

Strangeness in Neutron Stars

Theo F. Motta

University of Adelaide

July 14, 2021

What do we know about the
content of NS cores *beyond
significant doubt?*

- 1 A Brief History of the Hyperon Crisis
- 2 The Quark-Meson Coupling Model
- 3 The Nature of Strangeness Puzzle

A Brief History of the Hyperon Crisis

Brief History

Brief History

- 1932 The neutron is discovered (Chadwick)

Brief History

- 1932 The neutron is discovered (Chadwick)
- 1934 Baade & Zwicky propose the idea of a neutron star

Brief History

- 1932 The neutron is discovered (Chadwick)
- 1934 Baade & Zwicky propose the idea of a neutron star
- 1939 TOV mass limit

Brief History

- 1932 The neutron is discovered (Chadwick)
- 1934 Baade & Zwicky propose the idea of a neutron star
- 1939 TOV mass limit
- 1967 Jocelyn Bell & Anthony Hewish discover PSR B1919+21

In the beginning there was the TOV limit

$$\frac{dp}{dr} = -\frac{G}{r^2} \left(\epsilon(p(r)) + \frac{p}{c^2} \right) \left(M + 4\pi r^3 \frac{p}{c^2} \right) \left(1 - \frac{2GM}{c^2 r} \right)^{-1}.$$

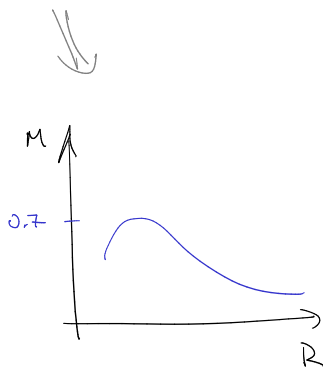
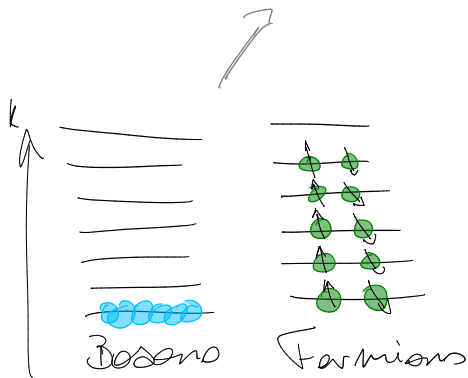
In the beginning there was the TOV limit

$$\frac{dp}{dr} = -\frac{G}{r^2} \left(\epsilon(p(r)) + \frac{p}{c^2} \right) \left(M + 4\pi r^3 \frac{p}{c^2} \right) \left(1 - \frac{2GM}{c^2 r} \right)^{-1}.$$

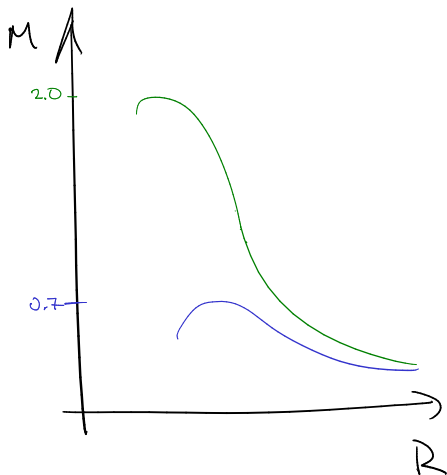


In the beginning there was the TOV limit

$$\frac{dp}{dr} = -\frac{G}{r^2} \left(\epsilon(p(r)) + \frac{p}{c^2} \right) \left(M + 4\pi r^3 \frac{p}{c^2} \right) \left(1 - \frac{2GM}{c^2 r} \right)^{-1}.$$

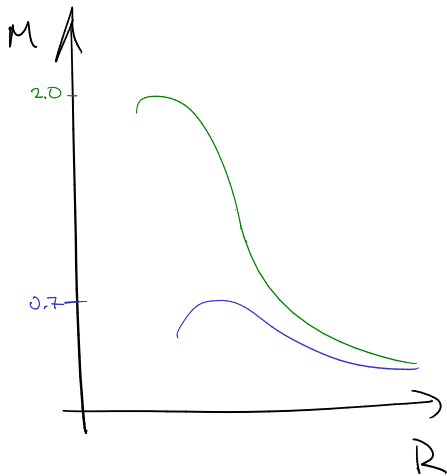


Beyond Degenerate Fermi Gas



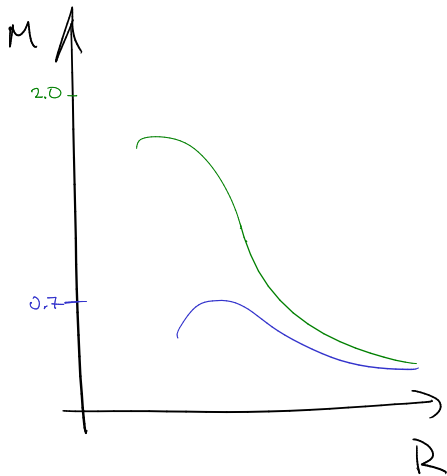
- So, if you include interactions in our degenerate Fermi Gas you get much more pressure and higher masses.

Beyond Degenerate Fermi Gas



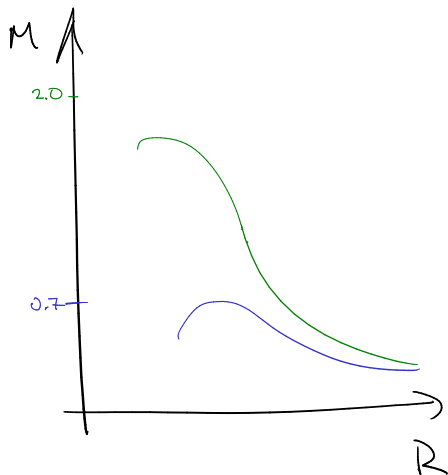
- So, if you include interactions in our degenerate Fermi Gas you get much more pressure and higher masses.
- However, if the mass is that high, energy density must be huge. Well enough to produce strangeness...

Beyond Degenerate Fermi Gas



- So, if you include interactions in our degenerate Fermi Gas you get much more pressure and higher masses.
- However, if the mass is that high, energy density must be huge. Well enough to produce strangeness...

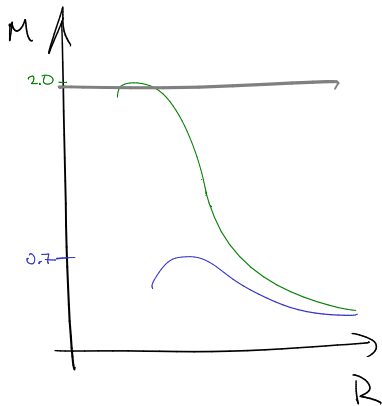
Beyond Degenerate Fermi Gas



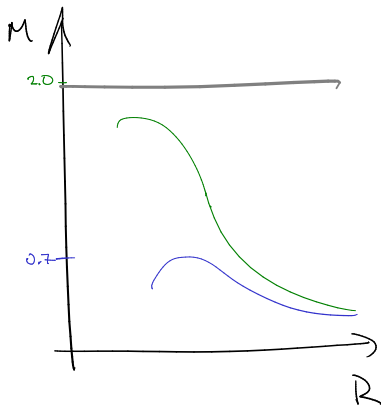
- PSR J1614-2230
 $M = 1.908 \pm 0.016 M_{\odot}$
- PSR J0348+0432
 $M = 2.01 \pm 0.04 M_{\odot}$
- PSR J0740+6620
 $M = 2.08 \pm 0.07 M_{\odot}$

Hyperon Puzzle in a Nutshell

Hyperon Puzzle in a Nutshell



Mass good
Content bad



Mass bad
Content good

Some Solutions?

Some Solutions?

- Modify gravity!

Some Solutions?

- Modify gravity!
- Extra hyperon repulsion

Some Solutions?

- Modify gravity!
- Extra hyperon repulsion
- Kaon condensation

Some Solutions?

- Modify gravity!
- Extra hyperon repulsion
- Kaon condensation
- Also the deconfinement question...

The Quark-Meson Coupling Model

- The QMC model is one of the relativistic phenomenological models. What we want, fundamentally, is

$$\mathcal{L} = \bar{\Psi}_B(i\not{\partial} - M_B)\Psi_B + \mathcal{L}_{\text{mesons}} + \mathcal{L}_{\text{interactions}}$$

- The QMC model is one of the relativistic phenomenological models. What we want, fundamentally, is

$$\mathcal{L} = \bar{\Psi}_B(i\not{\partial} - M_B)\Psi_B + \mathcal{L}_{\text{mesons}} + \mathcal{L}_{\text{interactions}}$$

- However, we also want to retain some information on the *baryon structure*.

- The QMC model is one of the relativistic phenomenological models. What we want, fundamentally, is

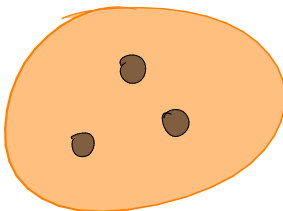
$$\mathcal{L} = \bar{\Psi}_B(i\not{\partial} - M_B)\Psi_B + \mathcal{L}_{\text{mesons}} + \mathcal{L}_{\text{interactions}}$$

- However, we also want to retain some information on the *baryon structure*.
- So before writing the baryon level lagrangian, let's take a look at the baryon structure, starting from a simple bag model perspective

- The QMC model is one of the relativistic phenomenological models. What we want, fundamentally, is

$$\mathcal{L} = \bar{\Psi}_B(i\not{\partial} - M_B)\Psi_B + \mathcal{L}_{\text{mesons}} + \mathcal{L}_{\text{interactions}}$$

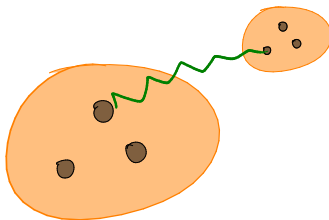
- However, we also want to retain some information on the *baryon structure*.
- So before writing the baryon level lagrangian, let's take a look at the baryon structure, starting from a simple bag model perspective



- The QMC model is one of the relativistic phenomenological models. What we want, fundamentally, is

$$\mathcal{L} = \bar{\Psi}_B(i\not{\partial} - M_B)\Psi_B + \mathcal{L}_{\text{mesons}} + \mathcal{L}_{\text{interactions}}$$

- However, we also want to retain some information on the *baryon structure*.
- We choose to model baryon-baryon interactions as a quark-meson interaction in the subhadronic level.



Bag Model

- Solve Dirac with the appropriate boundary conditions

$$(i\cancel{\partial} - m)\psi = 0$$

Bag Model

- Solve Dirac with the appropriate boundary conditions

$$(i\cancel{\partial} - m)\psi = 0$$

- In our case, however, the quarks interact with mesons.

Bag Model

- Solve Dirac with the appropriate boundary conditions

$$(i\cancel{\partial} - m)\psi = 0$$

- In our case, however, the quarks interact with mesons.
- In MFA we have:

$$(i\cancel{\partial} - m^*)\psi = 0$$

Bag Model

- Solve Dirac with the appropriate boundary conditions

$$(i\cancel{\partial} - m)\psi = 0$$

- In our case, however, the quarks interact with mesons.
- In MFA we have:

$$(i\cancel{\partial} - m^*)\psi = 0$$

where $m^* = m - g_\sigma^q \bar{\sigma}$

- We then calculate the energy of a stationary bag, that is, the mass of the baryon

- We then calculate the energy of a stationary bag, that is, the mass of the baryon

$$M_B$$

- We then calculate the energy of a stationary bag, that is, the mass of the baryon

$$M_B^*$$

- We then calculate the energy of a stationary bag, that is, the mass of the baryon

$$M_B^*(\vec{\sigma})$$

- We then calculate the energy of a stationary bag, that is, the mass of the baryon

$$M_B^*(\bar{\sigma}) = M_B - g_\sigma^B \bar{\sigma} + \frac{d}{2} (g_\sigma^B \bar{\sigma})^2$$

- We then calculate the energy of a stationary bag, that is, the mass of the baryon

$$M_B^*(\vec{\sigma}) = M_B - g_\sigma^B \vec{\sigma} + \frac{d}{2} (g_\sigma^B \vec{\sigma})^2$$

- We can now go back to the baryon level and construct a Lagrangian

- We then calculate the energy of a stationary bag, that is, the mass of the baryon

$$M_B^*(\vec{\sigma}) = M_B - g_\sigma^B \vec{\sigma} + \frac{d}{2} (g_\sigma^B \vec{\sigma})^2$$

- We can now go back to the baryon level and construct a Lagrangian

$$\mathcal{L} = \bar{\Psi}_B (i\not{\partial} - M_B) \Psi_B$$

- We then calculate the energy of a stationary bag, that is, the mass of the baryon

$$M_B^*(\vec{\sigma}) = M_B - g_\sigma^B \vec{\sigma} + \frac{d}{2} (g_\sigma^B \vec{\sigma})^2$$

- We can now go back to the baryon level and construct a Lagrangian

$$\mathcal{L} = \bar{\Psi}_B (i\not{\partial} - M_B) \Psi_B + g_\sigma^B \bar{\Psi} \Psi$$

- We then calculate the energy of a stationary bag, that is, the mass of the baryon

$$M_B^*(\bar{\sigma}) = M_B - g_\sigma^B \bar{\sigma} + \frac{d}{2} (g_\sigma^B \bar{\sigma})^2$$

- We can now go back to the baryon level and construct a Lagrangian

$$\mathcal{L} = \bar{\Psi}_B (i\not{\partial} - M_B) \Psi_B + g_\sigma^B \sigma \bar{\Psi} \Psi - \frac{d}{2} (g_\sigma^B \sigma)^2 \bar{\Psi} \Psi$$

- We then calculate the energy of a stationary bag, that is, the mass of the baryon

$$M_B^*(\bar{\sigma}) = M_B - g_\sigma^B \bar{\sigma} + \frac{d}{2} (g_\sigma^B \bar{\sigma})^2$$

- We can now go back to the baryon level and construct a Lagrangian

$$\mathcal{L} = \bar{\Psi}_B (i\not{\partial} - M_B) \Psi_B + g_\sigma^B \sigma \bar{\Psi} \Psi - \frac{d}{2} (g_\sigma^B \sigma)^2 \bar{\Psi} \Psi + \dots$$

Some QMC Highlights

Some QMC Highlights

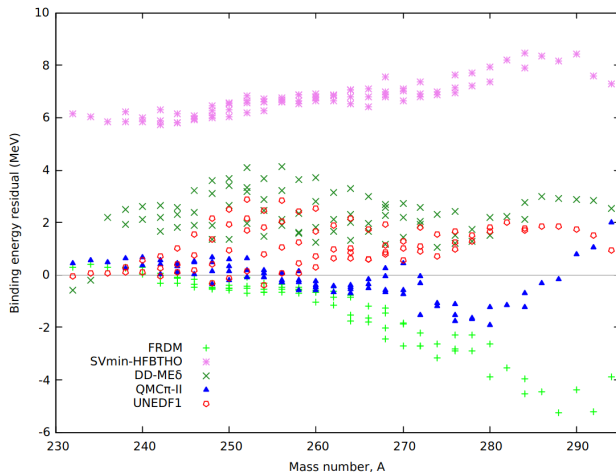


Figure: Nuclei binding energies from Martinez et al. [2019]

Some QMC Highlights

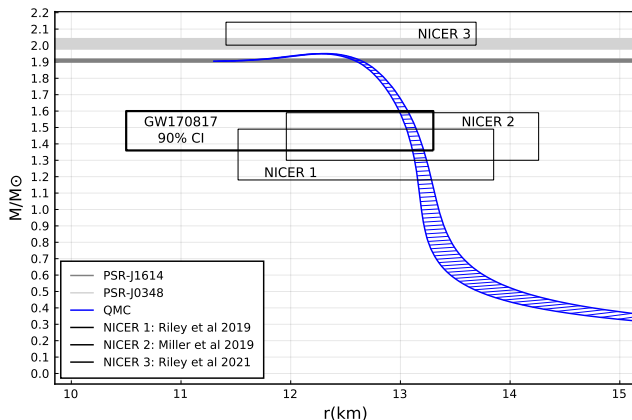


Figure: Stellar structure results from Motta et al. [2021] plus recent NICER results such as et al [2021]

The Nature of Strangeness Puzzle

- The QMC and other models have shown good agreement with masses and radii *including hyperons*

- The QMC and other models have shown good agreement with masses and radii *including hyperons*
- The Hyperon puzzle as it was first conceived of doesn't seem like too much of a puzzle anymore

- The QMC and other models have shown good agreement with masses and radii *including hyperons*
- The Hyperon puzzle as it was first conceived of doesn't seem like too much of a puzzle anymore
- However, one puzzle to be addressed is that of deconfinement.

- The QMC and other models have shown good agreement with masses and radii *including hyperons*
- The Hyperon puzzle as it was first conceived of doesn't seem like too much of a puzzle anymore
- However, one puzzle to be addressed is that of deconfinement.
- One interesting approach is to try the inverse problem, get NS constraints and from there infer what the EoS is. See Annala et al. [2020]

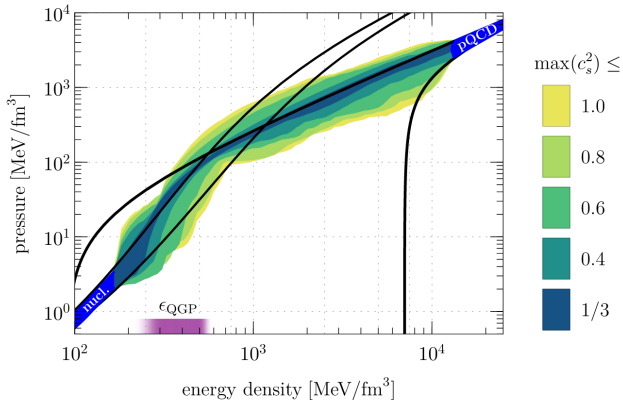


Figure: Results from Annala et al. [2020]

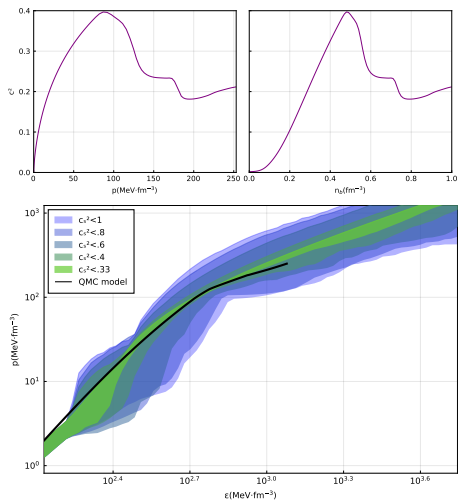
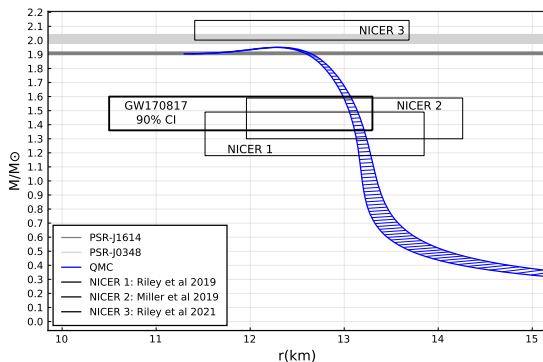


Figure: Motta et al. [2021] comparing QMC with Annala et al. [2020]

Conclusions

- The QMC model is in agreement with data
- There isn't much of a Hyperon puzzle anymore
- We still cannot discriminate between QM and Hyperonic matter



References

- Kay Marie L. Martinez, Anthony William Thomas, Jirina R. Stone, and Pierre A. M. Guichon. Parameter optimization for the latest quark-meson coupling energy-density functional. *Phys. Rev. C*, 100(2):024333, 2019. doi: [10.1103/PhysRevC.100.024333](https://doi.org/10.1103/PhysRevC.100.024333).
- T. F. Motta, P. A. M. Guichon, and A. W. Thomas. On the sound speed in hyperonic stars. *Nucl. Phys. A*, 1009:122157, 2021. doi: [10.1016/j.nuclphysa.2021.122157](https://doi.org/10.1016/j.nuclphysa.2021.122157).
- Thomas E. Riley et al. A nicer view of the massive pulsar psr j0740+6620 informed by radio timing and xmm-newton spectroscopy, 2021.
- Eemeli Annala, Tyler Gorda, Alekski Kurkela, Joonas Nättilä, and Alekski Vuorinen. Evidence for quark-matter cores in massive neutron stars. *Nature Phys.*, 16(9):907–910, 2020. doi: [10.1038/s41567-020-0914-9](https://doi.org/10.1038/s41567-020-0914-9).