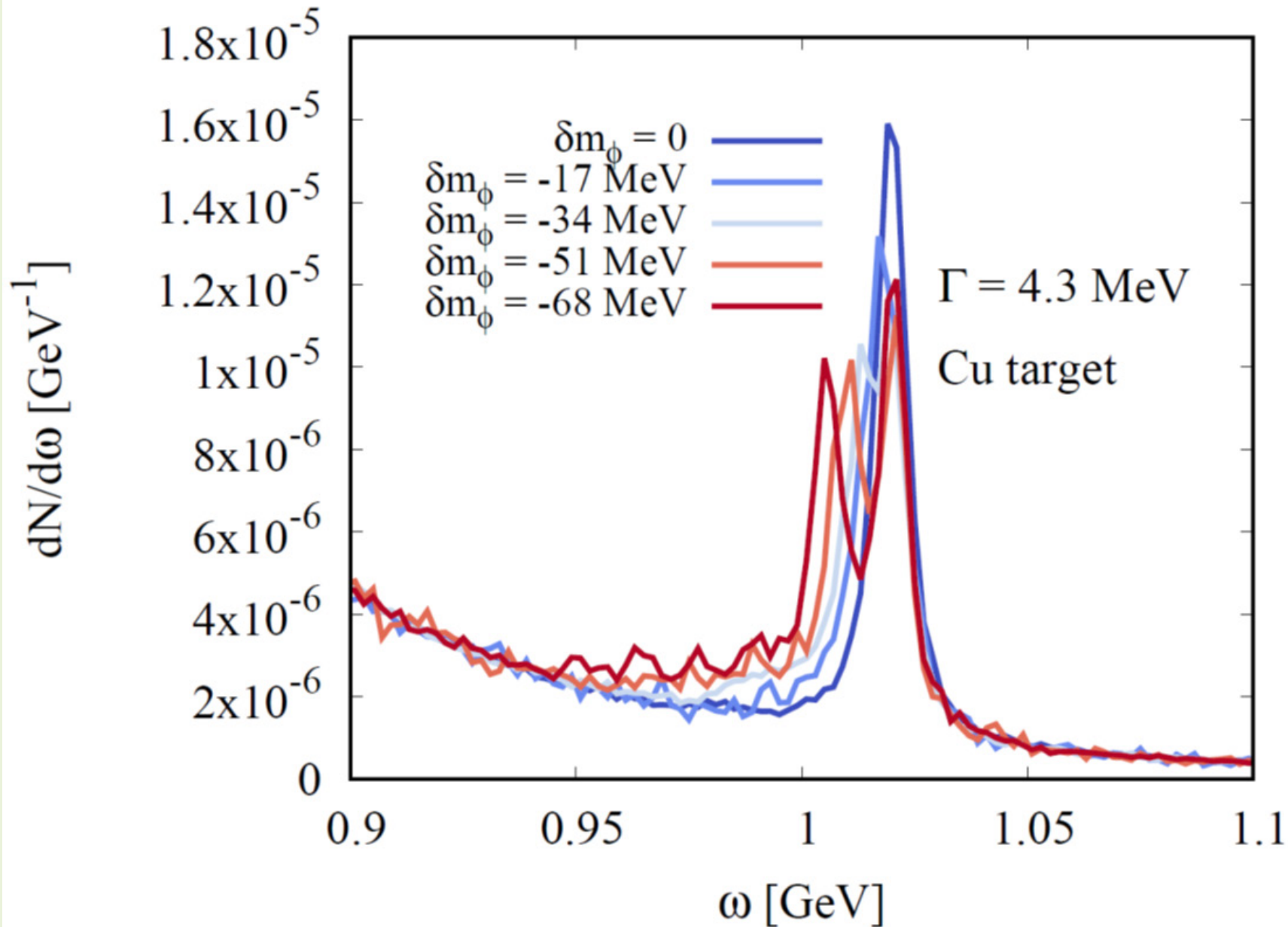


## Decay





# The dilepton spectrum in the $\phi$ meson region

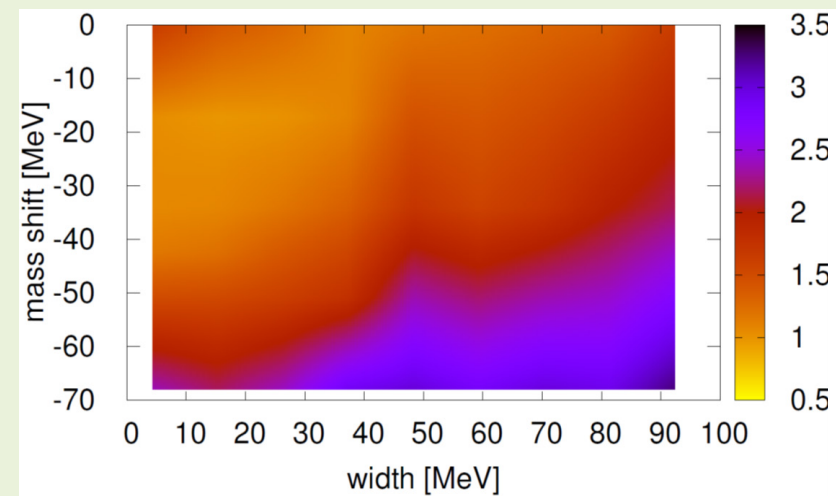
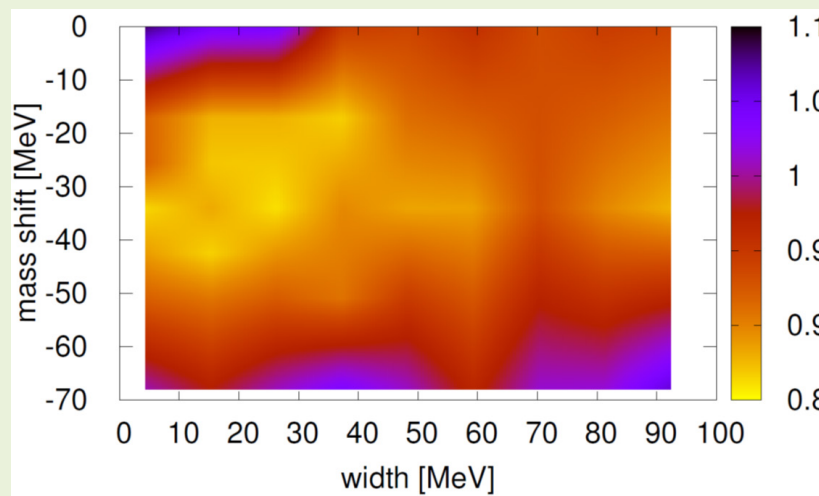
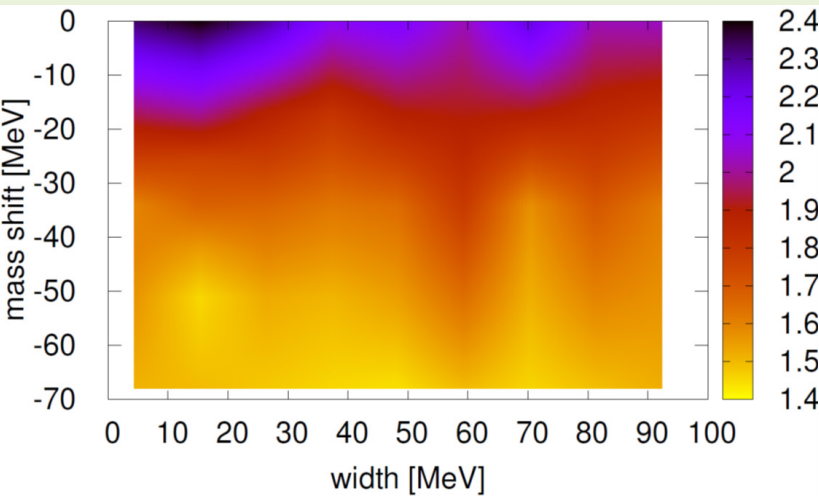
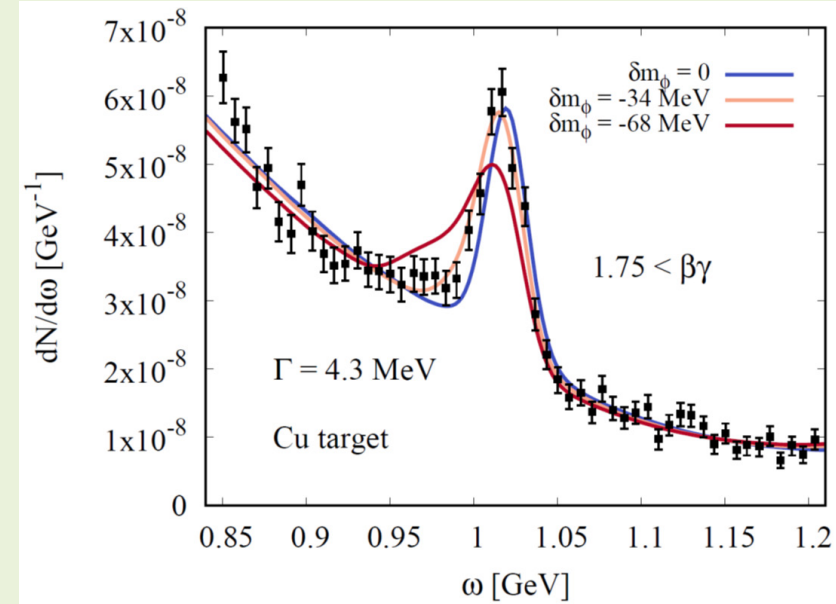
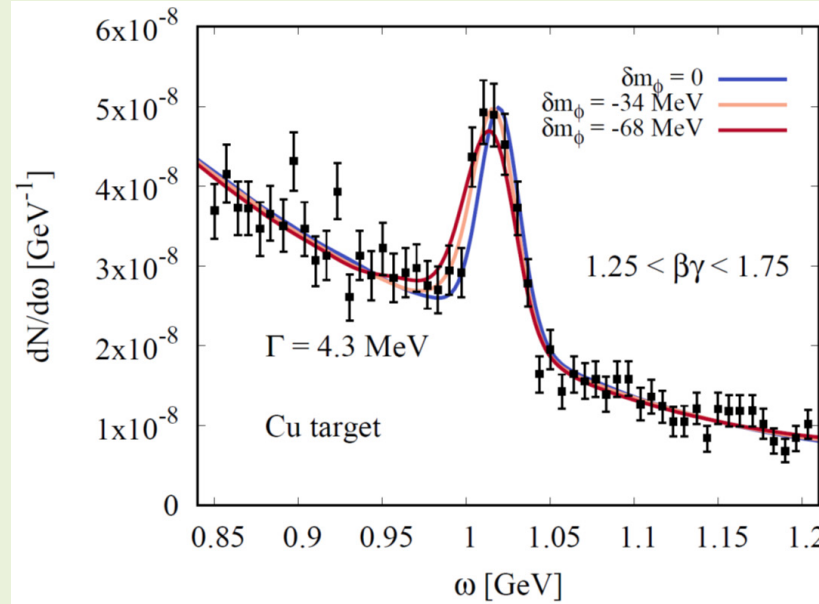
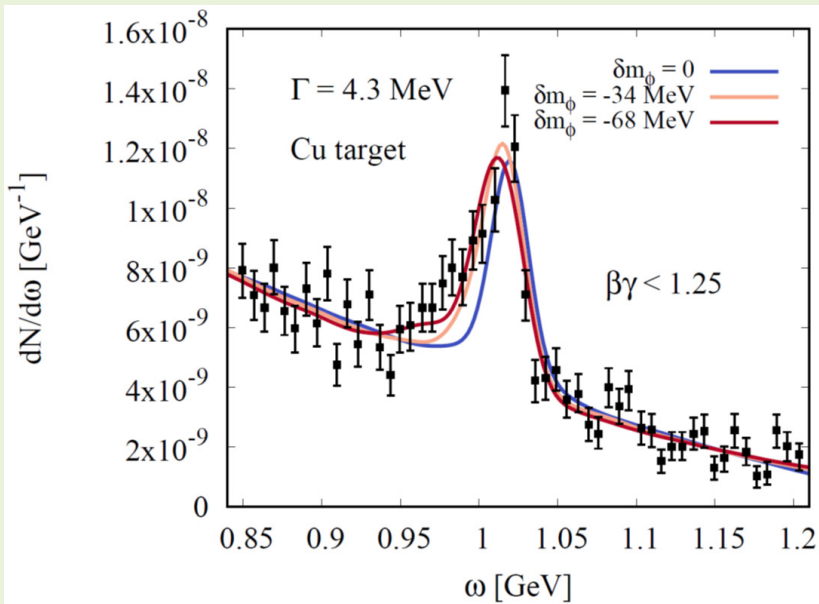


p + Cu at 12 GeV

No acceptance  
corrections!

No finite  
resolution effects!

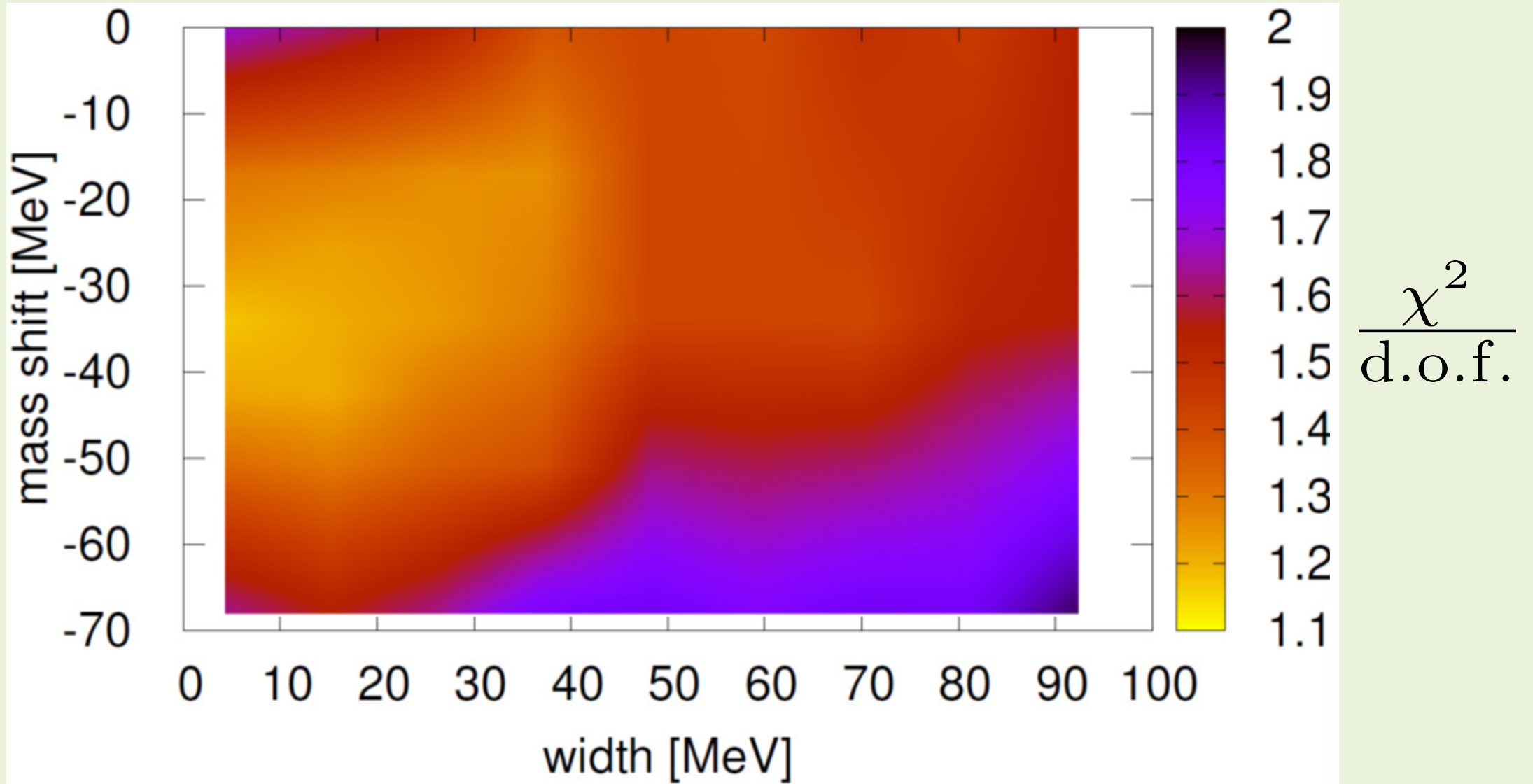
# Fits to experimental Copper target data (E325)



Favors relatively large negative mass shift!

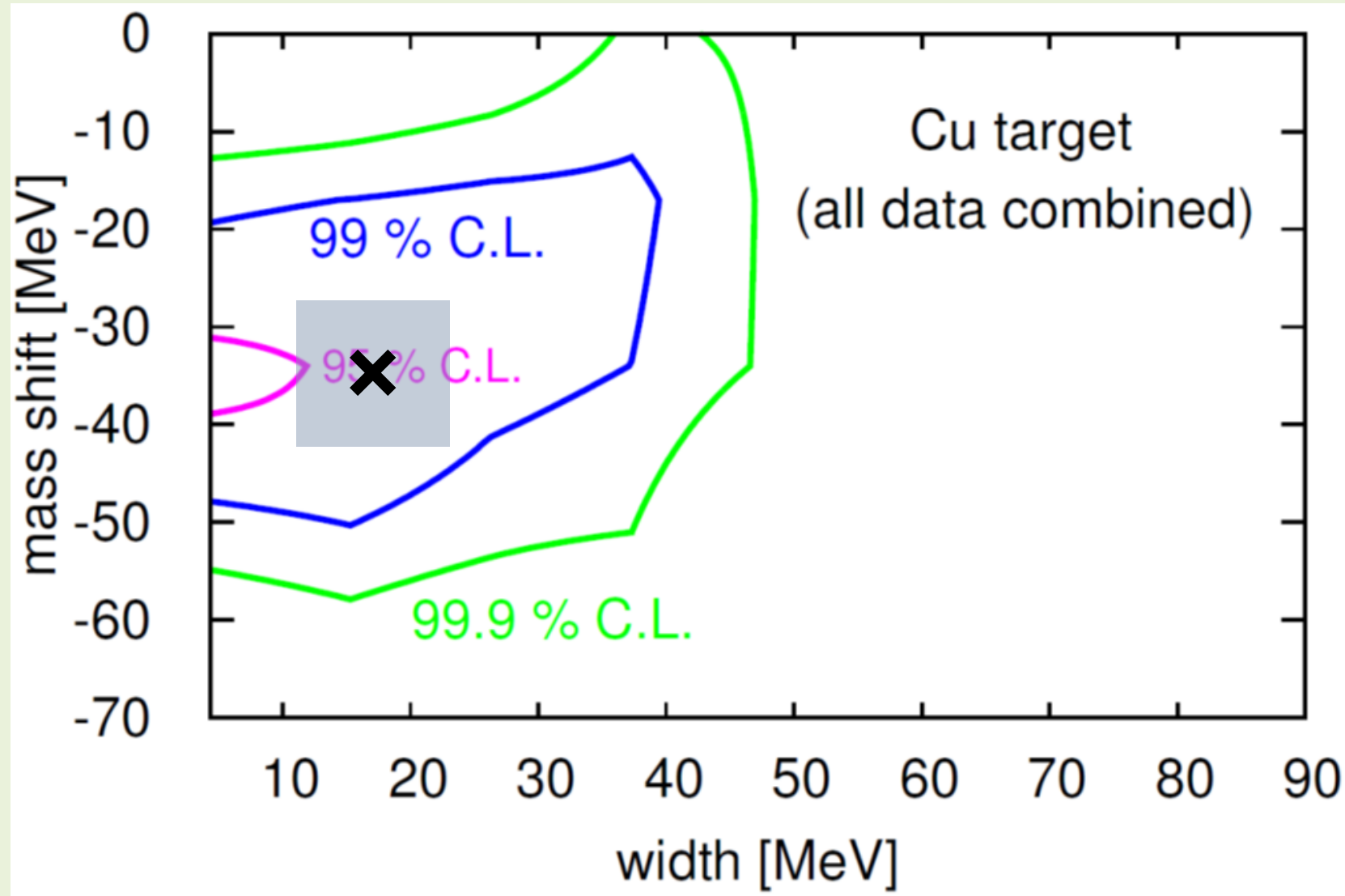
Favors small mass shift!

All Copper data combined

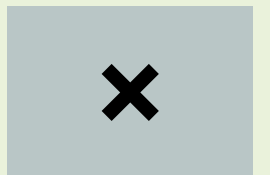


# Fits to experimental Copper target data (E325)

Confidence levels of combined Copper data

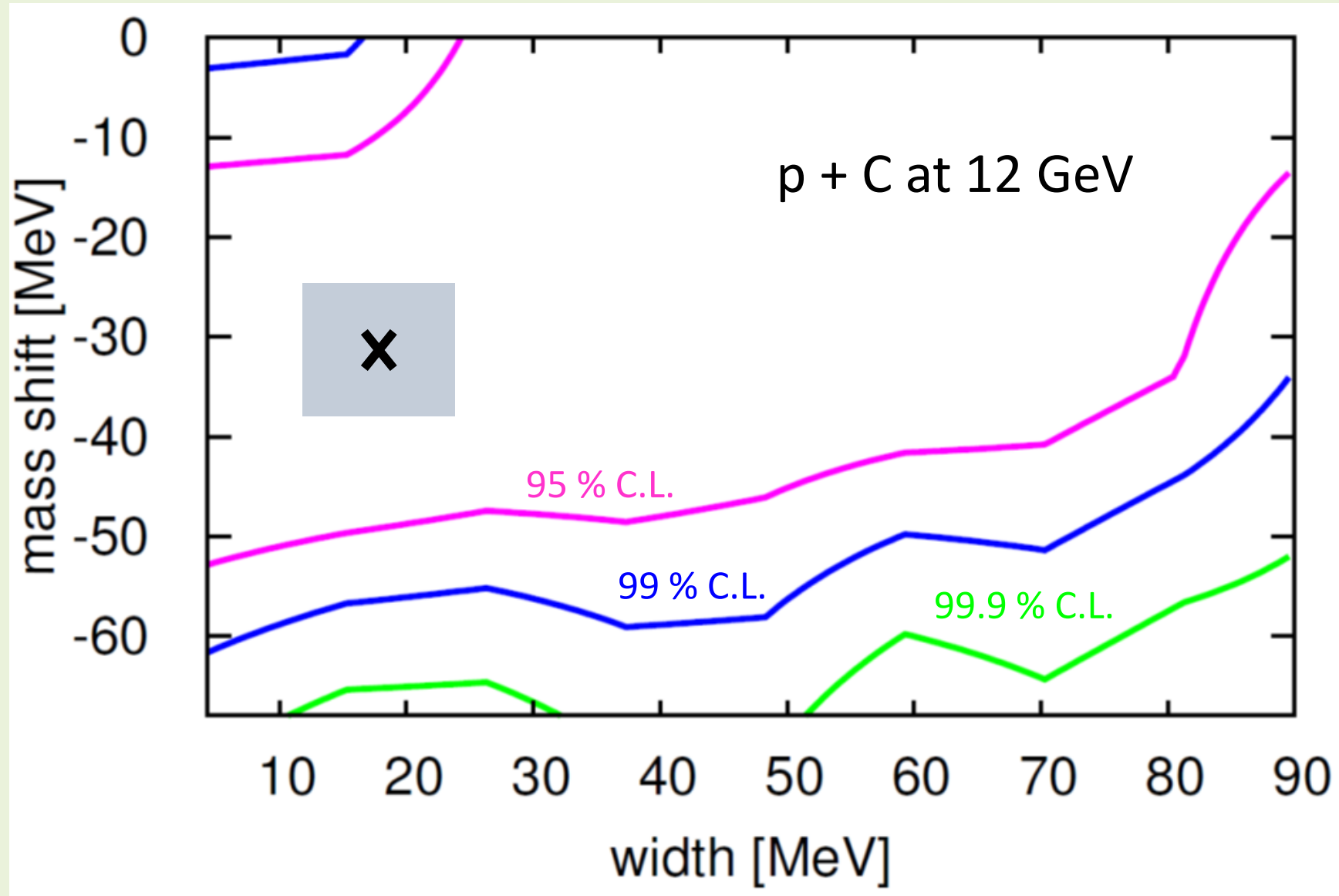


Conclusion of the  
E325 Collaboration

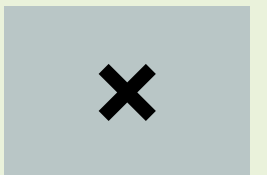


# Fits to experimental Carbon target data (E325)

Confidence levels of combined Carbon data



Conclusion of the  
E325 Collaboration

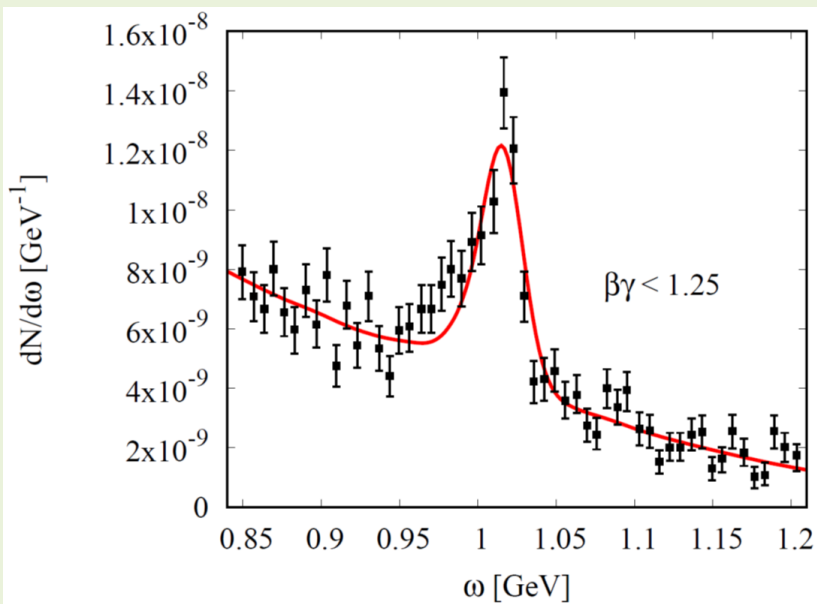


# Best fit to E325 data (p + Cu at 12 GeV)

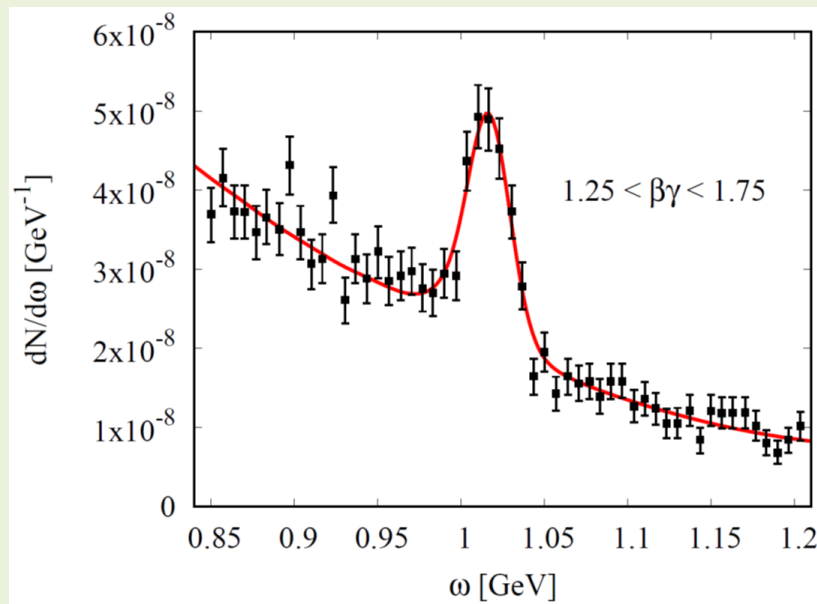
$$\delta m_\phi(\rho_0) = -34 \text{ MeV}$$

$$\Gamma(\rho_0) = 4.3 \text{ MeV}$$

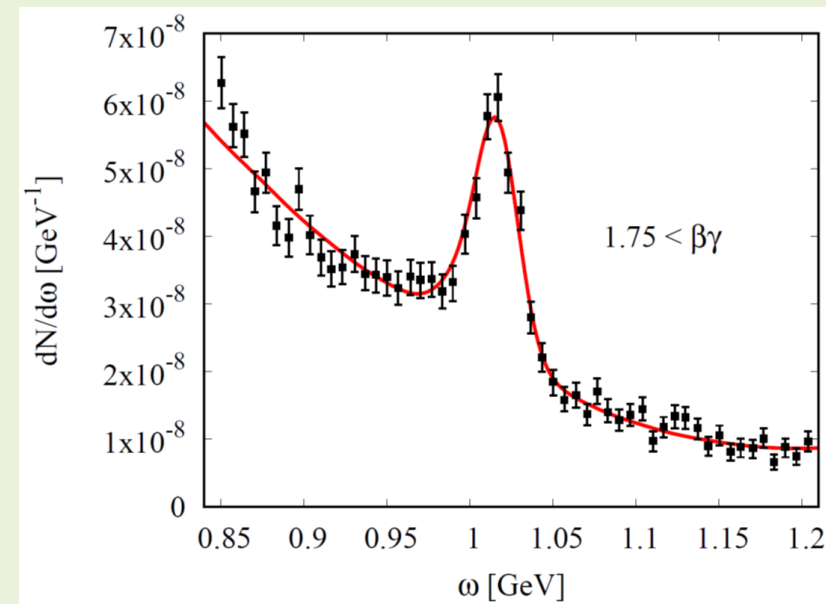
vacuum value



slow  $\phi$ s



intermediate  $\phi$ s



fast  $\phi$ s



# Summary and Conclusions

- ★ A lot of new experimental information about the  $\varphi$ N interaction is becoming available (HADES, LHC, J-PARC)
- ★ Studying the modification of the  $\varphi$  meson spectral function experimentally At finite density is non-trivial. A good understanding of the underlying reactions is needed!
- ★ Numerical simulations of the pA reactions to measured at the E325 experiment at KEK, using the PHSD transport code, are in progress.



First results indicate that the experimental data favor a **negative mass shift with none or only small broadening!**

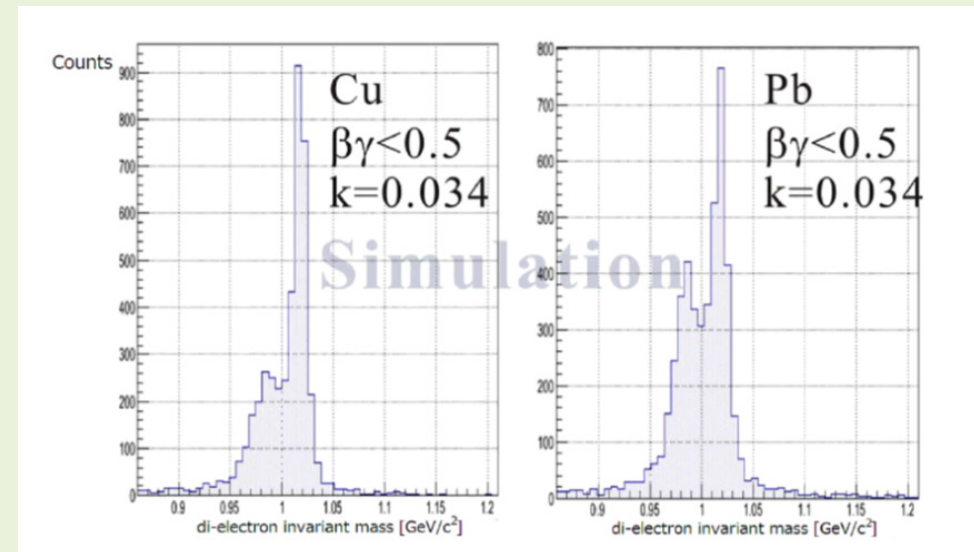
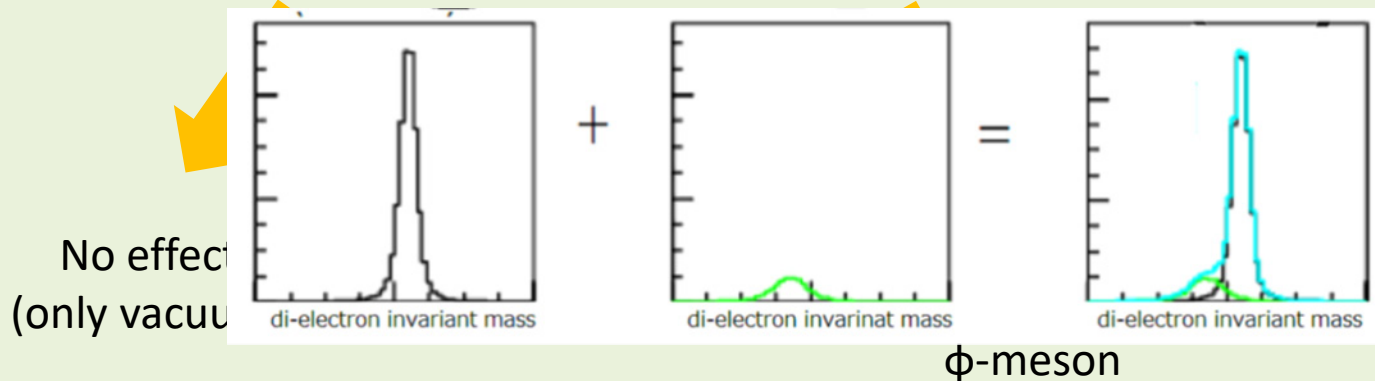
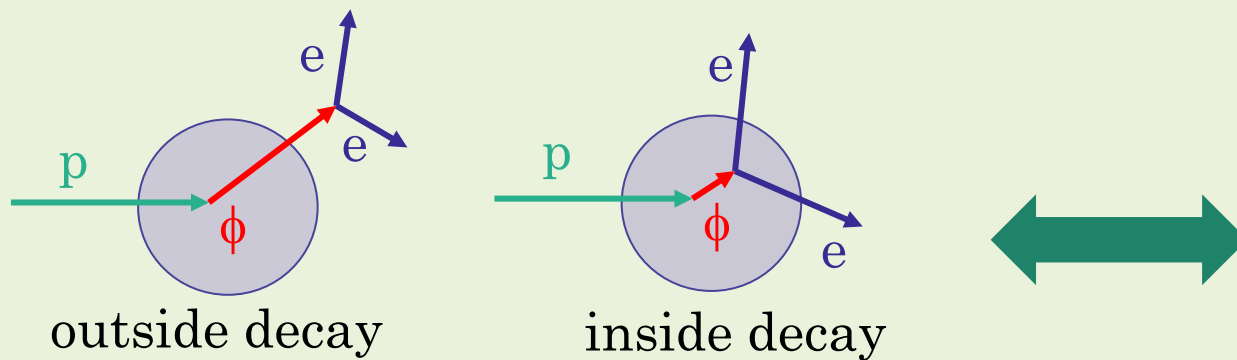
Backup slides



# The experimental situation

## The E325 Experiment (KEK)

Slowly moving  $\phi$  mesons are produced in 12 GeV  $p+A$  reactions and are measured through di-leptons.



Y. Morino et. al. (J-PARC E16 Collaboration),  
JPS Conf. Proc. 8, 022009 (2015).

# Our tool: a transport code PHSD (Parton Hadron String Dynamics)

W. Cassing and E. Bratkovskaya, Phys. Rev. C **78**, 034919 (2008).

Example:

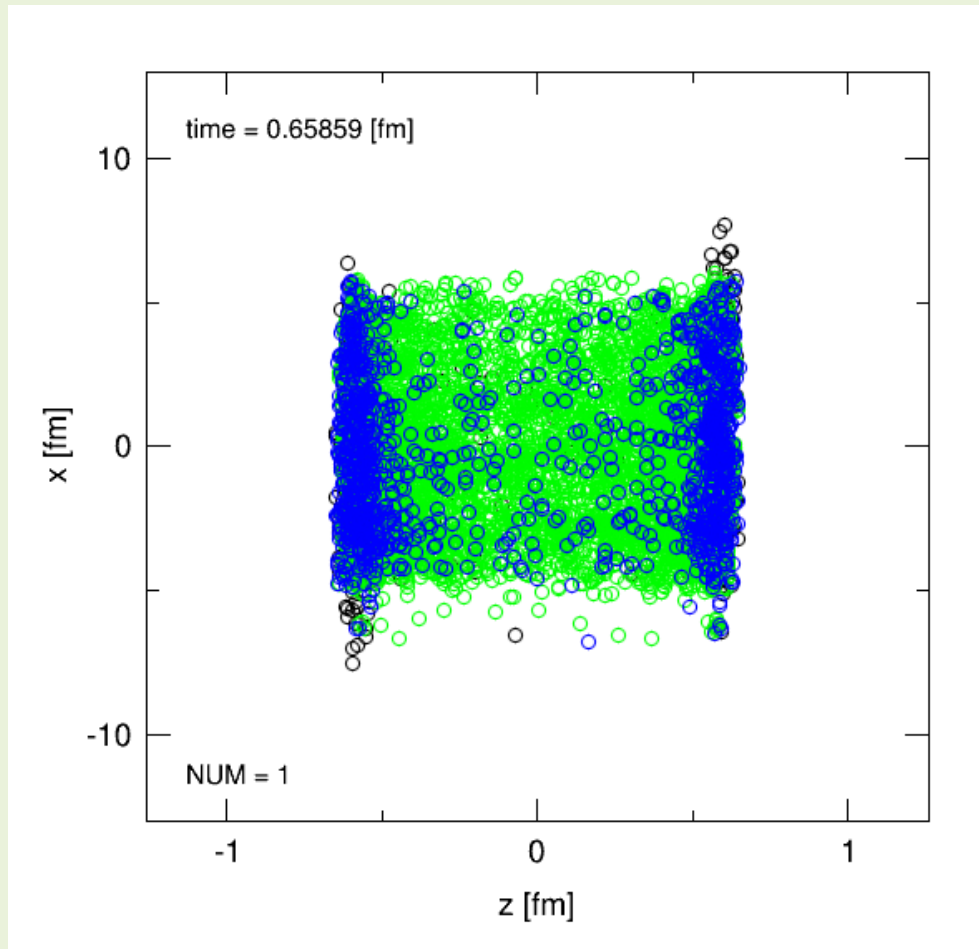
Au+Au collision at 200 GeV

$b = 2$  fm

nucleons

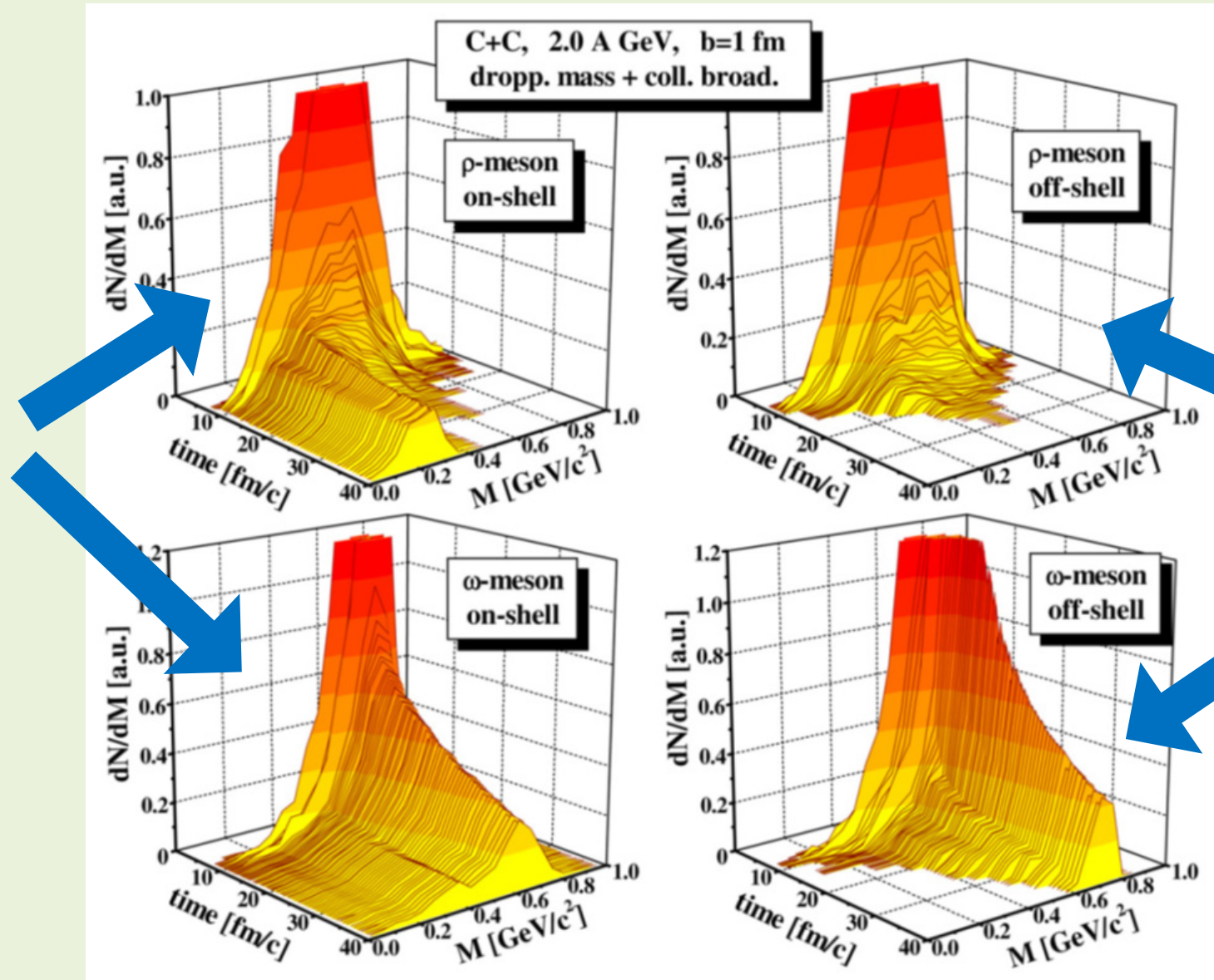
quarks

gluons



# The importance of off-shell contributions

Only on-shell contributions:  
Vacuum spectral function  
are not recovered at late  
time of the reaction



Off-shell  
contributions  
included:  
correct behavior

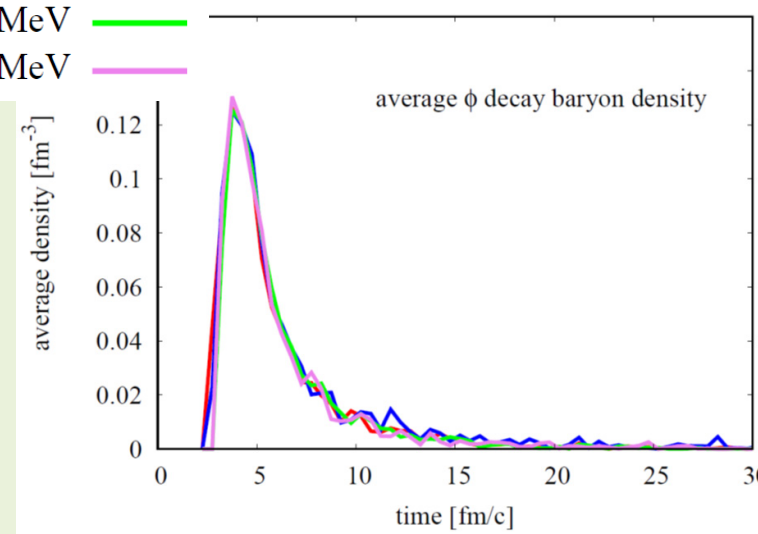
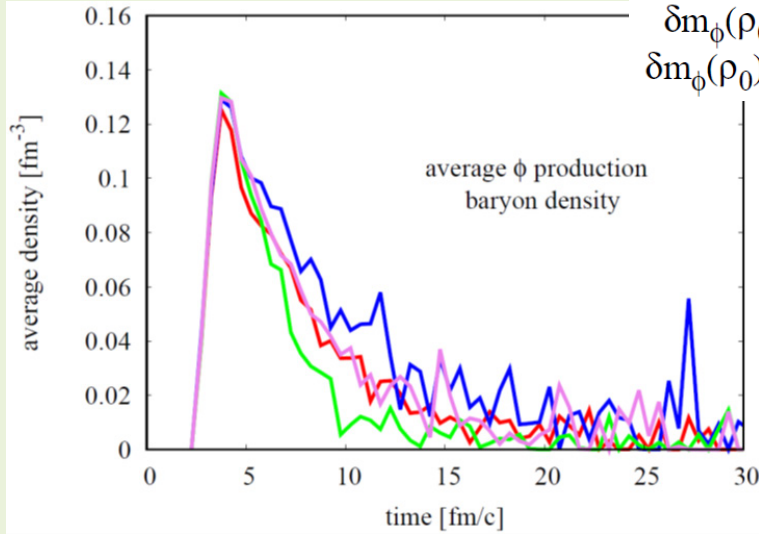
# What density does the $\phi$ feel in the reaction (p+Cu/C at 12 GeV)?

Production

Decay

- $\delta m_\phi(\rho_0) = 0, \Gamma(\rho_0) = 4.3 \text{ MeV}$  — (red line)
- $\delta m_\phi(\rho_0) = 0, \Gamma(\rho_0) = 59.3 \text{ MeV}$  — (blue line)
- $\delta m_\phi(\rho_0) = -68 \text{ MeV}, \Gamma(\rho_0) = 4.3 \text{ MeV}$  — (green line)
- $\delta m_\phi(\rho_0) = -68 \text{ MeV}, \Gamma(\rho_0) = 59.3 \text{ MeV}$  — (magenta line)

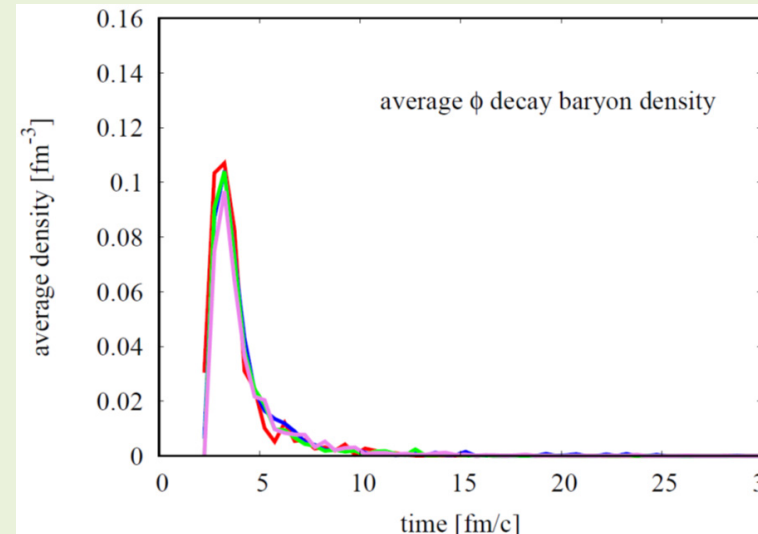
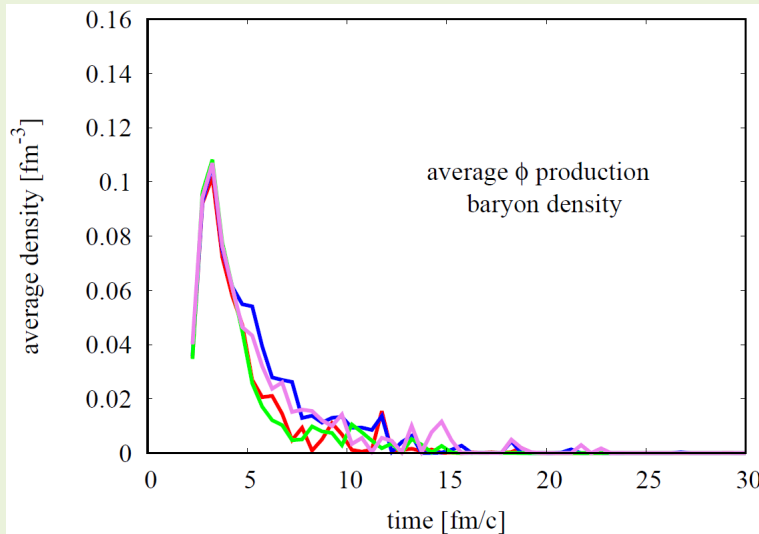
Cu  
target



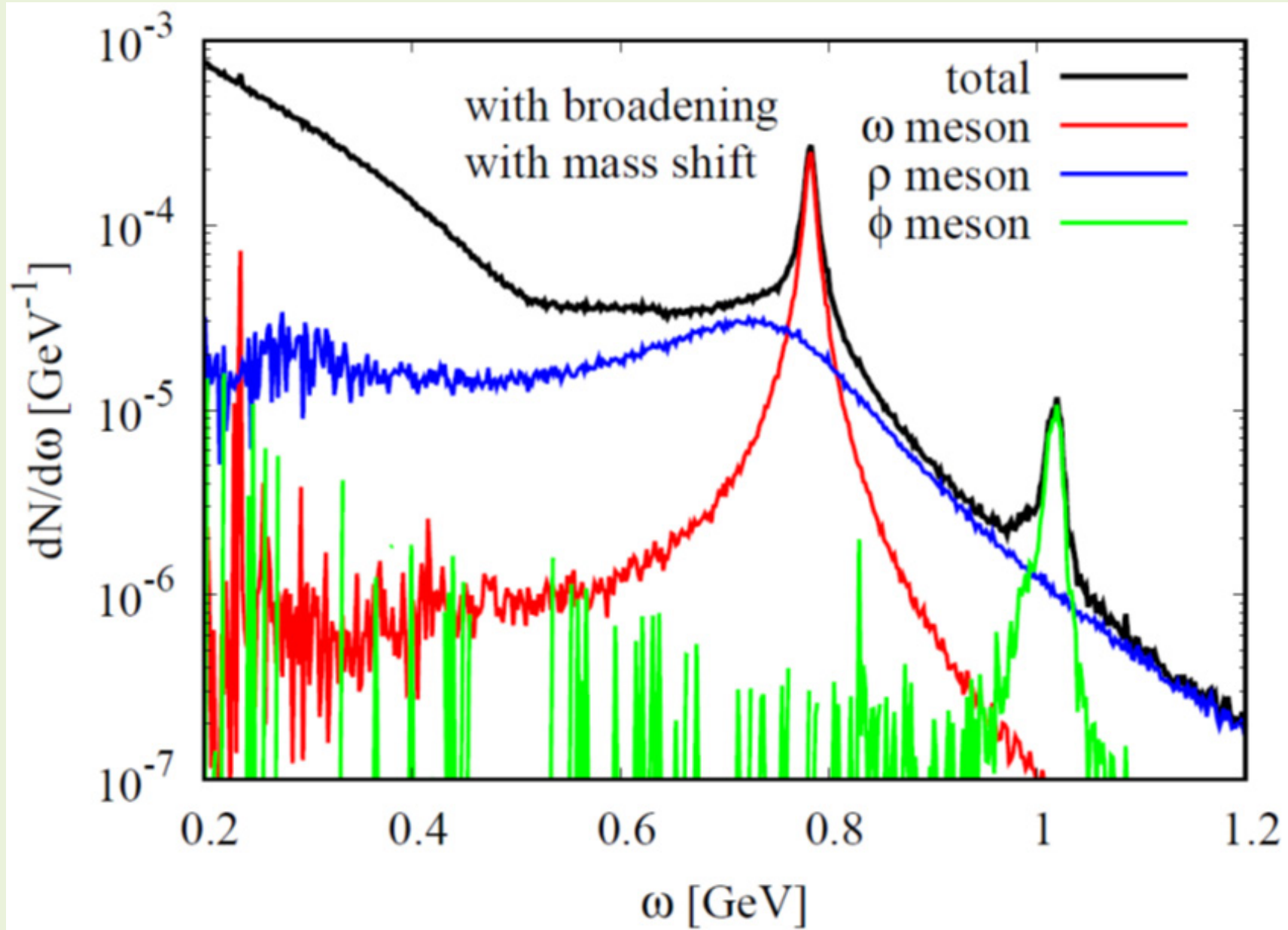
$\phi$  mesons are on average created at a density significantly below  $\rho_0$

$\phi$  mesons on average decay at a density significantly below  $\rho_0$

C  
target



# The dilepton spectrum

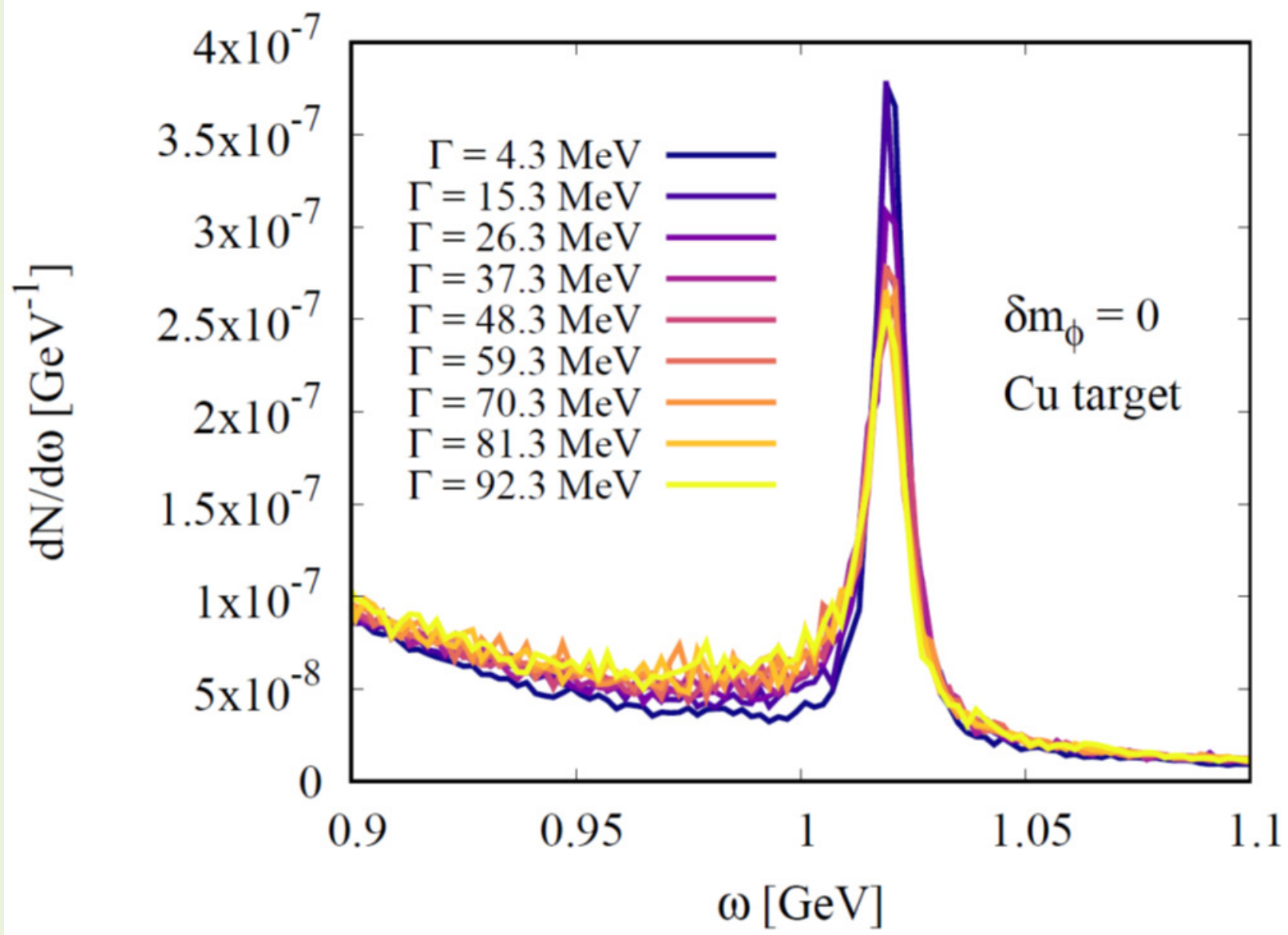


p+Cu at 12 GeV

The  $\phi$  meson peak  
is clearly visible.



# The dilepton spectrum in the $\phi$ meson region



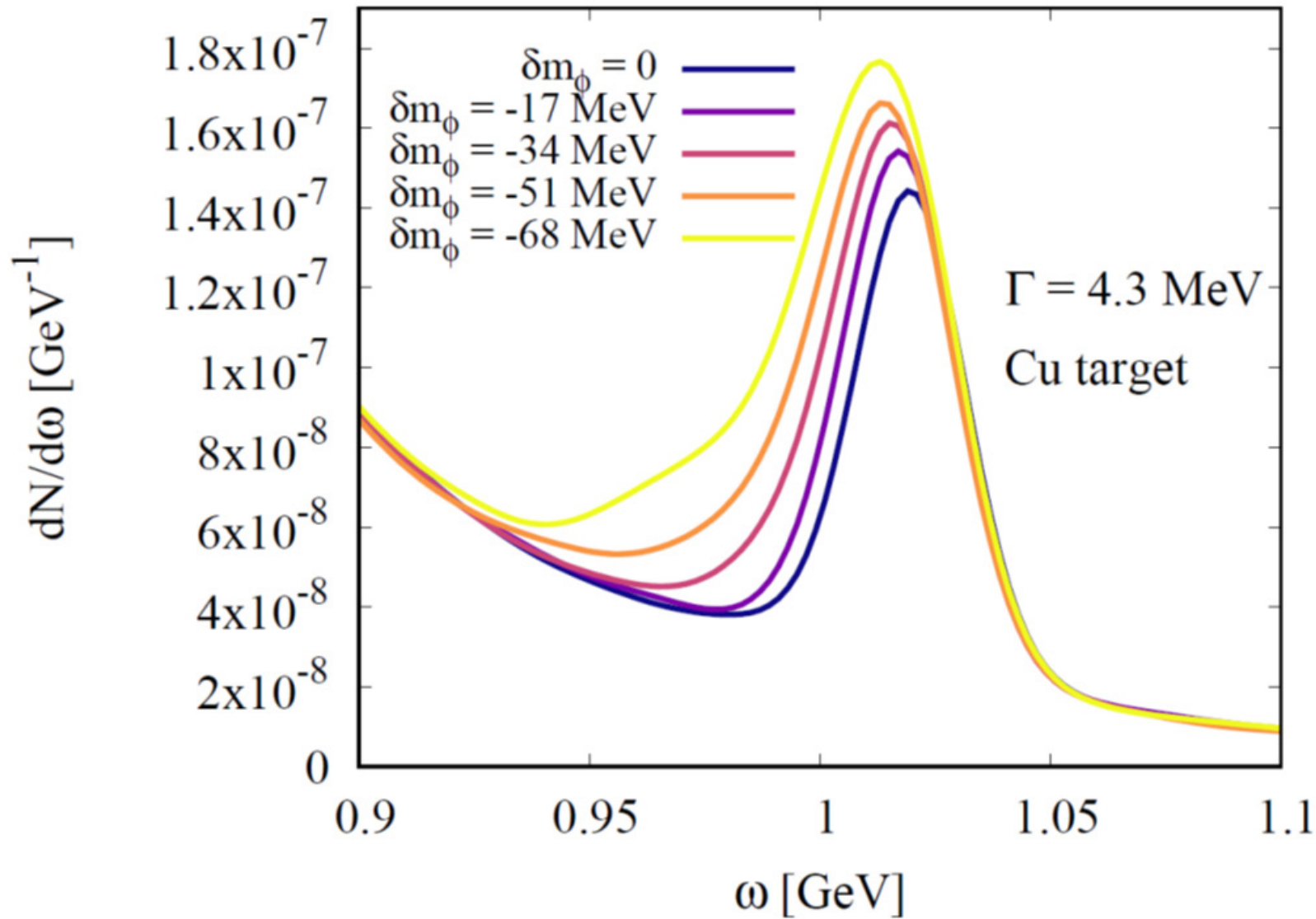
p + Cu at 12 GeV

No acceptance  
corrections!

No finite  
resolution effects!

Preliminary

# The dilepton spectrum in the $\phi$ meson region



p + Cu at 12 GeV

With acceptance  
corrections!

With finite  
resolution effects!

# Final step: comparison to experimental data

- Potential issues:
- ★ Experimental background is not included in the simulation
  - ★ Normalization of the experimental dilepton spectrum is not given



Fit to experimental data is necessary!

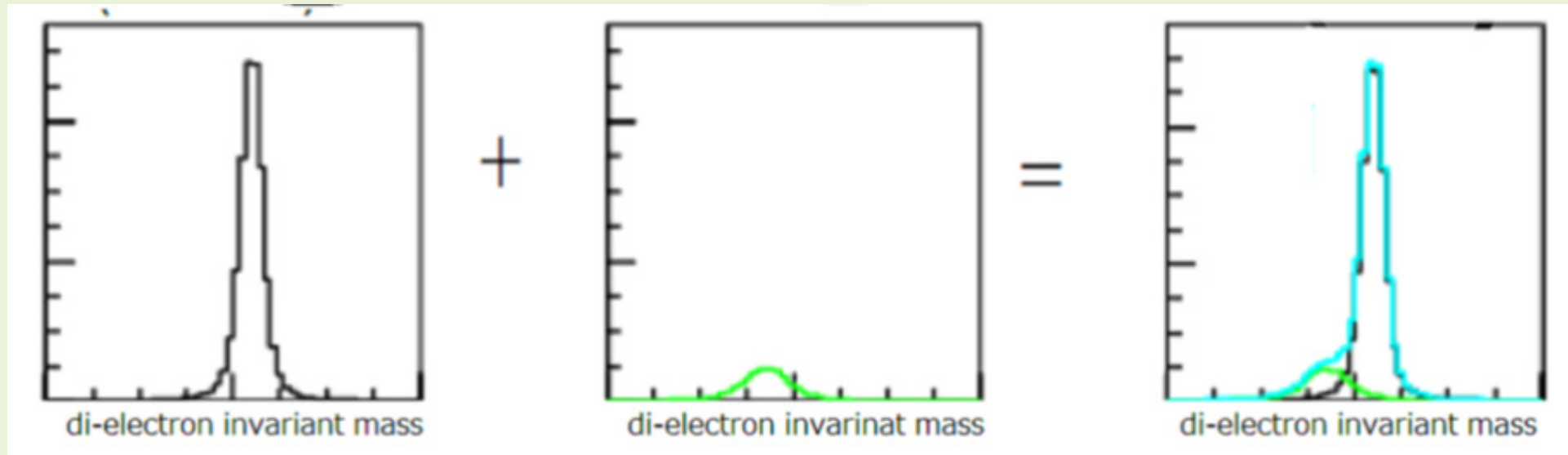
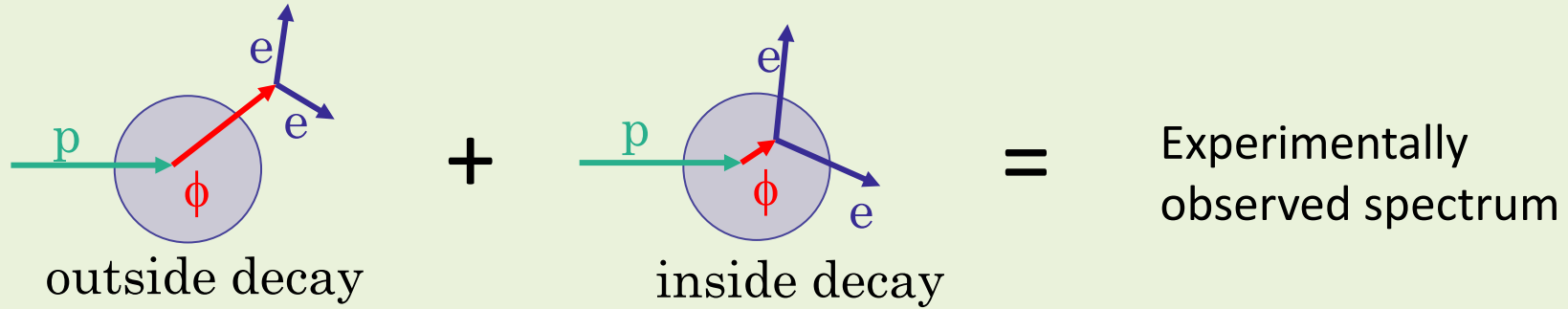
Dilepton spectrum:

$$\rho(\omega) = \underbrace{a\omega^2 + b\omega + c}_{\text{Background}} + \underbrace{A\rho_{\phi, \text{PHSD}}(\omega)}_{\phi \text{ meson signal}}$$

Fitted to the experimental dilepton spectrum independently for each  $\beta\gamma$ -region

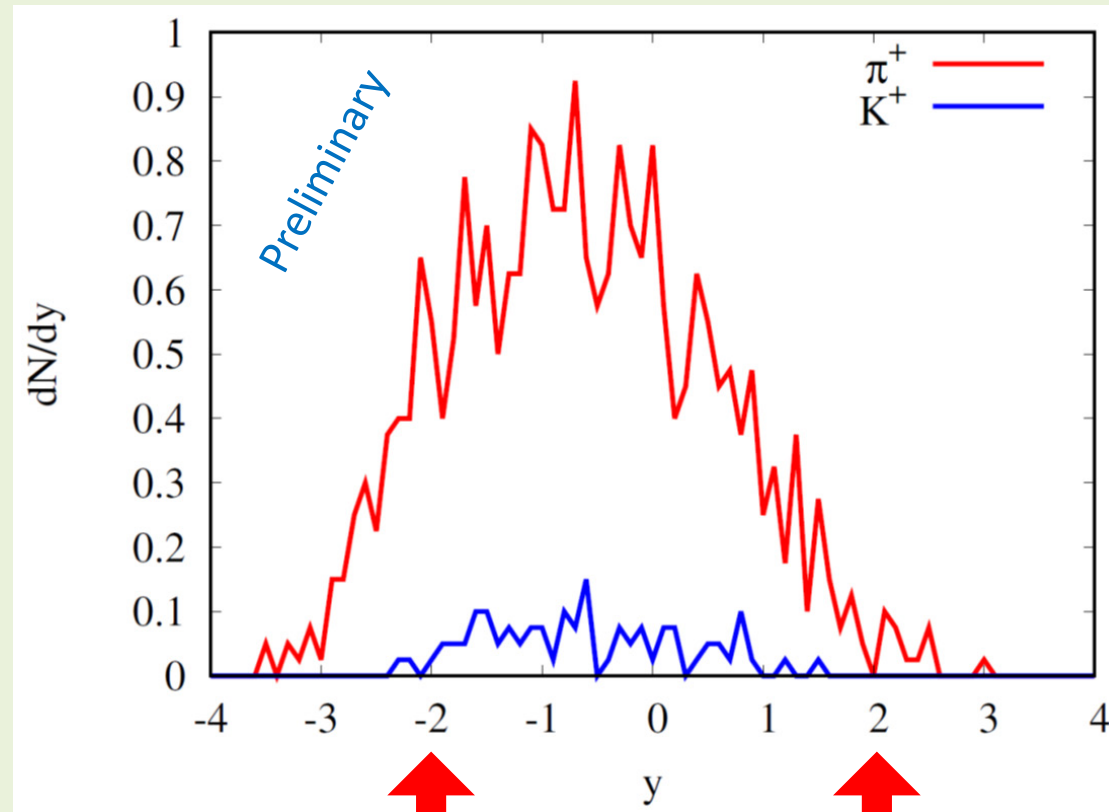


# Experimental di-lepton spectrum



# A first look at a reaction to be probed at J-PARC: pA collisions with initial proton energy of 30 GeV

A first look at the reaction:  
Rapidity distribution of  
protons/mesons



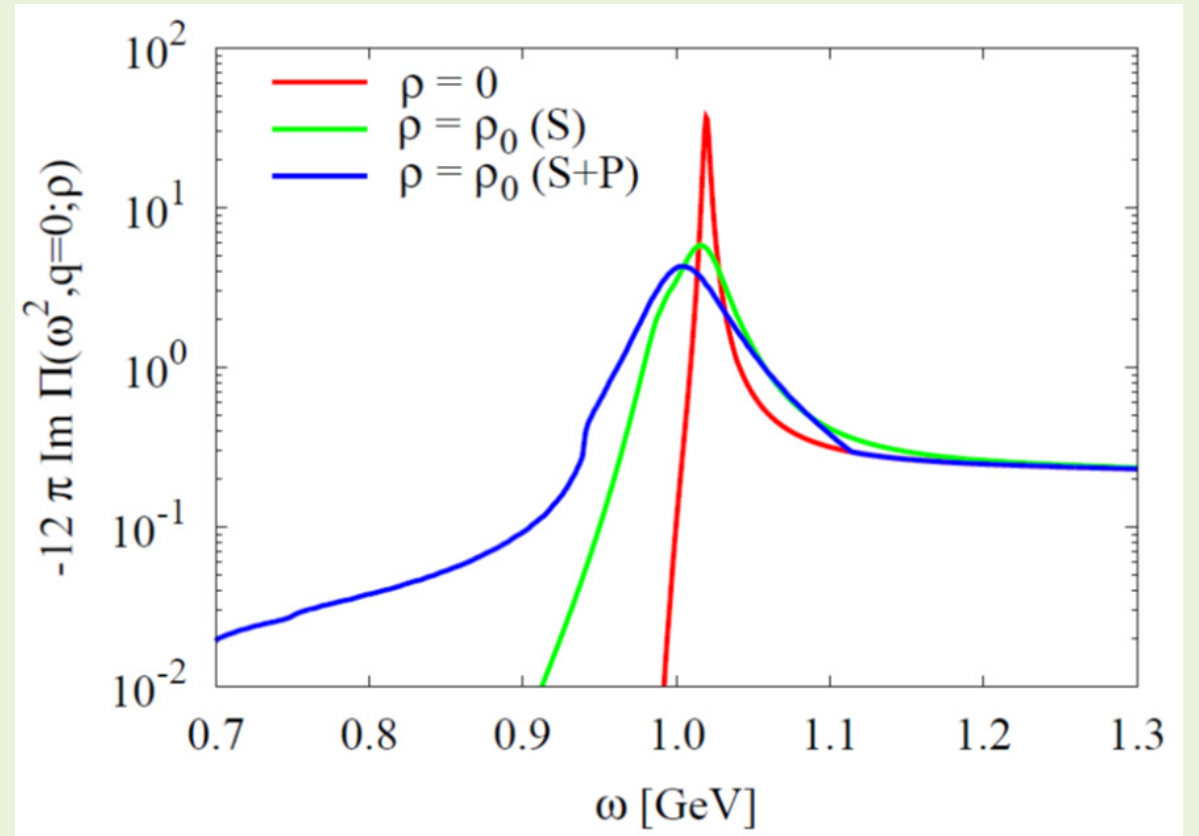
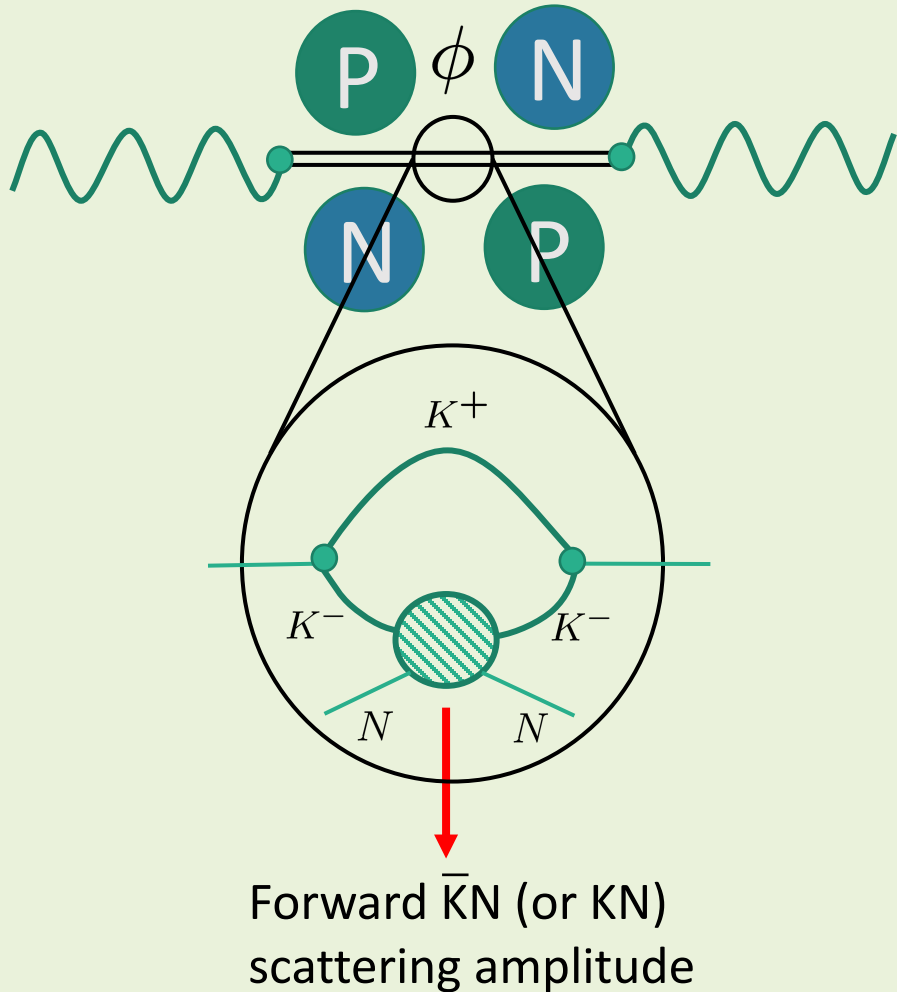
Due to the large collision  
energy, the incoming  
proton passes through the  
target nucleus

nucleon target  
after collision

projectile proton  
after collision

# Recent theoretical works about the $\phi$

based on hadronic models

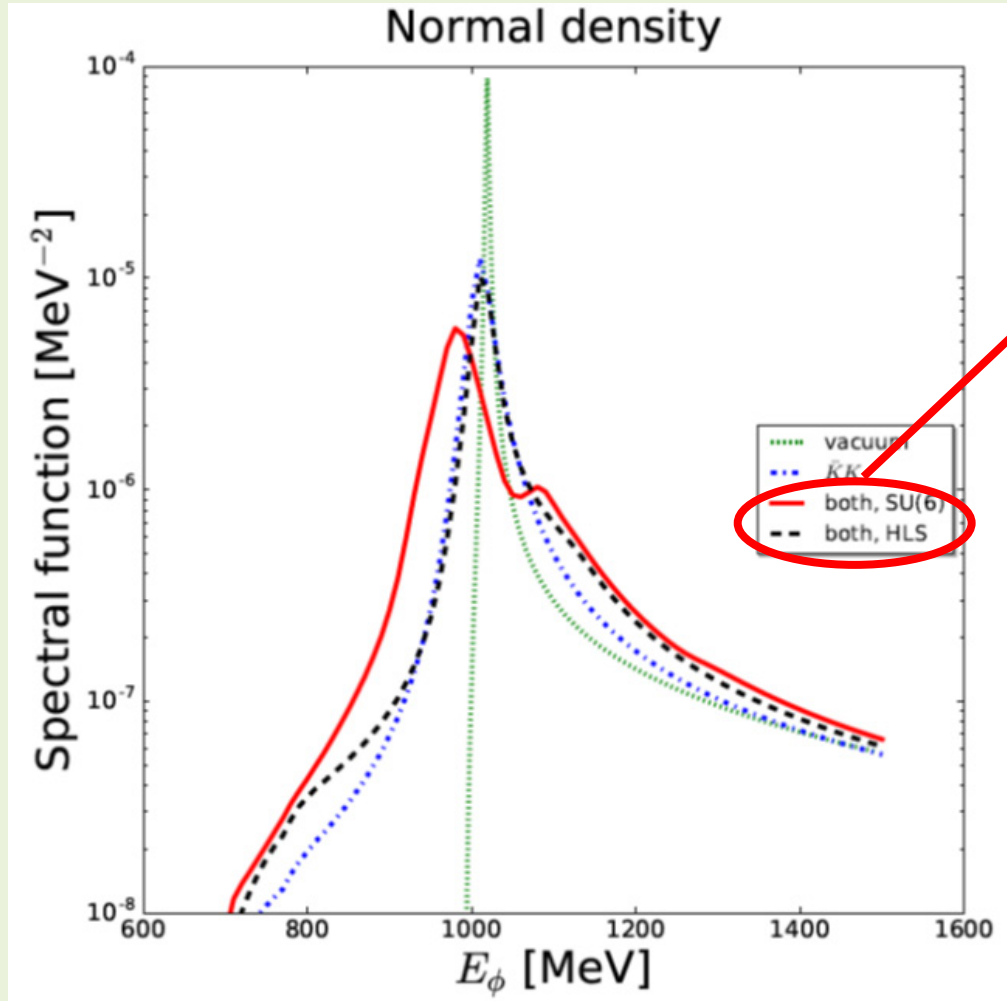


P. Gubler and W. Weise, Phys. Lett. B **751**, 396 (2015).

P. Gubler and W. Weise, Nucl. Phys. A **954**, 125 (2016).

# Recent theoretical works about the $\phi$

based on hadronic models



large dependence on details of the model incorporating Baryon - Vector meson interaction

SU(6): Spin-Flavor Symmetry extension of standard flavor SU(3)

HLS: Hidden Local Symmetry

Common features:

strong broadening, small negative mass shift

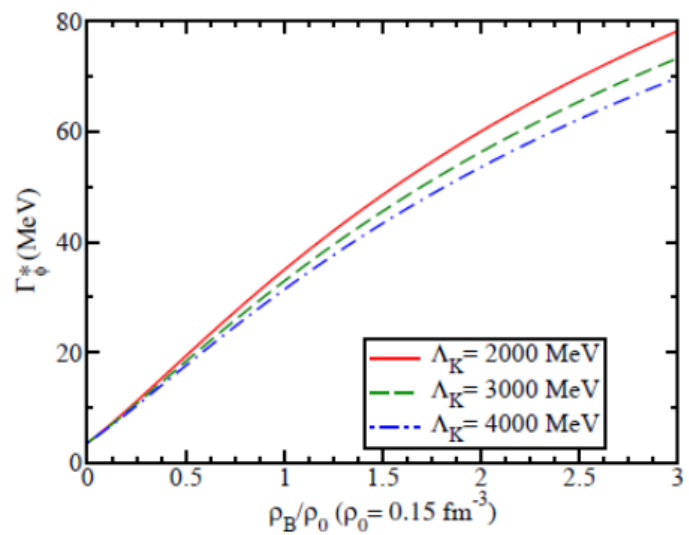
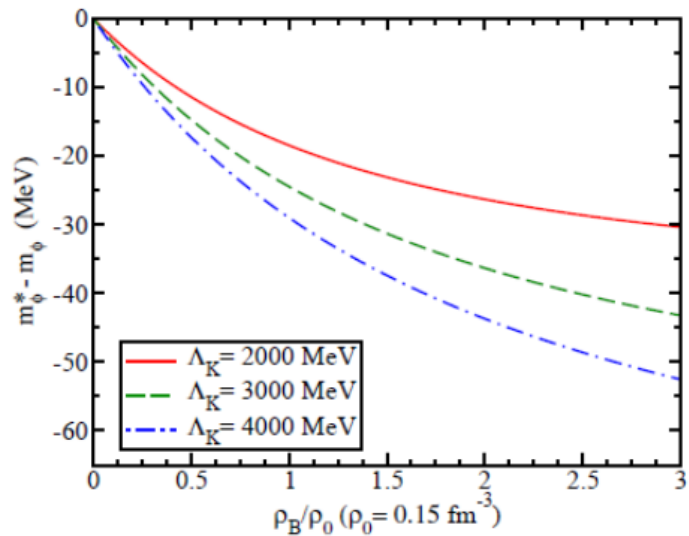
D. Cabrera, A.N. Hiller Blin and M.J. Vicente Vacas, Phys. Rev. C **95**, 015201 (2017).

See also:

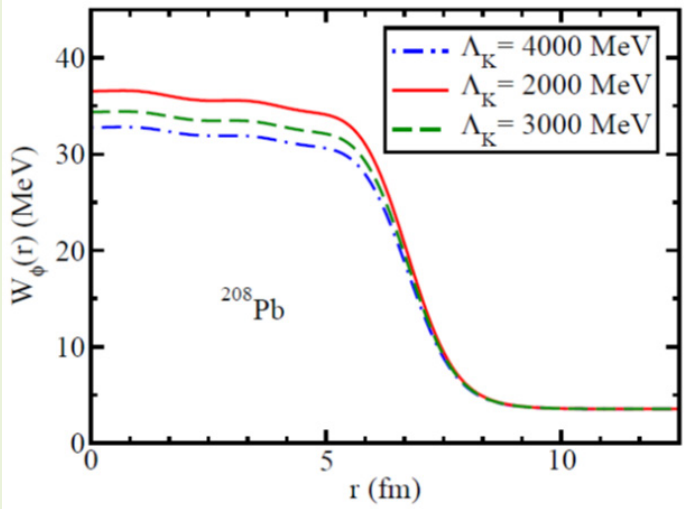
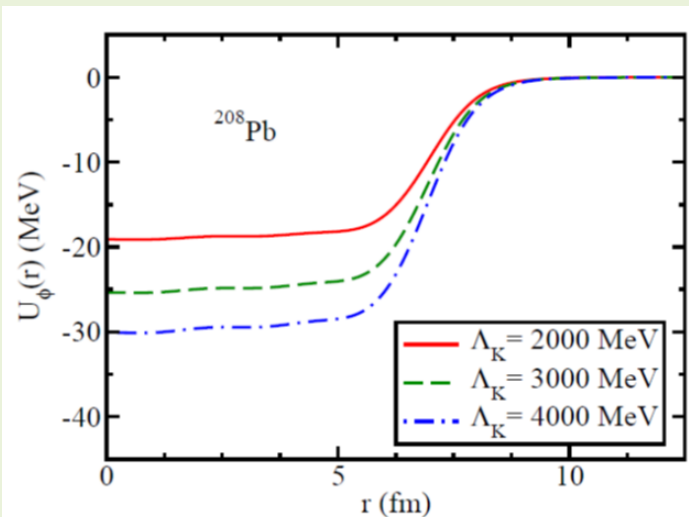
D. Cabrera, A.N. Hiller Blin and M.J. Vicente Vacas, Phys. Rev. C **96**, 034618 (2017).

# Recent theoretical works about the $\phi$

based on the quark-meson coupling model



$$V_{\phi A}(r) = U_\phi(r) - \frac{i}{2}W_\phi(r)$$



		$\Lambda_K = 2000$		$\Lambda_K = 3000$		$\Lambda_K = 4000$	
		$E$	$\Gamma/2$	$E$	$\Gamma/2$	$E$	$\Gamma/2$
$^4\text{He}$	1s	n (-0.8)	n	n (-1.4)	n	-1.0 (-3.2)	8.3
$^{12}\text{C}$	1s	-2.1 (-4.2)	10.6	-6.4 (-7.7)	11.1	-9.8 (-10.7)	11.2
$^{16}\text{O}$	1s	-4.0 (-5.9)	12.3	-8.9 (-10.0)	12.5	-12.6 (-13.4)	12.4
	1p	n (n)	n	n (n)	n	n (-1.5)	n
$^{40}\text{Ca}$	1s	-9.7 (-11.1)	16.5	-15.9 (-16.7)	16.2	-20.5 (-21.2)	15.8
	1p	-1.0 (-3.5)	12.9	-6.3 (-7.8)	13.3	-10.4 (-11.4)	13.3
	1d	n (n)	n	n (n)	n	n (-1.4)	n
$^{48}\text{Ca}$	1s	-10.5 (-11.6)	16.5	-16.5 (-17.2)	16.0	-21.1 (-21.6)	15.6
	1p	-2.5 (-4.6)	13.6	-7.9 (-9.2)	13.7	-12.0 (-12.9)	13.6
	1d	n (n)	n	n (-0.8)	n	-2.1 (-3.6)	11.1
$^{90}\text{Zr}$	1s	-12.9 (-13.6)	17.1	-19.0 (-19.5)	16.4	-23.6 (-24.0)	15.8
	1p	-7.1 (-8.4)	15.5	-12.8 (-13.6)	15.2	-17.2 (-17.8)	14.8
	1d	-0.2 (-2.5)	13.4	-5.6 (-6.9)	13.5	-9.7 (-10.6)	13.4
	2s	n (-1.4)	n	-3.4 (-5.1)	12.6	-7.4 (-8.5)	12.7
	2p	n (n)	n	n (n)	n	n (-1.1)	n
$^{208}\text{Pb}$	1s	-15.0 (-15.5)	17.4	-21.1 (-21.4)	16.6	-25.8 (-26.0)	16.0
	1p	-11.4 (-12.1)	16.7	-17.4 (-17.8)	16.0	-21.9 (-22.2)	15.5
	1d	-6.9 (-8.1)	15.7	-12.7 (-13.4)	15.2	-17.1 (-17.6)	14.8
	2s	-5.2 (-6.6)	15.1	-10.9 (-11.7)	14.8	-15.2 (-15.8)	14.5
	2p	n (-1.9)	n	-4.8 (-6.1)	13.5	-8.9 (-9.8)	13.4
	2d	n (n)	n	n (-0.7)	n	-2.2 (-3.7)	11.9

J.J. Cobos-Martinez, K. Tsushima, G. Krein and A.W. Thomas, Phys. Lett. B **771**, 113 (2017).  
 J.J. Cobos-Martinez, K. Tsushima, G. Krein and A.W. Thomas, Phys. Rev. C **96**, 035201 (2017).

Some  $\phi A$  bound states might exist, but they have a large width  
 → difficult to observe experimentally?

# Our tool: a transport code

## PHSD (Parton Hadron String Dynamics)

W. Cassing and E. Bratkovskaya, Phys. Rev. C **78**, 034919 (2008).

W. Cassing and E. Bratkovskaya, Phys. Rept. **308**, 65 (1999).

W. Cassing, V. Metag, U. Mosel and K. Niita, Phys. Rept. **188**, 363 (1990).

Basic Ingredient 1: Solve a Vlasov-Uehling-Uhlenbeck type equation for each particle type

$$\left( \frac{\partial}{\partial t} + \frac{\mathbf{p}_1}{m} \cdot \frac{\partial}{\partial \mathbf{r}} - \frac{\partial}{\partial \mathbf{r}} U_{\text{BHF}}(\mathbf{r}; t) \cdot \frac{\partial}{\partial \mathbf{p}_1} \right) f(\mathbf{r}, \mathbf{p}_1; t) = \left( \frac{\partial f}{\partial t} \right)_{\text{coll}}$$

mean field  
(tuned to reproduce  
nuclear matter properties)

particle distribution  
function

Basic Ingredient 2: „Testparticle“ approach



$$f_h(\mathbf{r}, \mathbf{p}; t) = \frac{1}{N_{\text{test}}} \sum_i^{N_h(t) \times N_{\text{test}}} \delta(\mathbf{r} - \mathbf{r}_i(t)) \delta(\mathbf{p} - \mathbf{p}_i(t))$$

# Structure of QCD sum rules for the phi meson

$$\frac{1}{M^2} \int_0^\infty ds e^{-\frac{s}{M^2}} \rho(s) = c_0(\rho) + \frac{c_2(\rho)}{M^2} + \frac{c_4(\rho)}{M^4} + \frac{c_6(\rho)}{M^6} + \dots$$

In Vacuum

$$\text{Dim. 0: } c_0(0) = 1 + \frac{\alpha_s}{\pi}$$

$$\text{Dim. 2: } c_2(0) = -6m_s^2$$

$$\text{Dim. 4: } c_4(0) = \frac{\pi^2}{3} \langle \frac{\alpha_s}{\pi} G^2 \rangle + 8\pi^2 m_s \langle \bar{s}s \rangle$$

$$\text{Dim. 6: } c_6(0) = -\frac{448}{81} \kappa \pi^3 \alpha_s \langle \bar{s}s \rangle^2$$

# Structure of QCD sum rules for the phi meson

$$\frac{1}{M^2} \int_0^\infty ds e^{-\frac{s}{M^2}} \rho(s) = c_0(\rho) + \frac{c_2(\rho)}{M^2} + \frac{c_4(\rho)}{M^4} + \frac{c_6(\rho)}{M^6} + \dots$$

## In Nuclear Matter

Dim. 0:  $c_0(\rho) = c_0(0)$

$$\langle \bar{s}s \rangle_\rho = \langle \bar{s}s \rangle_0 + \langle N | \bar{s}s | N \rangle \rho + \dots$$

Dim. 2:  $c_2(\rho) = c_2(0)$

Dim. 4: 
$$c_4(\rho) = c_4(0) + \rho \left[ -\frac{2}{27} M_N + \frac{56}{27} m_s \langle N | \bar{s}s | N \rangle \right. \\ \left. + \frac{4}{27} m_q \langle N | \bar{q}q | N \rangle + A_2^s M_N - \frac{7}{12} \frac{\alpha_s}{\pi} A_2^g M_N \right]$$

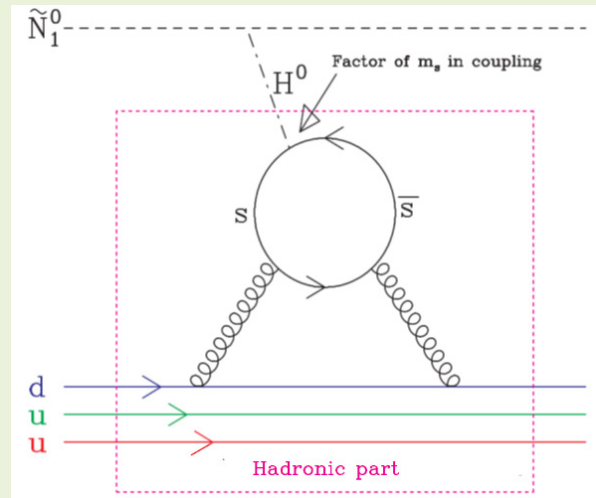
Dim. 6: 
$$c_6(\rho) = c_6(0) + \rho \left[ -\frac{896}{81} \kappa_N \pi^3 \alpha_s \langle \bar{s}s \rangle \langle N | \bar{s}s | N \rangle - \frac{5}{6} A_4^s M_N^3 \right]$$



# The strangeness content of the nucleon: $\sigma_{sN} = m_s \langle N | \bar{s}s | N \rangle$

Important parameter for dark-matter searches!

Neutralino:  
Linear superposition of the  
Super-partners of the Higgs, the  
photon and the Z-boson



Adapted from:  
W. Freeman and D. Toussaint (MILC Collaboration),  
Phys. Rev. D **88**, 054503 (2013).

$$\sigma_{\text{scalar}}^{(\text{nucleon})} = \frac{8G_F^2}{\pi} M_Z^2 m_{\text{red}}^2 \left[ \frac{F_h I_h}{m_h^2} + \frac{F_H I_H}{m_H^2} \frac{M_Z}{2} \sum_q \langle N | \bar{q}q | N \rangle \sum_i P_{\tilde{q}_i} (A_{\tilde{q}_i}^2 - B_{\tilde{q}_i}^2) \right]^2$$

most important contribution

$$I_{h,H} = k_{u\text{-type}}^{h,H} g_u + k_{d\text{-type}}^{h,H} g_d$$

dominates

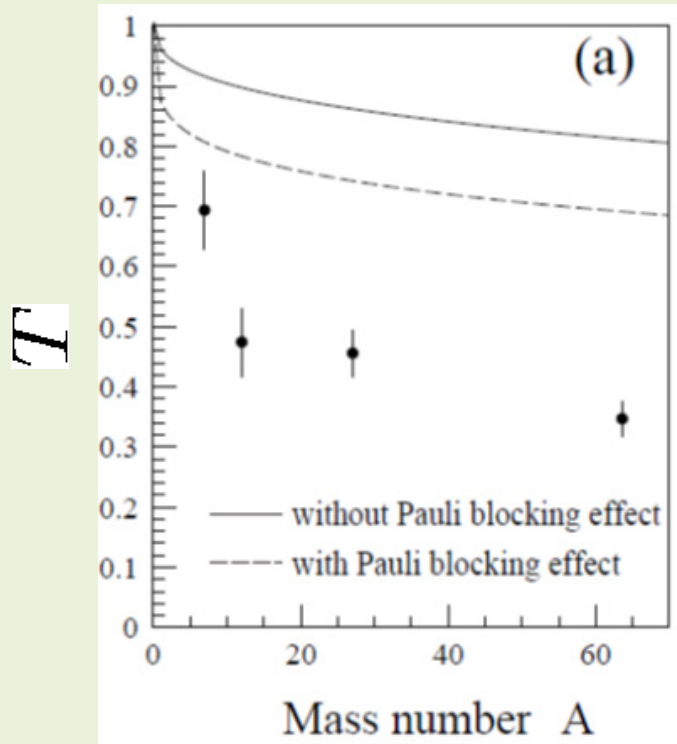
$$g_d = \frac{2}{27} \left( m_N + \frac{23}{4} \sigma_{\pi N} + \frac{25}{2} \sigma_{sN} \right)$$

# Other experimental results

There are some more experimental results on the  $\phi$ -meson width in nuclear matter, based on the measurement of the transparency ratio  $T$ :

$$T = \frac{\sigma_{\gamma A \rightarrow \phi X}}{A \sigma_{\gamma N \rightarrow \phi X}}$$

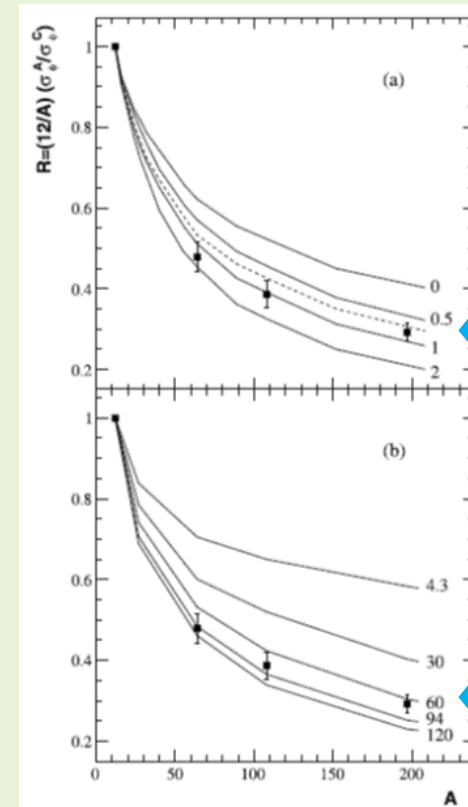
Measured at SPring-8 (LEPS)



$\Gamma_{\phi}(\rho_0) \simeq 30 \text{ MeV}$

Theoretical calculation:  
D. Cabrera, L. Roca, E. Oset,  
H. Toki and M.J. Vicente Vacas,  
Nucl. Phys. **A733**, 130 (2004).

Measured at COSY-ANKE



Theoretical calculation:  
V.K. Magas, L. Roca and E. Oset,  
Phys. Rev. C **71**, 065202 (2005).

$\Gamma_{\phi}(\rho_0) \simeq 27 \text{ MeV}$

Theoretical calculation:  
E. Ya. Paryev,  
J. Phys. G **36**, 015103 (2009).

$\Gamma_{\phi}(\rho_0) \simeq 73 \text{ MeV}$